

1 **Precipitation driving of droplet concentration**
2 **variability in marine low clouds**

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22 **June 14th 2012**
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24

25 **Abstract**

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27 The concentration N_d of cloud droplets in marine low clouds is a primary determinant of
28 their ability to reflect sunlight and modulates their ability to precipitate. Previous studies
29 have focused upon aerosol source variability as the key driver of variability in N_d . Here,
30 we use a highly simplified aerosol budget model to examine the impact of precipitation
31 on N_d . This model considers: precipitation (coalescence) scavenging, constrained using
32 new satellite measurements of light precipitation; entrainment of aerosol from above
33 cloud combined with constant aerosol concentration based on recent field observations of
34 aerosol particles in the free troposphere; and sea-surface aerosol production estimated
35 using a wind-speed dependent source function. Despite the highly simplified nature of
36 this model, it skillfully predicts the geographical variability of N_d in regions of extensive
37 marine low clouds. Inclusion of precipitation results in reduction in N_d by factors of 2-3
38 over the remote oceans. Within 500 km of coastlines the reduction in N_d due to
39 precipitation is weak but in these regions the model is not able to accurately predict N_d
40 because of strong pollution sources. In general, neither free-tropospheric nor surface
41 CCN sources alone are sufficient to maintain N_d against precipitation losses. The results
42 demonstrate that even the light precipitation rates typical of marine stratocumulus
43 profoundly impact the radiative properties of marine low clouds.

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47 **1. Introduction**

48 Anthropogenic activities have resulted in marked increases in the concentration of
49 aerosol particles in the atmosphere [*Kaufman et al. 2002, Isaksen et al. 2009*] and these
50 increases exert a significant but highly uncertain radiative forcing on the global climate
51 [*Isaksen et al. 2009, IPCC 2007*]. A large fraction of this forcing is attributed to the
52 effects that aerosol particles have on clouds by increasing the concentration N_d of cloud
53 droplets [*Martin et al. 1994; Ramanathan et al. 2001; Lohmann and Feichter 2005*],
54 reducing droplet size [*Breon et al. 2002*], thereby increasing the reflected solar radiation
55 [*Twomey 1974, Penner et al. 2004, Quaas et al. 2009*]. Systematic increases in N_d have
56 been observed downwind of east Asia over the past two decades [*Bennartz et al., 2011*]
57 and have been attributed to rapid industrialization. The magnitude of the so-called
58 “aerosol indirect effect” on climate depends not only upon present-day conditions, but
59 also upon the unperturbed microphysical state of the clouds prior to the addition of
60 anthropogenic aerosols [*Platnick and Twomey 1994, Oreopoulos and Platnick, 2008*]. It is
61 reasonable to argue that N_d is the single most important cloud microphysical variable that
62 must be accurately represented in models in order to accurately determine aerosol indirect
63 effects on climate. However, there are marked differences between values of N_d in
64 different climate models [*Quaas et al. 2009, Ming et al. 2006, Gettelman et al. 2008*]
65 demonstrating a clear lack of understanding of the key controls on N_d .

66 Satellite-based studies use the relationship between observed cloud droplet size or
67 concentration, and nearby clear-sky estimates of aerosol loading, to infer the role that
68 aerosols play in influencing clouds and climate [*Breon et al. 2002*]. Some even go so far
69 as to quantitatively estimate aerosol indirect effects globally [*Quaas et al. 2008, Jones et*

70 *al.* 2009]. Inherent in this approach is that correlations between cloud microphysical
71 properties and aerosols in the current climate are indicative of an aerosol influence on
72 cloud properties rather than vice versa. Aerosol-cloud correlative studies do not take the
73 possible effects of precipitation into account. One can make a reasonable case that
74 precipitation-induced aerosol changes will not significantly impact the inferences drawn
75 from these studies only if one assumes that the impacts of precipitation are localized,
76 intermittent, and relatively rare, and that the aerosol fields that interact with the majority
77 of clouds are not significantly affected by precipitation. However, recent observations
78 from the sensitive spaceborne radar on the CloudSat satellite are finding that precipitation
79 occurs more frequently over the globe than previously thought [*Leon et al.* 2008, *Haynes*
80 *et al.* 2009] prompting an examination of the role of precipitation in driving aerosol
81 variability.

82 In this study we use these state-of-the-art quantitative estimates of light
83 precipitation from CloudSat to constrain a simple budget model that predicts the mean
84 concentrations of cloud condensation nuclei (CCN) and cloud droplets over those parts of
85 the global oceans containing extensive low clouds. These clouds are confined within the
86 marine boundary layer (MBL) and are among the most susceptible to aerosol
87 perturbations [*Oreopoulos and Platnick*, 2008]. We build upon previous studies [*Baker*
88 *and Charlson* 1990, *Baker et al.* 1993] that used simplified budget models to provide
89 important insights into the factors controlling CCN, by constructing a steady-state budget
90 for CCN in the MBL appropriate for the regions of large-scale subsidence where
91 extensive marine stratocumulus clouds are favored [*Klein and Hartmann*, 1993].

92 This article is organized as follows. Section 2 describes the basis for a simple,
93 single-equation, steady-state budget model to predict CCN or N_d in the MBL. Section 3
94 describes how we determine the terms in the budget using a variety of observations
95 including satellites and in-situ data. Section 4 presents the key results from the model and
96 compares the model against satellite observations of N_d . Section 5 discusses the
97 implications of our findings.

98

99 **2. Steady state budget model**

100 The rate of increase of CCN concentration \dot{N} averaged over the depth of the MBL
101 can be written as the sum of various source and sink terms:

$$102 \quad \dot{N} = \dot{N}_{\text{FT}} + \dot{N}_{\text{S}} + \dot{N}_{\text{PROD}} + \dot{N}_{\text{P}} + \dot{N}_{\text{DRY}} + \dot{N}_{\text{ADV}} \quad [1]$$

103 where \dot{N}_{FT} , \dot{N}_{S} , \dot{N}_{PROD} , \dot{N}_{P} , \dot{N}_{DRY} and \dot{N}_{ADV} are the time tendencies due to
104 entrainment of CCN from the free-troposphere (FT), primary production at the surface
105 (i.e. sea spray), and secondary production, precipitation (i.e. coalescence scavenging), dry
106 deposition to the surface and horizontal advection, respectively.

107 Free-tropospheric air is constantly being mixed into the MBL by cloud top
108 entrainment, and this can either provide a net source of CCN to the MBL or can dilute
109 MBL aerosol concentrations. Modeling and observational studies suggest that the FT is a
110 primary source of CCN in the remote MBL [Clarke *et al.* 1998b, Capaldo *et al.* 1999,
111 Katoshevski *et al.* 1999]. The net source rate is $\dot{N}_{\text{FT}} = w_e(N_{\text{FT}} - N)/z_i$ where z_i is the depth of
112 the MBL and w_e is the entrainment rate.

113 The modeled surface source \dot{N}_{S} is assumed to be from primary production of sea-
114 spray aerosol (SSA) and we use a recent parameterization [Clarke *et al.* 2006] to provide

115 $\dot{N}_S = F(\sigma) U_{10}^{3.41} / z_i$ where U_{10} is the wind speed at a height of 10 m, and $F(\sigma)$ depends upon
116 the assumed peak supersaturation σ experienced in the clouds (see section 3.1 below).
117 We examine the sensitivity to the parameterization of sea-spray by comparing with the
118 frequently-used formulation of *Monahan et al.* [1986]. Since this formulation has the
119 same wind speed dependence as *Clarke et al.* [2006], the expression for \dot{N}_S given above
120 is identical for the two schemes, but the function $F(\sigma)$ is different. We discuss how the
121 supersaturation and $F(\sigma)$ are specified in Sections 3.1 and 3.3 below.

122 The model does not take into account CCN formation from the nucleation of new
123 particles in the MBL since it is unlikely that this contributes significantly to the mean
124 CCN number concentration over the oceans. In severely-scavenged ultraclean MBLs
125 evidence of new particle formation has been noted [*Clarke et al.* 1998, *Petters et al.*
126 2006, *Tomlinson et al.* 2007, *Wood et al.* 2008] . Such nucleation events appear to be
127 quite rare, with only one clear instance observed during four weeks of shipborne
128 sampling over the tropical southeastern Pacific ocean [*Tomlinson et al.* 2007]. In
129 addition, evidence that freshly nucleated particles can grow to sufficient sizes to increase
130 the population of CCN without being scavenged by existing cloud is lacking. Nucleation
131 events are therefore unlikely to compete with other source processes in determining the
132 mean state [*Capaldo et al.* 1999, *Katoshevski et al.* 1999]. In any case, formulations for
133 the rate of production of CCN from new particle formation in the MBL are highly
134 uncertain [*Capaldo et al.* 1999, *Kirkby et al.* 2011]. To maintain simplicity, and because
135 the production rates are highly uncertain, we do not include production of CCN from
136 other secondary processes like aqueous phase processing. We therefore set the secondary
137 production rate $\dot{N}_{\text{PROD}} = 0$.

138 The precipitation sink term \dot{N}_P depends upon the precipitation rate at cloud base
139 P_{CB} . We use a formulation that accounts for losses from the collection of cloud droplets
140 by precipitation drops in the cloud via accretion [Wood 2006]. This gives $\dot{N}_P = K N P_{CB}$
141 h/z_i , where $K = 2.25 \text{ m}^2 \text{ kg}^{-1}$ is a constant that depends upon the collection efficiency of
142 cloud droplets by drizzle drops [Wood 2006], and h is the cloud thickness.

143 Dry deposition of CCN to the ocean is estimated using deposition velocity
144 parameterization [Georgi 1988] for accumulation mode particles (0.05-1 μm diameter)
145 that make up the bulk of the CCN in the MBL. Rates are in the range 0.001-0.01 cm s^{-1}
146 with the higher values occurring at higher wind speeds and for the larger of the
147 accumulation mode particles. Given these deposition velocities, we find that dry
148 deposition constitutes only a very weak sink for CCN under most circumstances, with
149 loss rates unlikely to exceed $2 \text{ cm}^{-3} \text{ d}^{-1}$ for most values of N observed, and most wind
150 speeds, over the oceans. We are therefore justified in setting $\dot{N}_{\text{DRY}} = 0$ in the budget
151 model.

152 We also set $\dot{N}_{\text{ADV}} = 0$ to avoid the complication of calculating spatial gradients
153 and to facilitate the interpretation of the key physical processes controlling N_d . It is
154 possible to estimate the magnitude of the advection term \dot{N}_{ADV} using the observed cloud
155 droplet concentration from satellite [George and Wood 2010]. Over the remote oceans
156 the magnitude of \dot{N}_{ADV} is generally $10 \text{ cm}^{-3} \text{ d}^{-1}$ or less, while higher values can be found
157 in near-coastal regions. Including advection in the steady state model introduces an
158 additional level of complexity since it involves taking spatial gradients. To preserve
159 simplicity we omit it from the model. Since precipitation is the dominant control on N_d in

160 the model, the geographic pattern of the advection term largely follows the spatial
161 gradient in precipitation rate.

162 In this study we examine the time-mean CCN budget by setting $\dot{N}=0$ in
163 equation (1) inserting the expressions discussed above for the various terms, and
164 rearranging to obtain an expression for the steady-state value N_{eq} of the CCN
165 concentration in the MBL as

166

$$N_{eq} = \frac{\left(N_{FT} + \frac{F(\sigma)U_{10}^{3.41}}{Dz_i} \right)}{\left(1 + \frac{hKP_{CB}}{Dz_i} \right)} \quad [2]$$

167 Here we have also assumed that entrainment is in balance with the large-scale
168 subsidence rate, so that $w_e = Dz_i$, where D is the large scale divergence, appropriately
169 assumed to be constant with height over the depth of the MBL [Wood *et al.* 2009]. In
170 practice, the entrainment rate exceeds the subsidence rate by 10-40% over the subtropical
171 stratocumulus regions [Wood and Bretherton 2004], but estimating its precise value is
172 itself a major challenge [Stevens 2002] and is not attempted here.

173

174 3. Model constraints and inputs

175 The CCN budget model (equations 1 and 2) implicitly assumes that the cloud
176 droplet concentration N_d and the CCN concentration are one and the same. This is
177 reasonable provided that (a) we choose an appropriate peak supersaturation σ in order to
178 convert the aerosol sources, which are provided as a function of aerosol size, into

179 tendencies of CCN; (b) the cloud droplet concentration throughout the cloud is equal to
180 that determined by the aerosol activation process.

181

182 3.1 Supersaturation

183 For marine stratocumulus clouds, observations suggest that values of σ in the
184 range 0.1-0.8% are typical with mean values around 0.3% consistent with observations
185 [*Martin et al.* 1994, *Snider et al.* 2003]. Here we assume a constant value of $\sigma=0.3\%$ for
186 all calculations. The peak supersaturation $\sigma = 0.3\%$ assumed in the model is assumed to
187 be constant everywhere. Understanding how this changes systematically over the remote
188 oceans is complex, as it depends upon variations in the strength of the turbulent updrafts
189 and upon the size distribution of the aerosol being activated [*Martin et al.* 1994, *Snider et*
190 *al.* 2003]. We use recent field measurements of mean aerosol size distributions at
191 different distances from the Chilean coast from the VOCALS Regional Experiment
192 (REx) [*Wood et al.* 2011] to estimate the likely systematic geographical variability in σ .
193 In the marine boundary layer, the Hoppel minimum [*Hoppel* 1986] in the size distribution
194 is indicative of the minimum size of aerosols that are activated. Assuming a hygroscopic
195 aerosol, this minimum size is directly related to the supersaturation. A systematic
196 reduction from 0.1 to 0.07 μm in the minimum diameter was observed to occur from
197 70°W to 78°W moving westwards along 20°S away from the Chilean coast [*Kleinmann*
198 *et al.* 2012]. This implies that the mean supersaturation increases by ~75% from the coast
199 offshore. However, the coastal CCN concentrations are high consistent with pollution
200 aerosol impacts [*Allen et al.* 2011], and so these are not representative of the

201 supersaturations we are attempting to represent in the model. No systematic shift in the
202 Hoppel minimum is observed over the cleaner region from 74°W to 78°W, suggesting
203 that systematic changes in supersaturation over the remoter regions may be considerably
204 smaller than occur near the coasts.

205 3.2 Free-tropospheric CCN

206 We constrain N_{FT} based on aerosol measurements from field data taken in the
207 Southern and Northern Hemisphere remote subtropical FT. Two data sources are used:

208 (a) aircraft measurements of FT CCN from a thermal diffusion CCN counter taken on
209 the NSF/NCAR C-130 aircraft during the VOCALS-REx. Details of the
210 instrument are provided in *Snider et al. (2006)*, and flight plans are described in
211 *Wood et al. [2011]* A total of 638 CCN measurements were made in the FT (61,
212 117, 147 and 313 at supersaturations of approximately 0.1, 0.25, 0.5 and 0.9%
213 respectively) along the 20°S latitude line from the coast (70°W) to approximately
214 1500 km offshore (85°W) over the month-long campaign. CCN concentrations
215 are corrected to an assumed mean MBL pressure of 925 hPa, which is the typical
216 pressure at which clouds form in the region [*Bretherton et al. 2010*];

217 (b) a composite time-mean size distribution measured over 17 days with large scale
218 subsidence at a remote FT station on Mauna Loa during July 1992. These data are
219 described in *Weber and McMurry [1996]* and are also corrected to a pressure of
220 925 hPa. We convert the size distribution into a CCN spectrum using a plausible
221 range of aerosol hygroscopicity following the ‘kappa’ parameter approach
222 [*Petters and Kreidenweis 2007*]. A reasonable lower boundary ($\kappa=0.3$) is
223 approximately the lowest value of CCN-based hygroscopicity measured in the

224 northeastern Pacific FT from *Roberts et al.* (2010). The upper boundary ($\kappa=0.98$)
225 is the geometric mean value from *Roberts et al.* (2010). Values of κ significantly
226 higher than unity were inferred from CCN measurements in *Roberts et al.* (2010)
227 but seem implausible given that even the most hygroscopic compounds have κ
228 values of about unity [*Petters and Kreidenweis 2007*]. Ammonium sulfate has a κ
229 value of approximately 0.7 [*Petters and Kreidenweis 2007*].

230

231 Figure 1a compares the measured CCN concentrations in the southeastern Pacific FT
232 with those from Mauna Loa. For the southeastern Pacific, we take data from west of
233 75°W (>500 km off the Chilean coast) that is minimally impacted by coastal pollution
234 [*Allen et al. 2010*]. Time-mean CCN concentrations west of 75°W over the southeastern
235 Pacific Ocean (Fig. 1) are in remarkably good agreement with those derived from the
236 Mauna Loa size distribution measurements. This is perhaps surprising because although
237 the southeastern Pacific and Hawaii are in similar tropical meteorological regimes, one
238 might expect marked differences in the mean size distribution due to the different array of
239 sources and landmasses in the Northern and Southern Hemisphere. The mean
240 distributions are made up of a mixture of different aerosol populations from a number of
241 different sources. Previous studies [*Clarke et al. 1998b, Friedlander 1977, Raes 1995*]
242 suggest that new particle formation from naturally-produced sulfuric acid in the upper
243 troposphere constitutes one major source of clean FT CCN. This aerosol subsequently
244 subsides in the descending branches of large scale atmospheric systems where its
245 distribution is expected to reach a quasi-steady state [*Friedlander 1977, Raes 1995*]. In

246 addition, the remote marine FT also includes air masses transported long distances from
247 continents that likely contain some pollution aerosol.

248 We derive a value of N_{FT} in the range 100-175 cm^{-3} active at $\sigma=0.3\%$ from the
249 mean Mauna Loa size distribution (assuming a plausible range of aerosol hygroscopicity,
250 see caption, Fig. 1a). This is in remarkable agreement with mean values from CCN
251 spectrum measurements made in the southern hemisphere subtropical FT during
252 VOCALS-REx (see Fig. 2a). For the model base case we therefore assume a constant
253 mean concentration $N_{FT} = 125 \text{ cm}^{-3}$ everywhere, which is within 20% of the time-mean
254 values derived from the remote subtropical data in both hemispheres. We also force the
255 model with observed values of FT CCN from VOCALS-REx and conduct additional
256 sensitivity tests, as described in Section 4.

257

258 3.3 Sea surface source

259 For the surface source, we use a size-resolved sea spray generation function
260 [Clarke *et al.* 2006] to estimate the rate of particle generation for particles active at
261 $\sigma=0.3\%$ supersaturation. This is determined numerically by integrating the size-resolved
262 surface source function from the largest particles down to the critical dry diameter for a
263 given supersaturation. This provides a supersaturation-dependent $F(\sigma)$ curve shown in
264 Fig. 1b. For $\sigma=0.3\%$ the critical diameter is 50 nm, which yields a value
265 of $214 \text{ m}^{-3} (\text{m s}^{-1})^{-2.41}$ for $\sigma=0.3\%$ (Fig 1b). For a wind speed of 8 m s^{-1} this yields $\dot{N}_S = 22$
266 $\text{cm}^{-3} \text{ d}^{-1}$ averaged over an MBL that is 1 km deep. The widely-used source function of
267 Monahan *et al.* [1986] has the same wind speed dependence but a rate that is over a
268 factor of two lower (Fig. 1b).

269 To drive SSA production we use daily mean wind speed estimates from the
270 QuikScat satellite and average the production rates up to monthly averages. Including
271 sub-daily timescale variability in wind speed increases the surface production of SSA, but
272 we find from reanalysis data that including 6 hourly estimates increases SSA production
273 by less than 10%.

274

275 3.4 Precipitation sink

276 The model sink term is driven by new precipitation rate estimates from the
277 profiling W-band radar on the CloudSat satellite [*Lebsock and L'Ecuyer* 2011]. The
278 cloud base precipitation rate needed to calculate the coalescence scavenging [*Wood* 2006]
279 is estimated as the maximum value in each radar profile. Mean precipitation rates from
280 low clouds are estimated for $5\times 5^\circ$ gridboxes globally by removing profiles with
281 detectable echoes above the 3 km level. For the regions considered, the results are not
282 strongly sensitive to the choice of this level since the majority of the clouds are situated
283 below 2 km. Gridded precipitation rates are produced on a monthly basis for data from
284 2006-2009.

285 We also use precipitation measurements from VOCALS-REx to compare against
286 those from CloudSat. For this, we use both aircraft in-situ observations from an optical
287 array probe, and aircraft radar measurements from the University of Wyoming Cloud
288 Radar. The VOCALS-REx precipitation dataset is described in *Bretherton et al.* [2010].
289 The majority of the aircraft flights were conducted during the later part of the night (03-
290 09 local time) when precipitation rates are at their diurnal maximum.

291

292 3.5 Boundary layer depth, cloud thickness, wind speed, and surface divergence.

293 Equation (2) indicates that we also need to estimate MBL depth z_i , cloud
294 thickness h , surface wind speed U_{10} , and surface divergence D . MBL depth is estimated
295 using from MODIS cloud top temperature retrievals [Wood and Bretherton 2004]. Cloud
296 thickness h is estimated with an adiabatic assumption [e.g. Albrecht et al. 1990] using
297 MODIS retrievals of cloud liquid water path (LWP). This relationship is $h = (2LWP/\Gamma)^{1/2}$,
298 where Γ is a weak function of temperature and pressure [Albrecht et al. 1990], evaluated
299 as described in Wood and Bretherton [2004]. Both cloud top temperature and LWP are
300 taken from $1 \times 1^\circ$ gridded daily Level 3 MODIS products. The results are not strongly
301 sensitive to these parameters. Wind speed and surface divergence estimates are from the
302 QuikScat satellite [see Wood et al. 2009 for details].

303

304 3.5. Cloud droplet concentration

305 Satellite estimates of cloud droplet concentration are used to compare against
306 model-derived N_{eq} from Equation (2). Model estimates are produced globally on a $1 \times 1^\circ$
307 grid on a month by month basis. Cloud droplet concentration estimates from the
308 Moderate Resolution Imaging Spectroradiometer (MODIS) on the NASA Terra satellite
309 are produced from daily Level 3 data for $1 \times 1^\circ$ boxes using a visible/near infrared
310 approach [Bennartz 2007]. To minimize problems of retrievals in broken clouds, we only
311 include in our averages those daily boxes where the cloud cover from liquid clouds
312 exceeds 0.8. These are then averaged together to provide monthly mean N_d estimates.

313

314 3.6. Model estimates of N_d

315 We use the budget model Equation (2) to predict monthly mean values of N_d by
316 forcing with monthly-mean values of the input variables discussed above. Annual means
317 are then derived from the monthly means only for those months with (a) mean subsidence
318 (positive mean surface divergence); (b) mean boundary layer depth lower than 4 km; and
319 (c) with mean liquid cloud fractions exceeding 0.3. Annual mean data are only analyzed
320 for those locations where at least four months pass the acceptance criteria.

321

322 4. Results

323 4.1 Model assessment over the southeastern Pacific

324 To assess the quality of the model, we use recent field measurements from VOCALS-
325 REx that extensively sampled the lower troposphere over the tropical southeastern Pacific
326 Ocean [Wood *et al.* 2011]. The measurements focused upon characterizing the largest
327 semi-permanent subtropical sheet of stratocumulus on Earth that extends westward from
328 the Chilean and Peruvian coasts. Extensive survey sampling was carried out along 20°S
329 from the Chilean coast at 70°W to ~1400 km offshore at 85°W using a combination of
330 different research aircraft. Measurements were made in both the MBL and the lower FT.
331 Figure 2a shows observations of the two key inputs to the budget model, namely the
332 time-mean FT CCN concentration (section 3.2) and the precipitation rate close to the
333 cloud base (section 3.4). CCN concentrations in the FT fall off sharply within 500 km of
334 the coast but, despite considerable day to day variability, the time-mean campaign values

335 remain relatively constant for over 1000 km out to 85°W (Fig 2a). Several lines of
336 evidence point to anthropogenic pollution being responsible for the high values very
337 close to the coast [Allen *et al.* 2011, Yang *et al.* 2011, Saide *et al.* 2012].

338 The cloud base precipitation rate increases markedly with distance from the
339 Chilean coast, from an essentially nonprecipitating state with $<0.1 \text{ mm d}^{-1}$ near the coast
340 to $>1 \text{ mm d}^{-1}$ at 85°W (Fig. 2a). Rates of a few tenths of a mm d^{-1} are sufficient to drive
341 significant coalescence scavenging of CCN [Feingold *et al.* 1996, Wood 2006]. This
342 gradient in precipitation is driven to a significant extent by thickening clouds and a
343 deeper boundary layer to the west [Bretherton *et al.* 2010], but is likely also modulated
344 by aerosol in the MBL [Terai *et al.* 2012]. The cause of the precipitation is not the focus
345 of this study. Precipitation production in marine stratocumulus maximizes at night [Leon
346 *et al.* 2008], and is typically heaviest in the early morning hours when the clouds are at
347 their thickest [Bretherton *et al.* 2004, Wood *et al.*, 2002]. This is apparent in the
348 observations where precipitation rates are lowest at 1:30pm (CloudSat, daytime overpass,
349 Fig. 2a), take intermediate values at 1:30am (CloudSat nighttime overpass) and are
350 largest during 03-09 am (VOCALS field data). Since the timescale for CCN removal due
351 to precipitation is typically at least 1 day given these precipitation rates [Wood 2006], we
352 use a daily mean estimate as the mean of the two CloudSat overpasses to drive the model.

353 The model, when forced with CloudSat observed mean precipitation rates and
354 observed FT CCN, captures the observed increase in N_d as the coast is approached (Fig.
355 2b) with remarkable fidelity given the model's simplicity. The model predicts a factor of
356 two increase in N_d from 90°W to 75°W even when the model is forced with a fixed FT
357 CCN concentration of 125 cm^{-3} (consistent with mean values over the remote region

358 away from the coast). With fixed CCN, however, the model is unable to reproduce the
359 highest concentrations within 500 km of the coast (Fig. 2b). However, these high values
360 are obtained when the model is forced with the observed longitudinally-varying FT CCN
361 concentration increase that includes the near-coastal enhancement due to pollution
362 sources [Allen *et al.* 2011] (Fig 2b). The general behavior of decreasing N_d westward
363 from 75-90°W is caused by increasing precipitation scavenging (Fig. 2a), which can be
364 seen by comparing the model estimates with fixed FT CCN and either no precipitation or
365 precipitation fixed at a constant value of 1 mm d⁻¹ (Fig 2b). A critical finding here is that
366 a precipitation rate of as little as 1 mm d⁻¹ is sufficient to drive down N_d by a factor of
367 three over the remote ocean, which further serves to emphasize how important
368 precipitation from low clouds is in controlling mean cloud droplet concentrations over the
369 remote ocean.

370 Primary production from sea-spray constitutes a weaker, but nevertheless
371 significant, source than entrainment from the FT consistent with a previous study with the
372 same parameterization [Clarke *et al.* 2006]. However, it is important to note that we are
373 using a source function that is one of the more prolific available [de Leeuw *et al.* 2011],
374 although experimentation with different primary production parameterizations only
375 changes the modeled N_d values by less than 20% (Fig. 3). The CCN concentration from
376 sea-salt (the difference between the solid and the dotted line in Fig. 3) ranges from <10
377 cm⁻³ close to the Chilean coast, where wind speeds are low, to around 40 cm⁻³ further
378 afield, where wind speeds are higher. These findings are consistent with preliminary
379 measurements of sea-salt particles from aircraft during VOCALS-REx (Anthony Clarke,
380 personal communication). The choice of sea-spray parameterization does not affect our

381 conclusions regarding the importance of precipitation at driving the offshore gradient in
382 cloud droplet concentration over the southeastern Pacific.

383

384 4.2 Application to marine low cloud regions globally

385 Given the ability of the budget model with fixed FT CCN concentrations to
386 reproduce with some skill the gradient in N_d over the remote southeastern Pacific Ocean
387 (more than 500 km from the coast), we apply the model more generally to regions of
388 extensive marine low cloud under conditions of large scale subsidence (Fig. 4). Aerosol
389 concentrations in the FT vary significantly both regionally and in response to variations
390 in natural and anthropogenic sources [Allen *et al.* 2011]. Because there are no global
391 observational constraints on the time-mean FT CCN, we draw on the consistency
392 between mean FT CCN spectra at Mauna Loa and over the southeastern Pacific Ocean
393 (Fig. 1 and section 3.2) and fix $N_{FT} = 125 \text{ cm}^{-3}$ everywhere for the base case. The satellite
394 observations show that N_d values in excess of 150 cm^{-3} tend to be located near the
395 continental coastlines (e.g. California, Chile/Peru, Europe), with values reducing toward
396 the remote oceans, where they are as low as $30\text{-}60 \text{ cm}^{-3}$ (Fig. 4a). The base case model is
397 able to reproduce well the mean values (Table 1) and geographical variability in N_d (Fig.
398 4) for low cloud regions, especially for the remote subtropical/tropical regions 35°S -
399 35°N . The model underestimates N_d close to coastlines (Table 1) consistent with a lack of
400 continental sources. The model also captures the low values ($<60 \text{ cm}^{-3}$) over the remote
401 North Pacific and Atlantic and the Southern Ocean north of 45°S . Removing the
402 precipitation sink increases mean N_d in the model by a factor of 2-3 over the remote
403 oceans, but only 15% in the near-coastal regions where precipitation rates are very low

404 (Table 1). A doubling of N_{FT} from 80 to 160 cm^{-3} leads to a 50-70% increase in N_d
405 (Table 1) because FT CCN is partly buffered by surface sources.

406 The skill of the base case model in predicting N_d variability is remarkable given
407 that there is no variation whatsoever in the FT CCN source in Fig. 4. To examine the key
408 factors controlling the geographical variability of N_d in the model, we conduct additional
409 model sensitivity experiments. In each experiment, only one (or two) of the variables in
410 equation (2) is allowed to vary. All other variables are fixed by setting them to their
411 respective mean values over time and space (Table 2). It is clear that precipitation
412 variability is required in order to produce the strong correlation with observations seen in
413 the base case ($r = 0.65$). No other variable can alone explain more than 15% of the
414 observed geographical variance in N_d . Divergence and cloud thickness variations also
415 lead to model fields with significant positive correlations ($r = 0.21$ and 0.37 respectively),
416 but the geographical variability in N_d driven by these variables is far too weak to explain
417 the observed variability (Table 2). The correlations are positive because divergence and
418 cloud thickness correlate quite well with precipitation itself. We find that wind speed
419 variability alone explains an insignificant amount of the model N_d variability ($r = -0.06$),
420 from which we conclude that variability in SSA is not a significant contributor to the
421 observed geographical variability in N_d . This is especially true in the subtropics and
422 tropics where wind speeds are relatively modest and the FT source is greater than the
423 surface source (Fig. 5). However, SSA does contribute to the mean N_d (Table 1) despite
424 not substantially impacting its geographical variability.

425 Overall, these results lead us to argue that a large fraction of the observed
426 geographical variability in cloud droplet concentration in extensive marine low clouds

427 over the remote oceans is driven by precipitation losses rather than aerosol source
428 variability. This is further demonstrated by noting the striking similarities between the
429 maps of the mean observed precipitation rates from low clouds (Fig. 6) and the observed
430 N_d field (Fig. 4a).

431 Frequency distributions of monthly mean N_d (Fig. 7) show that the base case
432 model ($N_{FT} = 125 \text{ cm}^{-3}$) can represent satellite-observed N_d variability well. As we might
433 expect from Fig. 2, the model is unable to capture the very highest concentrations
434 observed by the satellite ($N_d > 200 \text{ cm}^{-3}$) that are mostly regions within a few hundred
435 kilometers of coastlines. This is because neither the advection of continentally-influenced
436 MBL air nor elevated near-coastal FT concentrations (e.g., Fig. 2) are considered in the
437 model. When the model is forced by removing either FT CCN or SSA production, the
438 model is unable to represent the distribution of observed N_d and underestimates the mean
439 N_d (Fig. 7 and see also Table 1). This further emphasizes that the surface and FT are both
440 important contributors to the cloud droplet concentration over the remote oceans
441 [*Katoshevski et al. 1999, Capaldo et al. 1999, Clarke et al. 2006*]. The shape of the
442 model N_d distribution is relatively insensitive to plausible variations in the assumed mean
443 FT CCN concentration (Fig. 8).

444

445 **5. Implications and conclusions**

446 Our results have a number of important implications. First, if CCN and cloud
447 droplet concentration variability over much of the global ocean are determined by
448 precipitation variability rather than aerosol source variability, this calls into question the
449 interpretation of correlative studies [*Quaas et al. 2008, Jones et al. 2009*] linking cloud

450 properties to aerosol properties as providing useful information on anthropogenic aerosol
451 indirect effects on climate. It also suggests that the notion of there being a ‘background’
452 aerosol concentration in the unperturbed marine boundary layer may not be a useful one
453 because MBL CCN concentrations are strongly modulated by precipitation processes that
454 vary strongly both geographically and temporally. The FT CCN over remote oceanic
455 regions is known to reflect a complex mixture of different sources, some of which are
456 natural and some anthropogenic [*Clarke and Kapustin 2010*]. Our finding that a constant
457 time-mean FT CCN supply is sufficient to explain a significant fraction of the time-mean
458 gradients in the cloud droplet concentration over the remote oceans should not therefore
459 be interpreted as indicating that the remote oceanic regions are devoid of anthropogenic
460 influence. The increasing concentrations observed within about 500 km of continents
461 most likely reflects a lower tropospheric pathway for the transport of continental aerosols
462 to the MBL, whereas the more remote anthropogenic contributions are associated with
463 aerosol or precursors lofted higher into the troposphere that can then be transported long
464 distances before subsidence carries them into the MBL.

465 Here we have shown that MBL cloud droplet concentrations are impacted by
466 precipitation generated by the clouds themselves, but we note that an increasing body of
467 evidence shows that precipitation in low clouds typically decreases with cloud droplet
468 concentration [*Stevens and Feingold 2009*]. There is then the potential for a significant
469 positive feedback whereby modest increases in CCN reduce the precipitation sink,
470 amplifying the initial perturbation. Although we do not claim evidence for bistability in
471 the system [*Baker and Charlson 1990*], our results do suggest that pollution-driven CCN
472 increases may be amplified by precipitation suppression and that this warrants further

473 exploration with more sophisticated models. There is modeling evidence for this in the
474 recent literature [*Yang et al. 2012*].

475 Finally, we note that climate models tend to impose arbitrary fixed limits on cloud
476 droplet concentration minima [*Quaas et al. 2009*] suggesting deficiencies in modeling the
477 processes responsible for low concentrations over the remote ocean. Across models, there
478 is a significant correlation between this fixed lower limit on N_d and the strength of the
479 aerosol indirect effect [*Quaas et al. 2009*]. A closer focus on the role of precipitation is
480 therefore needed to better understand whether climate models are able to produce light
481 precipitation in the marine boundary layer, and whether the models are impacted
482 appropriately by it. This study helps to highlight that we now have the satellite
483 measurements of light precipitation and cloud microphysical properties to begin to
484 explore this critical control on cloud microphysical properties. Further work could
485 include an examination of the effects of seasonal precipitation variability on the seasonal
486 cycle of cloud droplet concentration, as suggested by previous work [*Liu 2010*].

487

488 **Acknowledgements:**

489 The authors would like to thank the staff and crew of the NSF/NCAR C-130 aircraft
490 whose dedication resulted in the in-situ observational VOCALS Regional Experiment
491 dataset. The CloudSat data were distributed by the CloudSat Data Processing Center at
492 Colorado State University. MODIS data were obtained from the NASA Goddard Land
493 Processes data archive. QuikScat data were produced by Remote Sensing Systems and
494 sponsored by the NASA Ocean Vector Winds Science Team. This work was supported
495 by NASA award numbers NNX10AN78G and NNX10AM29G and NSF award numbers
496 ATM-0745702 and ATM-0745368.

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500 **References**

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684
685

686 **Figure Captions**

687

688 **Figure 1: (a) Free-tropospheric CCN spectra from the southeastern Pacific and**
689 **Hawaii.** Observations from the southeastern Pacific are from CCN spectra taken in the
690 remote FT west of 75°W using the NSF/NCAR C-130 aircraft in VOCALS-REx [Wood
691 *et al.* 2011], corrected to an assumed mean MBL pressure of 925 hPa. Box-whisker plots
692 show the 10th, 25th, 50th, 75th, and 90th percentile concentrations for four supersaturations.
693 Shaded region shows a plausible range of CCN concentration estimated using the
694 composite size distribution for subsiding FT air measured on Mauna Loa in Hawaii
695 [Weber and McMurry 1996], corrected to an assumed mean MBL pressure of 925 hPa,
696 with the spread representing a plausible range of hygroscopicity κ parameters [Petters
697 and Kreidenweis 2007] for clean FT air (see text); **(b) Sea-spray source functions $F(\sigma)$**
698 as a function of supersaturation, for the Clarke *et al.* (2006) parameterization used in this
699 study, and from Monahan *et al.* (1986) for comparison.

700

701 **Figure 2: Model inputs and results from southeastern Pacific stratocumulus region**
702 **from 70-90°W along 20°S.** (a) Free-tropospheric (FT) aerosol concentrations (left axis)
703 showing range of mean CCN concentrations corresponding to supersaturations relevant
704 for cloud formation (gray shading), and total non-volatile particle concentration (open
705 circles). Green bar shows estimated CCN for 0.2-0.5% supersaturation from
706 measurements of FT aerosol size distributions during subsiding conditions on Mauna
707 Loa, Hawaii, see the Methods section; cloud base precipitation rates (right axis) estimated
708 from CloudSat satellite (mean for October/November 2006-2009 between 22°S and 18°S,
709 red shading showing 1:30am and 1:30pm local time overpasses) and from the VOCALS-
710 REx field experiment (black and blue squares from aircraft radar and in-situ precipitation
711 probes respectively, in the latitude range 18-22°S); **(b)** observed (solid circles: aircraft
712 during VOCALS [Bretherton *et al.* 2010, Wood *et al.* 2011], diamonds: satellite estimates
713 from MODIS, 18-22°S) and modeled mean cloud droplet concentration N_d for different
714 model scenarios as denoted in legend and discussed in the text.

715

716 **Figure 3: Effects of different assumptions regarding primary production of sea-salt**
717 **in the model.** Same as Fig. 2b, but showing sensitivity to sea-salt aerosol
718 parameterization used. Observed (solid circles: aircraft during VOCALS [taken from
719 Bretherton *et al.* 2010], diamonds: satellite estimates from MODIS, 18-22°S) and
720 modeled mean cloud droplet concentration N_d . Solid line: standard model set up with
721 constant FT CCN. Dashed line: Monahan sea-salt parameterization [Monahan *et al.*
722 1986] in place of Clarke *et al.* [2006]. Dotted line: No primary production of sea-salt at
723 all.

724 **Figure 4: Cloud droplet concentrations in regions of extensive marine low clouds**
725 **observed by satellite and from the budget model.** Annual mean cloud droplet
726 concentration N_d for extensive marine low clouds under conditions of large scale
727 subsidence, (a) from MODIS (see Methods section); (b) from the CCN budget model for
728 the same regions.

729

730 **Figure 5: Mean precipitation rate at cloud base from low clouds (cloud top height**
731 **$z_{top} < 3$ km) estimated with spaceborne radar measurements from CloudSat [Lebsock**

732 and L'Ecuyer 2011]. Data are screened to display regions of extensive marine low clouds
733 under conditions of mean subsidence as in Fig. 4.

734

735 **Figure 6:** Ratio of CCN flux from surface to that from entrainment from the free
736 troposphere (FT) in the model. The FT CCN concentration is set to 125 cm^{-3} and the
737 surface source depends upon daily wind speed. In the subtropics and tropics, the majority
738 of the CCN originate from the FT, but in the midlatitudes where winds are stronger, the
739 surface source can exceed that from the FT.

740

741 **Figure 7: Frequency distributions of observed and modeled monthly mean cloud**
742 **droplet concentration.** Only months that meet the criteria needed to contribute to the
743 means shown in Fig. 4 (regions with extensive low clouds under divergent conditions) are
744 shown. Shown here are the base version of the model (i.e. that used to construct Fig. 4,
745 solid black), together with estimates with no FT contribution to CCN (dashed), and no
746 sea salt contribution (dotted). Error bars show the 95% confidence interval in the
747 frequency estimates for the observations and for the model base case due to sampling
748 limitations.

749

750 **Figure 8:** Sensitivity to assumed free tropospheric CCN concentration N_{FT} of the
751 frequency distributions modeled monthly mean cloud droplet concentration. Only months
752 that meet the criteria needed to contribute to the means shown in Fig. 4 (regions with
753 extensive low clouds under divergent conditions) are shown. Shown here are the base
754 version of the model (i.e. that used to construct Fig. 4, solid black), together with model
755 estimates with low ($N_{\text{FT}} = 80 \text{ cm}^{-3}$) and high ($N_{\text{FT}} = 160 \text{ cm}^{-3}$) estimates of the FT
756 contribution to CCN (dashed blue and orange respectively). Error bars show the 95%
757 confidence interval in the frequency estimates for the observations and for the model base
758 case due to sampling limitations.

759

760

761 **TABLES**

762

763 Table 1: The effects on the mean cloud droplet concentration in various geographical
 764 regions of changing the primary source and sink terms
 765

Observed or model estimate	Mean cloud droplet concentration [cm^{-3}]		
	35°S-35°N	35°S-35°N within 300 km of coastlines	60°S-60°N
Observations (MODIS)	88	152	74
Model, base case ($N_{\text{FT}} = 125 \text{ cm}^{-3}$)	88	129	74
Model, $N_{\text{FT}} = 0$	20	21	25
Model, no SSA	68	109	46
Model, no precipitation	169	150	245
Model, $N_{\text{FT}} = 80 \text{ cm}^{-3}$	63	90	55
Model, $N_{\text{FT}} = 160 \text{ cm}^{-3}$	107	154	85

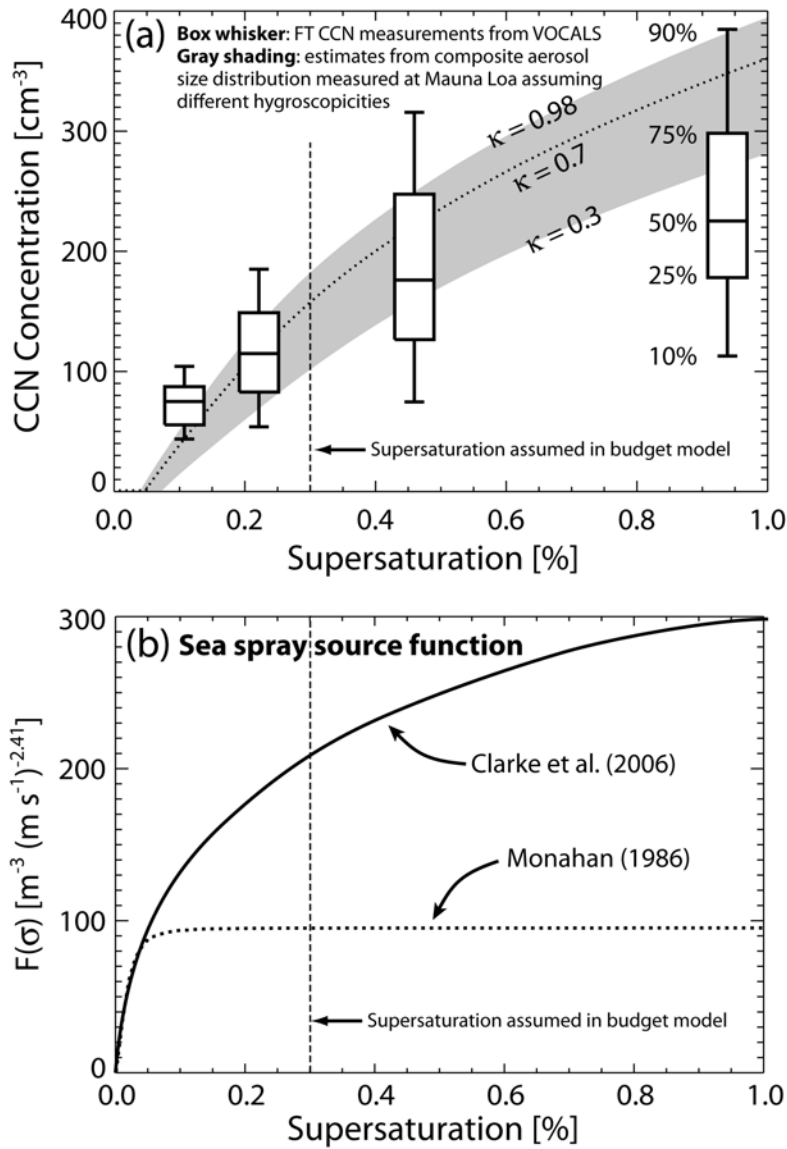
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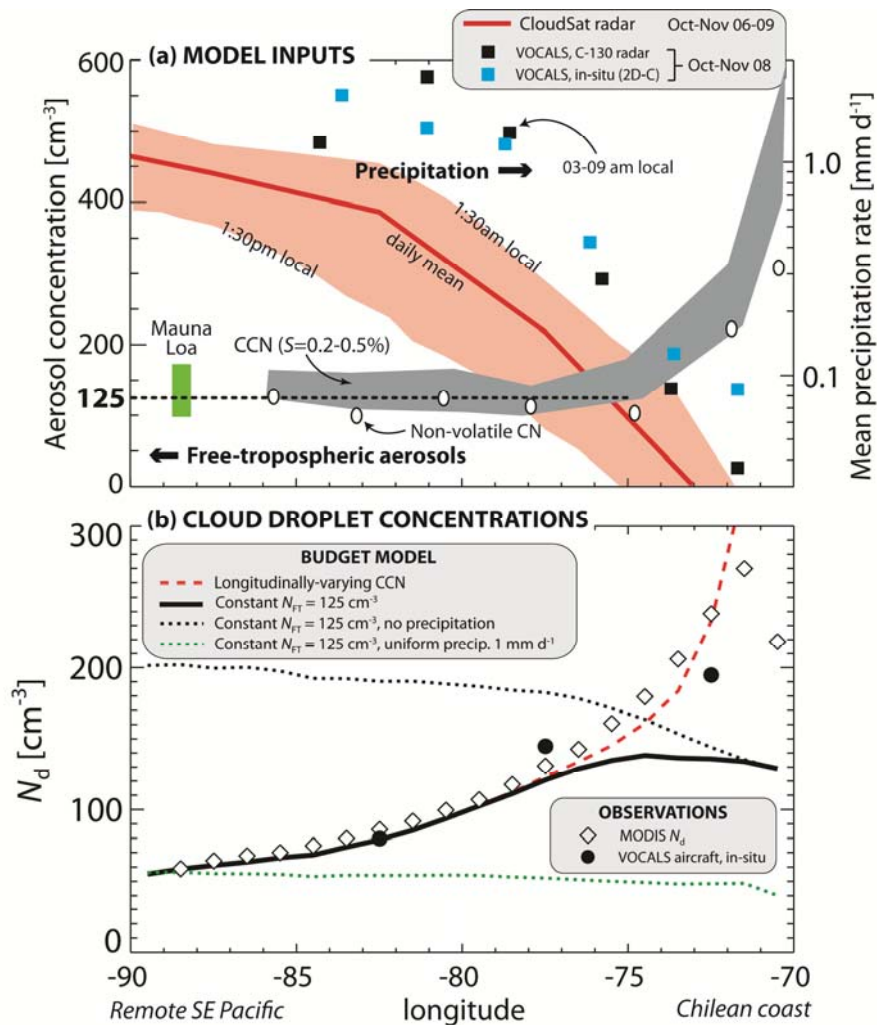
768 Table 2: Geographical variability of cloud droplet concentration for various model
 769 configurations. Results show correlation coefficients between annual mean MODIS
 770 observed and model estimates of annual mean N_d and the ratio of the model and observed
 771 standard deviations ($\sigma_{\text{model}}/\sigma_{\text{obs}}$). Correlations not significant at the 2σ level are italicized.
 772 All results are for the tropics and subtropics (35°S-35°N).
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Model configuration	r (obs, model)	$\sigma_{\text{model}}/\sigma_{\text{obs}}$
Base case	0.65	0.92
<i>P_{CB} variability only</i>	<i>0.77</i>	1.01
<i>P_{CB} and U_{10} variability only</i>	<i>0.70</i>	0.97
<i>U_{10} variability only</i>	<i>-0.06</i>	0.37
<i>D variability only</i>	<i>0.21</i>	0.18
<i>h variability only</i>	<i>0.37</i>	0.13
<i>z_i variability only</i>	<i>-0.32</i>	0.10

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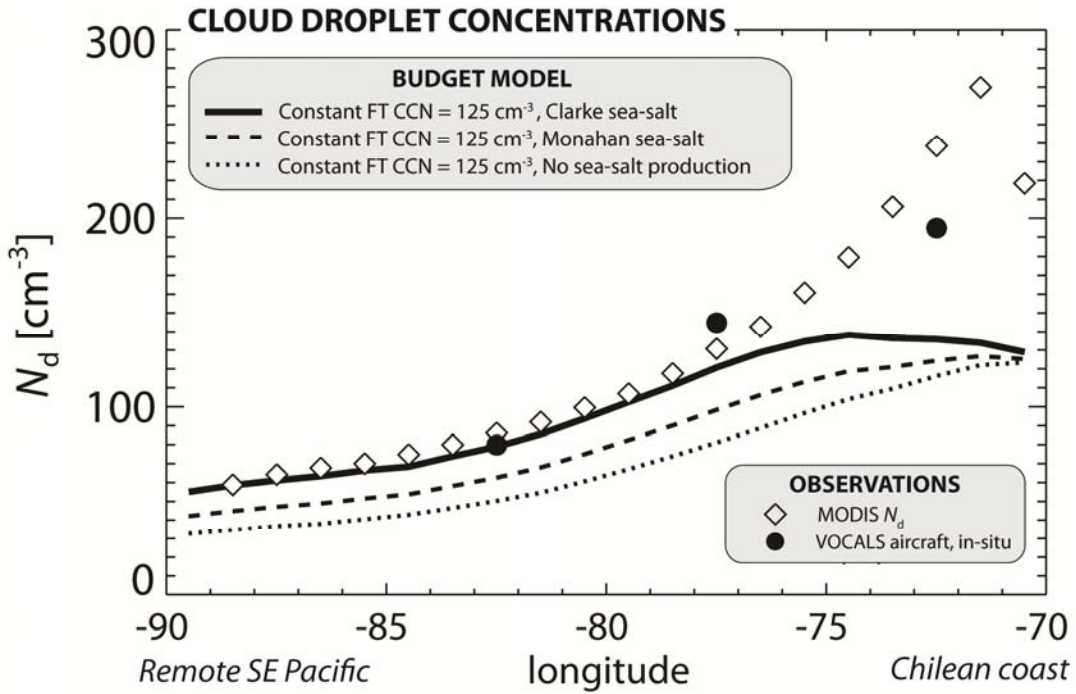
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 779 **Figure 1: (a) Free-tropospheric CCN spectra from the southeastern Pacific and**
 780 **Hawaii.** Observations from the southeastern Pacific are from CCN spectra taken in the
 781 remote FT west of 75°W using the NSF/NCAR C-130 aircraft in VOCALS-REx [Wood
 782 *et al.* 2011], corrected to an assumed mean MBL pressure of 925 hPa. Box-whisker plots
 783 show the 10th, 25th, 50th, 75th, and 90th percentile concentrations for four supersaturations.
 784 Shaded region shows a plausible range of CCN concentration estimated using the
 785 composite size distribution for subsiding FT air measured on Mauna Loa in Hawaii
 786 [Weber and McMurry 1996], corrected to an assumed mean MBL pressure of 925 hPa,
 787 with the spread representing a plausible range of hygroscopicity κ parameters [Petters
 788 and Kreidenweis 2007] for clean FT air (see text); (b) **Sea-spray source functions $F(\sigma)$**
 789 as a function of supersaturation, for the Clarke et al. (2006) parameterization used in this
 790 study, and from Monahan et al. (1986) for comparison.



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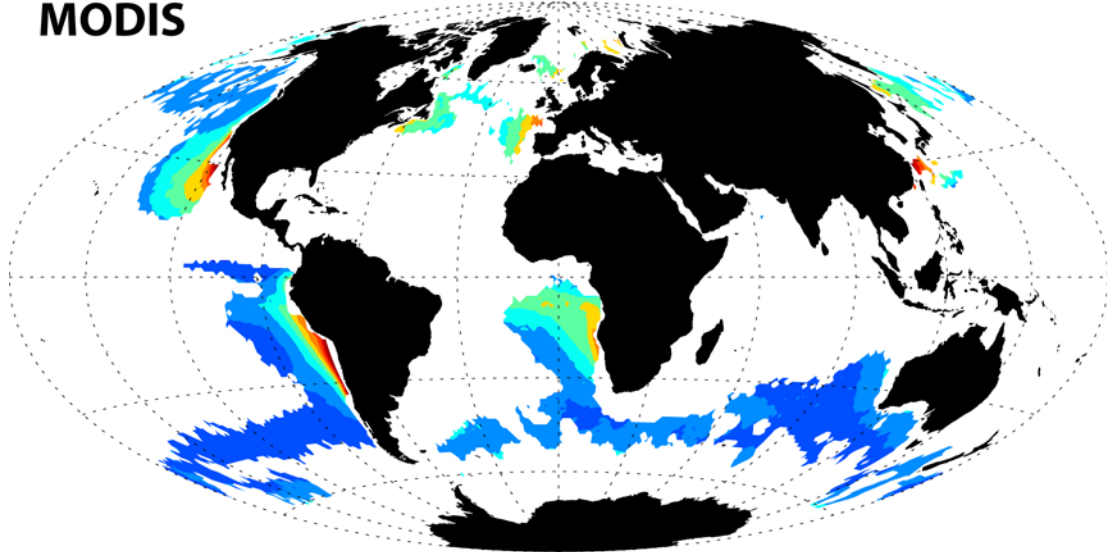
794 **Figure 2: Model inputs and results from southeastern Pacific stratocumulus region**
 795 **from 70-90°W along 20°S. (a)** Free-tropospheric (FT) aerosol concentrations (left axis)
 796 showing range of mean CCN concentrations corresponding to supersaturations relevant
 797 for cloud formation (gray shading), and total non-volatile particle concentration (open
 798 circles). Green bar shows estimated CCN for 0.2-0.5% supersaturation from
 799 measurements of FT aerosol size distributions during subsiding conditions on Mauna
 800 Loa, Hawaii, see the Methods section; cloud base precipitation rates (right axis) estimated
 801 from CloudSat satellite (mean for October/November 2006-2009 between 22°S and 18°S,
 802 red shading showing 1:30am and 1:30pm local time overpasses) and from the VOCALS-
 803 REx field experiment (black and blue squares from aircraft radar and in-situ precipitation
 804 probes respectively, in the latitude range 18-22°S); **(b)** observed (solid circles: aircraft
 805 during VOCALS [Bretherton *et al.* 2010, Wood *et al.* 2011], diamonds: satellite estimates
 806 from MODIS, 18-22°S) and modeled mean cloud droplet concentration N_d for different
 807 model scenarios as denoted in legend and discussed in the text.
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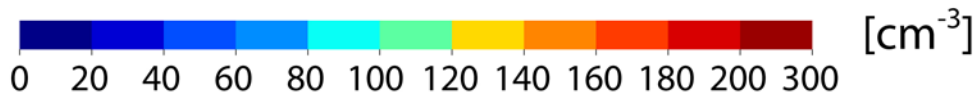
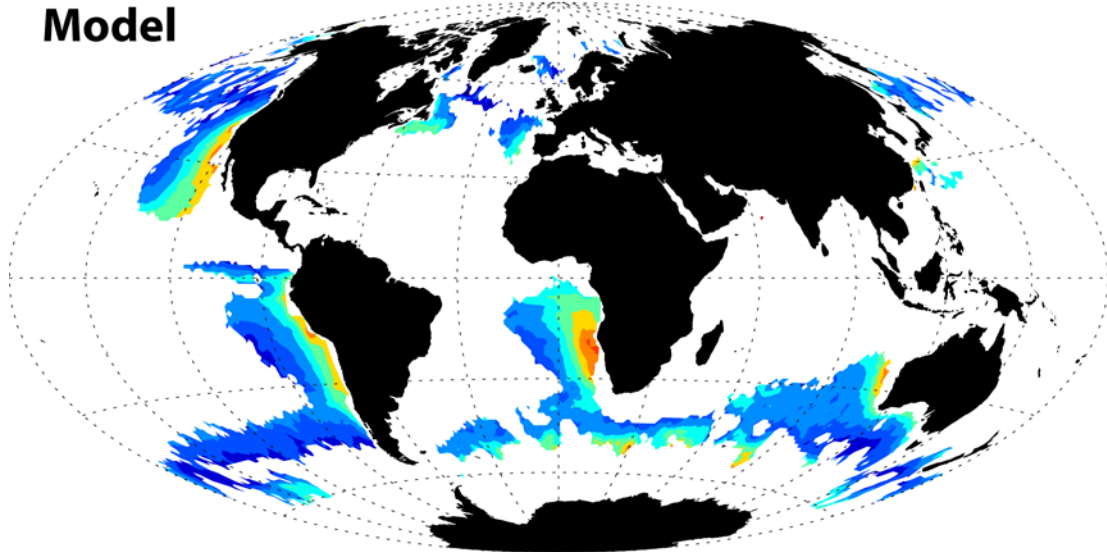
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Figure 3: Effects of different assumptions regarding primary production of sea-salt in the model. Same as Fig. 2b, but showing sensitivity to sea-salt aerosol parameterization used. Observed (solid circles: aircraft during VOCALS [taken from *Bretherton et al.* 2010], diamonds: satellite estimates from MODIS, 18-22°S) and modeled mean cloud droplet concentration N_d . Solid line: standard model set up with constant FT CCN. Dashed line: Monahan sea-salt parameterization [*Monahan et al.* 1986] in place of *Clarke et al.* [2006]. Dotted line: No primary production of sea-salt at all.

MODIS

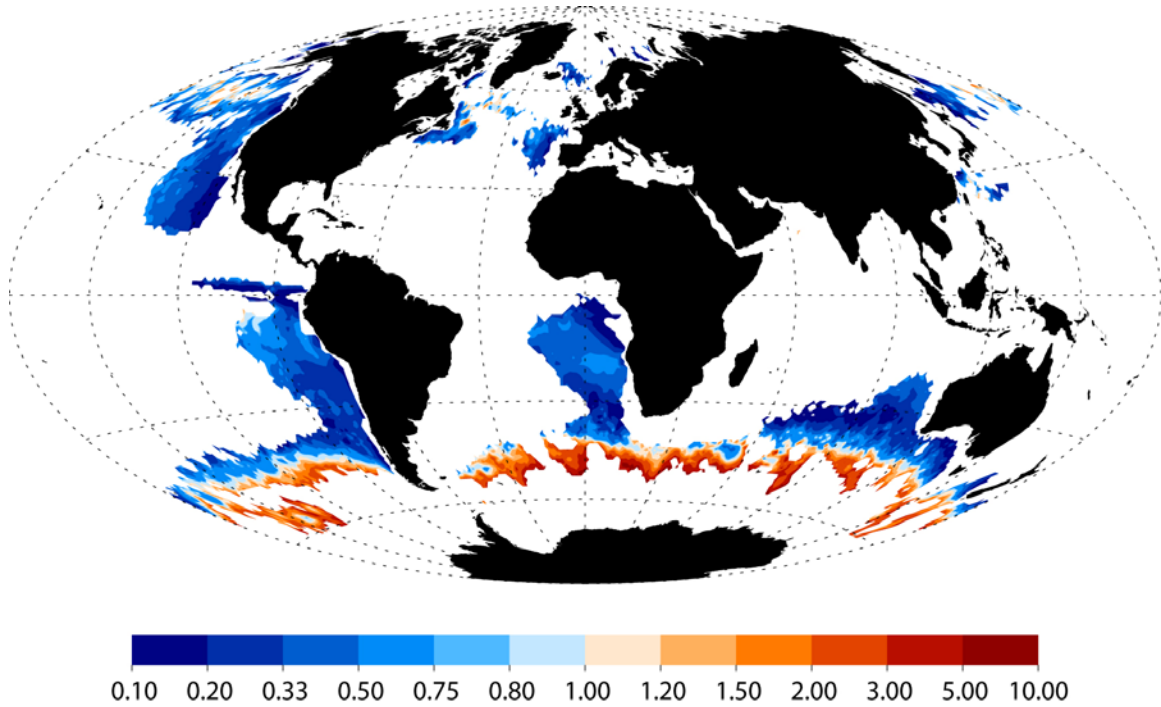


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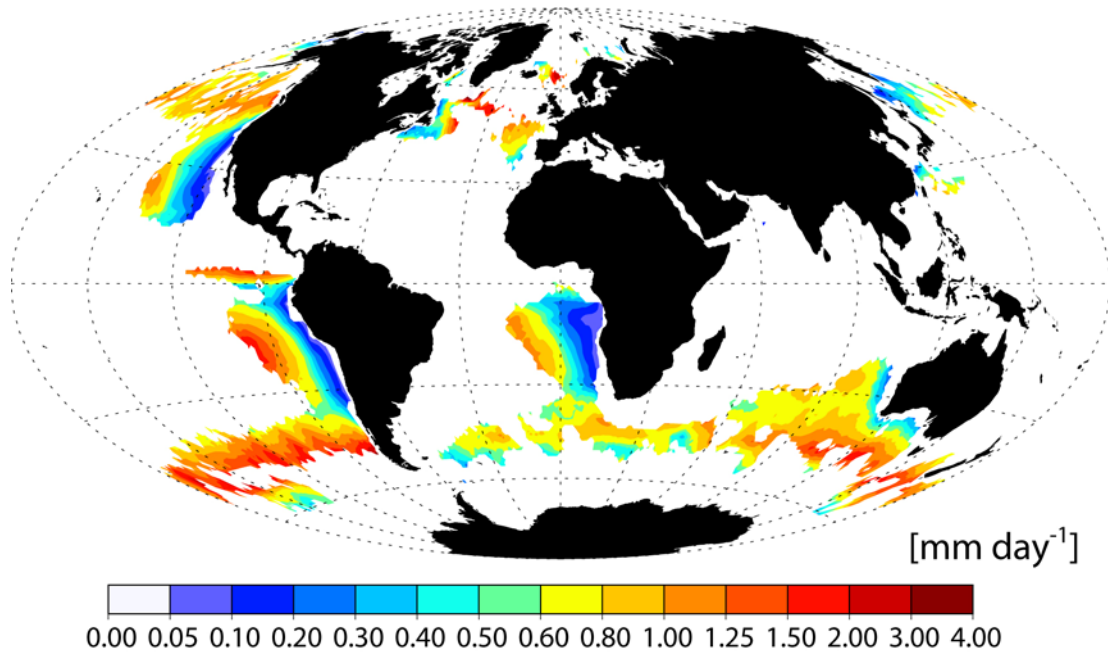
Figure 4: Cloud droplet concentrations in regions of extensive marine low clouds observed by satellite and from the budget model. Annual mean cloud droplet concentration N_d for extensive marine low clouds under conditions of large scale subsidence, (a) from MODIS (see section 3.5); (b) from the steady-state CCN budget model for the same regions.



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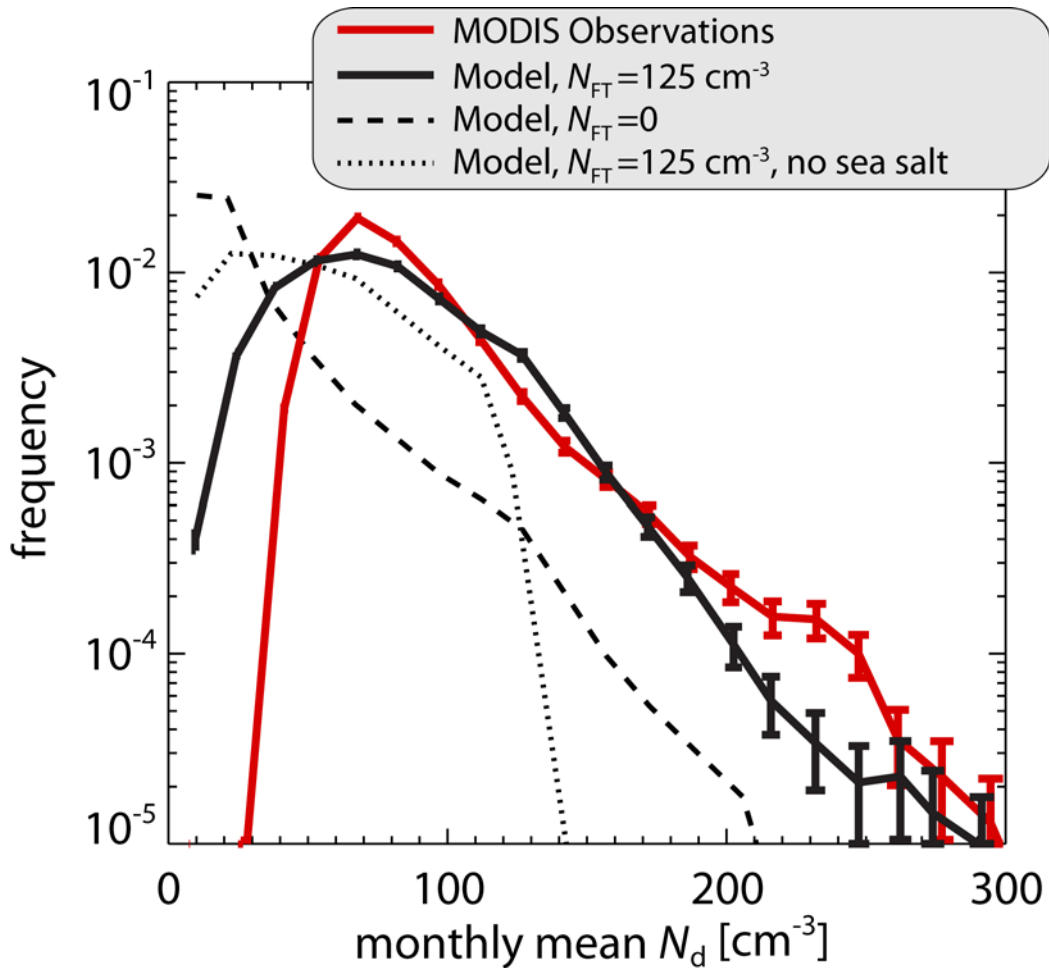
Figure 5: Ratio of CCN flux from surface to that from entrainment from the free troposphere (FT) in the model. The FT CCN concentration is set to 125 cm^{-3} and the surface source depends upon surface wind speed. In the subtropics and tropics, the majority of the CCN originate from the FT, but in the midlatitudes where winds are stronger, the surface source can exceed that from the FT.

CloudSat precipitation rate ($z_{\text{top}} < 3 \text{ km}$)



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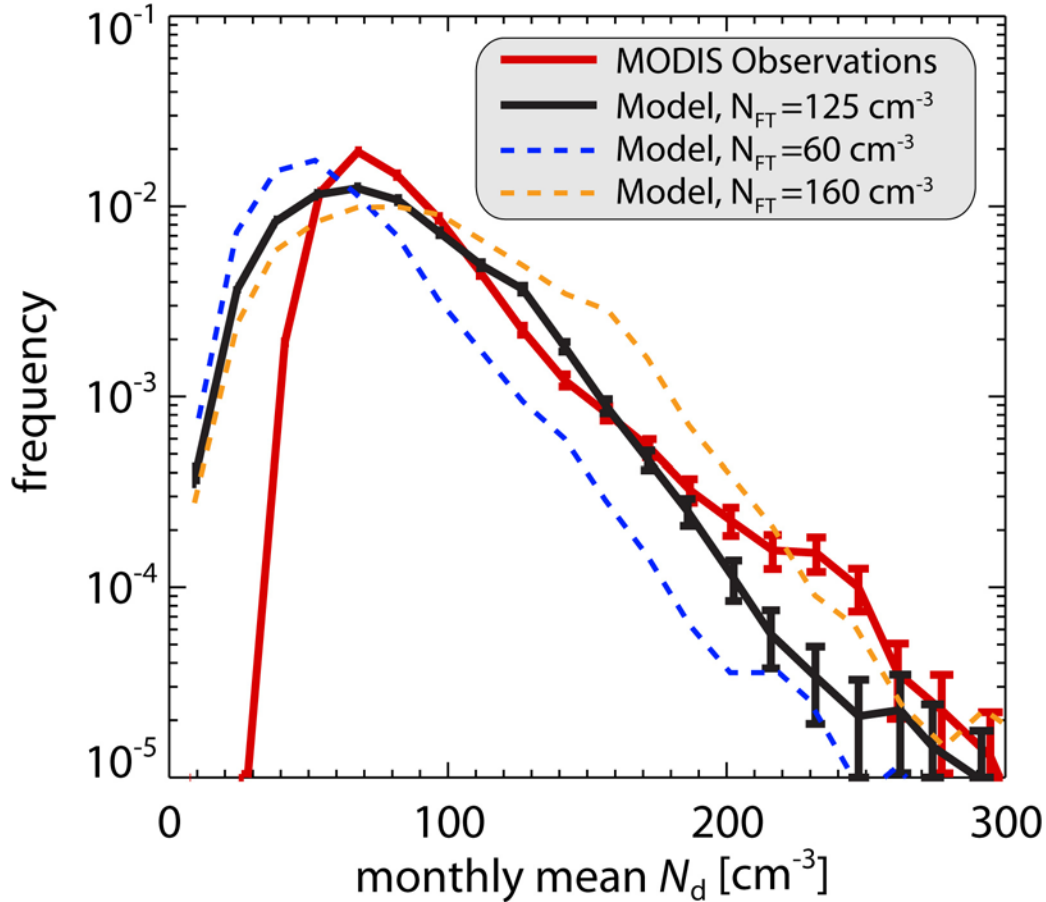
Figure 6: Mean precipitation rate at cloud base from low clouds (cloud top height $z_{\text{top}} < 3 \text{ km}$) estimated with spaceborne radar measurements from CloudSat [Lebsock and L'Ecuyer 2011]. Data are screened to display regions of extensive marine low clouds under conditions of mean subsidence as in Fig. 4.



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Figure 7: Frequency distributions of observed and modeled monthly mean cloud droplet concentration. Only months that meet the criteria needed to contribute to the means shown in Fig. 4 (regions with extensive low clouds under divergent conditions) are shown. Shown here are the base version of the model (i.e. that used to construct Fig. 4, solid black), together with estimates with no FT contribution to CCN (dashed), and no sea salt contribution (dotted). Error bars show the 95% confidence interval in the frequency estimates for the observations and for the model base case due to sampling limitations.

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Figure 8: Sensitivity to assumed free tropospheric CCN concentration N_{FT} of the frequency distributions modeled monthly mean cloud droplet concentration. Only months that meet the criteria needed to contribute to the means shown in Fig. 4 (regions with extensive low clouds under divergent conditions) are shown. Shown here are the base version of the model (i.e. that used to construct Fig. 4, solid black), together with model estimates with low ($N_{FT} = 80 \text{ cm}^{-3}$) and high ($N_{FT} = 160 \text{ cm}^{-3}$) estimates of the FT contribution to CCN (dashed blue and orange respectively). Error bars show the 95% confidence interval in the frequency estimates for the observations and for the model base case due to sampling limitations.