

1 **On the Nature and Extent of Optically-Thin Marine Low Clouds**

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22 **Abstract**

23 Macrophysical properties of optically-thin marine low-clouds over the non-polar oceans  
24 (60°S-60°N) are measured using two years of full-resolution nighttime data from the Cloud-  
25 Aerosol Lidar with Orthogonal Polarization (CALIOP). Optically-thin clouds, defined as  
26 the subset of marine low clouds that do not fully attenuate the lidar signal, comprise almost  
27 half of the low-clouds over the marine domain. Regionally, the fraction of low clouds that  
28 are optically-thin ( $f_{\text{thin,cld}}$ ) exhibits a strong inverse relationship with the low-cloud cover,  
29 with maxima in the tropical trades ( $f_{\text{thin,cld}} > 0.8$ ) and minima in regions of persistent marine  
30 stratocumulus and in midlatitudes ( $f_{\text{thin,cld}} < 0.3$ ). Domain-wide, a power  
31 law fit describes the cloud length distribution, with exponent  $\beta = 2.03 \pm 0.06$  ( $\pm 95\%$   
32 confidence interval). On average, the fraction of a cloud that is optically-thin decreases  
33 from  $\sim 1$  for clouds smaller than 2 km, to  $< 0.3$  for clouds larger than 30 km. This  
34 relationship is found to be independent of region, so that geographical variations in the  
35 cloud length distribution explain three-quarters of the variance in  $f_{\text{thin,cld}}$ . Comparing  
36 collocated trade cumulus observations from CALIOP and the airborne High Spectral  
37 Resolution Lidar reveals that clouds with lengths smaller than are resolvable with CALIOP  
38 contribute approximately half of the low clouds in the region sampled. A bounded cascade  
39 model is constructed to match the observations from the trades. The model shows that the  
40 observed optically-thin cloud behavior is consistent with a power law scaling of cloud  
41 optical depth and suggests that most optically-thin clouds only partially fill the CALIOP  
42 footprint.

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## 46 1.0 Introduction

47 The planetary albedo is a quantity determined by the reflection of sunlight by the  
48 Earth's surface and atmosphere. It is a critical determinant of the Earth's temperate and  
49 climate, and it is therefore imperative that its controlling factors be understood and  
50 adequately quantified. Without clouds, the planetary albedo would be about 0.15, due to the  
51 low reflectance of the Earth's surface (most of which is ocean). Highly reflective surfaces,  
52 such as those characteristic of the polar regions and unvegetated deserts, are notable  
53 exceptions but cover a small part of the planet and do not contribute strongly to the average  
54 surface albedo. However, the average albedo of the Earth (~0.3) is roughly doubled by the  
55 presence of clouds [Harrison *et al.*, 1990], and low clouds are the dominant type by  
56 contribution to albedo [Hartmann *et al.*, 1992].

57 Among albedo-influencing components of the Earth system, clouds are unique in  
58 both their extreme transience and in their remarkable range of sizes and forms. Although  
59 we have a reasonably accurate assessment of the global impact of clouds on the planetary  
60 albedo, our understanding of how clouds of different sizes and types contribute to the  
61 albedo remains rather poor. Although it is known that low clouds are the dominant  
62 contributor to the albedo globally [Hartmann *et al.*, 1992], it is uncertain what the relative  
63 contributions of the different low cloud types (cumulus, stratocumulus, stratus) are to the  
64 overall albedo of low clouds.

65 Moist convection over the oceans has a varied morphology. Marine low clouds  
66 range from expansive stratocumulus decks stretching hundreds of kilometers, to small trade  
67 wind cumulus a few tens to hundreds of meters in extent [Hozumi *et al.*, 1982; Wood and  
68 Field, 2011]. We are currently unable to quantify how different low-cloud systems  
69 distribute condensate horizontally and vertically to determine their albedo. This is partly  
70 because we do not fully understand how albedo scales with cloud cover, as the latter is  
71 frequently difficult to measure with the moderate-resolution (~1 km) passive sensors used to  
72 generate our cloud cover climatologies [Zhao and Di Girolamo, 2006; Jones *et al.*, 2012].  
73 We still do not have proven measures of cloud condensate amounts, particularly for broken  
74 and optically-thin cloud fields [Turner *et al.*, 2007]. Many questions are therefore still  
75 largely unanswered. For example, per unit of cloud cover, are trade cumulus clouds  
76 optically thinner than stratocumulus? What fraction of low clouds is optically-thin and how  
77 are these optically-thin clouds partitioned by cloud size and type? What resolution do our  
78 passive sensors need to have in order to accurately determine cloud cover?

79 Satellite data show that the albedo of extensive marine low clouds is quite low  
80 [Bender *et al.*, 2011], and that their optical thickness is highly variable (e.g., Rossow *et al.*  
81 [2002]). It is reasonable to infer that a significant fraction of them must be optically-thin.  
82 These optically-thin clouds are not currently well resolved observationally, but are globally  
83 pervasive and potentially important for climate studies [Garret and Zhao, 2006; Turner *et al.*  
84 *et al.*, 2007; Zuidema *et al.*, 2012]. For this work, we define operationally an optically-thin  
85 cloud as a cloud that does not fully attenuate a lidar backscatter signal. This corresponds to  
86 unbroken clouds with optical depth ( $\tau_{\text{cld}}$ ) less than ~3. Analysis of  $\tau_{\text{cld}}$  from Moderate  
87 Resolution Imaging Spectroradiometer (MODIS) Level 3 liquid cloud retrievals over an  
88 ocean region spanning ( $\pm 45^\circ$  latitude) suggests that clouds with  $\tau_{\text{cld}} < 3$  comprise ~30% of  
89 marine low clouds, and contribute perhaps 15% of their albedo (Fig. 1). But since many of  
90 these clouds are known to be smaller than the resolution of the MODIS pixels (e.g., Zhao  
91 and Di Girolamo [2007]), it is important to appreciate that these estimates are likely to be a  
92 function of sensor resolution. How does the picture change when clouds are viewed with an  
93 active sensor with a much smaller footprint size? Indeed, does the use of metrics such as

94 cloud cover even make sense if a significant fraction of the clouds have very low optical  
95 thicknesses?

96       Optically-thin and broken clouds are not only important because of their  
97 contribution to the planetary albedo, but also because aerosol retrievals are problematic  
98 where clear regions are in close proximity to clouds and contamination of clear pixels by  
99 sub-pixel scale optically-thin clouds is common (e.g., *Zhang and Reid* [2005], *Charlson et*  
100 *al.* [2007], *Koren et al.* [2007]). A trade wind cumulus study reported that a third of  
101 MODIS pixels diagnosed by the cloud mask as being clear actually contained sub-pixel  
102 scale cloud [*Zhao and Di Girolamo*, 2006]. Over the global ice-free oceans, 50% of marine  
103 low clouds are separated from each other by less than 5 km [*Várnai and Marshak*, 2011],  
104 suggesting that a high fraction of clear sky over ocean is in close proximity to cloud. Other  
105 recent studies of the clear-cloud transition region using both *in situ* [*Twohy et al.*, 2009] and  
106 remote sensing measurements [*Su et al.*, 2008; *Tackett and Di Girolamo*, 2009; *Redemann*  
107 *et al.*, 2009] suggest that aerosol properties in the vicinity of clouds are different from those  
108 2-5 km from the cloud. Some of this difference may be attributed to aerosol hygroscopic  
109 growth in the high relative humidity air surrounding cloud, but the relative magnitudes of  
110 cloud contamination and aerosol hygroscopic growth effects on retrieved aerosol properties  
111 is not known. Observations show that enhancement of near-cloud reflectances due to  
112 undetected or unresolved cloud contamination of clear pixels leads to passive sensor aerosol  
113 optical depth retrieval overestimates of 10-20% [*Zhang and Reid*, 2005].

114       Although passive sensors provide excellent coverage, it is fundamentally  
115 challenging to identify which atmospheric constituents are producing the measured top of  
116 atmosphere (TOA) signals, and at what altitudes they reside. Spaceborne lidar provides a  
117 metric of the TOA reflectance, integrated over the specified layers of interest, and offers a  
118 closer link to the atmospheric constituent producing that signal than passive sensors  
119 provide. A frequency distribution of such a metric, integrated attenuated backscatter  
120 (IABS), derived from high-resolution (90 m telescope footprint) Cloud-Aerosol Lidar with  
121 Orthogonal Polarization (CALIOP) is presented in Fig. 2. A bimodal distribution of IABS  
122 is observed with modes representing clear and cloudy atmospheric states as indicated (Fig.  
123 2a). The non-zero minimum between clear and cloudy states, which exemplifies the  
124 fundamental challenge in separating clouds from clear sky, has been termed the  
125 “continuum” region [*Charlson et al.*, 2007]. With detailed classification with CALIOP’s  
126 full-resolution cloud clearing algorithm [*Vaughan et al.*, 2009]), spaceborne lidar is able to  
127 inform us that this “continuum”, or clear-cloud transition region, is populated almost  
128 exclusively with clouds (Fig. 2b). These clouds must, by their virtue of having low visible  
129 reflectance, be optically-thin. IABS can be small either because the cloud is optically-thin  
130 or because the laser was attenuated fully some distance into an optically-thick cloud. The  
131 latter case is manifest in Fig. 2b as the smaller values of IABS categorized as optically-thick  
132 clouds. A frequency distribution alone cannot inform us of their spatial structure, but  
133 previous work suggests a prevalence of detraining cloud elements, dissipating, or nascent  
134 clouds in marine low-cloud systems [*Koren et al.*, 2009]. However, the CALIOP data are  
135 showing us that optically-thin clouds are not only limited to the “continuum” region  
136 (approximately defined as the IABS region spanning  $5 \times 10^{-2} \text{ sr}^{-1}$  to  $3 \times 10^{-1} \text{ sr}^{-1}$ ), but  
137 contribute to TOA reflectance throughout the entire cloud IABS distribution (Fig. 2b).  
138 Indeed, over non-polar oceans, we find that almost half of the marine low clouds are  
139 optically-thin by our definition of surface-detectability.

140       Despite the apparent prevalence of optically-thin clouds, very little is currently  
141 known about their macrophysical properties (spatial distribution, length, vertical

142 distribution, geometric depth) on a global scale. This study aims to provide insight into  
143 these properties, and thereby provide new information on a poorly characterized and  
144 understood subset of marine low clouds. Compared to satellite observations, large scale  
145 models tend to overestimate the frequency of optically-thick low clouds and underestimate  
146 the frequency of optically-thin low clouds (e. g., *Kay et al.*, [2012] a Community  
147 Atmosphere Model study). At the low end of the optical depth distribution, model  
148 frequencies are close to an order of magnitude too low compared to passive remote sensing  
149 satellite observations [*Kay et al.*, 2012]. However, it is important to appreciate that there  
150 are major differences between the various observational datasets [see Fig. 3 in *Kay et al.*,  
151 2012], making accurate assessment of model representation of optically-thin low clouds  
152 difficult.

153 The primary dataset employed for this study is the CALIOP full-resolution cloud  
154 mask (60 m vertical resolution and 90 m footprint diameter, spatially sampled every 335 m)  
155 at 1064 nm. In addition to providing new near-global observations of the macrophysical  
156 properties of optically-thin marine low clouds, we also investigate possible sensor-  
157 resolution detection issues by comparing CALIOP with data from NASA's airborne High  
158 Spectral Resolution Lidar (HSRL). The ensuing sections of this manuscript are ordered as  
159 follows. Dataset description and analysis methodology are detailed in Section 2. Results  
160 are presented in Section 3. A discussion of the implications is augmented and informed by  
161 the use of a simple fractal model in Section 4. Our findings are summarized in Section 5.

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## 164 2.0 Dataset Description and Analysis Method

### 165 2.1 CALIOP Dataset

166 Two years (September 2006-August 2008) of vertically-resolved cloud data from the  
167 full-resolution CALIOP Vertical Feature Mask (VFM) Version 3.01 [Vaughan *et al.*, 2009]  
168 over the non-polar oceans ( $\pm 60^\circ$  latitude) are used in this study. The VFM provides  
169 vertically-resolved cloud and aerosol layer information generated by the CALIOP multi-  
170 resolution layer identification and scene classification algorithms. Full-resolution layer  
171 information is obtained from the high-resolution cloud clearing tool within the layer  
172 detection algorithm. In the lowest 4 km of the atmosphere the 1064 nm channel (60 m  
173 vertical resolution) is used for full-resolution cloud-layer identification, and any feature  
174 detected above the surface is classified as a cloud [Vaughan *et al.*, 2009]. This wavelength  
175 is less sensitive to aerosol than the 532 nm channel and this improves cloud-aerosol  
176 discrimination. A detailed mission overview and comprehensive instrument details are  
177 presented by Winker *et al.* [2009] and Hunt *et al.* [2009], respectively. CALIOP has high  
178 sensitivity at low cloud optical depths: nighttime minimum detectable particulate extinction  
179 at 1064 nm is  $0.165 \text{ km}^{-1}$  for layers greater than 180 m in depth, equivalent to  $\tau_{\text{cld}}$  greater  
180 than 0.03, and is  $1.65 \text{ km}^{-1}$  for layers 60-180 m (minimum 60 m), equivalent to  $\tau_{\text{cld}}$  greater  
181 than 0.1 (Chepfer, H., G. Cesana, D. Winker, B. Getzewich, and M. Vaughan, Comparison  
182 of two different cloud climatologies derived from CALIOP-Level 1 observations: the  
183 CALIPSO-ST and the CALIPSO-GOCCP, submitted to *Journal of Atmospheric and*  
184 *Oceanic Technology*, 2012). CALIOP provides near-global ( $\pm 82^\circ$  latitude), vertically-  
185 resolved, atmosphere-only backscatter data, with no contamination from surface return or  
186 3-D radiation/cloud adjacency effects, unlike radiances measured from passive sensors  
187 [Wen *et al.*, 2007; Várnai and Marshak, 2009]. However, sampling is sparse. The lidar  
188 return signal is limited to a 90 m diameter receiver footprint, so no off-track atmospheric  
189 information can be provided. Results are only representative of a single sampling time for  
190 every location. In addition, consecutive laser footprints are not contiguous; horizontal  
191 separation between each footprint is 335 m and consecutive nighttime (and daytime) orbits  
192 are  $\sim 2,600$  km apart at the Equator, requiring a temporally extensive dataset for  
193 representative sampling.

194 Our focus is optically-thin marine low clouds. In this study, a low cloud is defined  
195 as having a cloud top height (CTH) less than or equal to 3 km. A high cloud is any cloud  
196 higher than this. We define an optically-thin cloud as a cloud detected at full-resolution that  
197 does not fully attenuate the lidar signal (i.e., one for which the ocean surface was also  
198 detected in the same profile), corresponding to  $\tau_{\text{cld}}$  less than or equal to 3. Optically-thick  
199 clouds are all other low cloud not classified as optically-thin. We screen the data for both  
200 high clouds, and incidences of low clouds where more than one cloud-layer is detected, the  
201 latter representing a small fraction of the marine low cloud data, 0.06 on average. To ensure  
202 the highest fidelity in our results, we use only nighttime data for optimal detection of  
203 optically-thin clouds; the solar background signal degrades the daytime signal-to-noise ratio  
204 (SNR) [Hunt *et al.*, 2009]. The daytime minimum detectable particulate extinction at 1064  
205 nm is  $0.28 \text{ km}^{-1}$ , equivalent to  $\tau_{\text{cld}}$  greater than 0.05 for layers greater than 180 m in depth  
206 and is  $2.8 \text{ km}^{-1}$  for layers 60-180 m in depth, equivalent to  $\tau_{\text{cld}}$  greater than 0.17. Inspection  
207 of daytime IABS distributions does not reveal large differences between daytime and  
208 nighttime data, but it is not possible to separate real diurnal differences from potential  
209 differences in detectability and other noise-related artifacts. We note that some aerosol  
210 layers (e.g., Saharan dust) may be dense enough to be misclassified as cloud at full-  
211 resolution, however, this occurs mostly over land and we assume the effects over ocean to

212 be negligible. In this work a profile is defined as a column of vertically-resolved VFM clear  
213 or cloud classifications derived from a single lidar pulse. We further define a cloudy profile  
214 as a full-resolution profile with one or more 60 m vertically-adjacent cloud layers. Multi-  
215 layered low clouds, comprising less than 6% of the high-cloud screened low cloud dataset,  
216 are excluded from this analysis. In summary, our primary dataset is comprised of clear  
217 profiles and single-layer cloudy profiles with CTH less than or equal to 3 km detected at  
218 full-resolution at night over non-polar oceans.

219 To provide an example of optically-thin clouds and their spatial context, we present  
220 collocated Wide Field Camera (WFC) band-averaged radiance (620 – 670 nm) for a 10 km  
221 orbit segment at 125 m horizontal resolution, and CALIOP full-resolution Level 1 daytime  
222 IABS at 532 nm integrated from 0 to 20 km, for a broken cloud field in the northwest  
223 subtropical Atlantic (Fig. 3). Along with CALIOP, the WFC is part of the Cloud Aerosol  
224 Lidar with Integrated Pathfinder Satellite Observation (CALIPSO) satellite’s payload. The  
225 WFC context (Fig. 3a) suggests that some detected optically-thin cloud profiles are cloud  
226 edges that partially fill the lidar field of view. We explore this issue further in Section 4. In  
227 general, peaks and troughs observed in both the radiance and IABS data track each other  
228 remarkably well (Fig. 3b). Optically-thick cloud profiles (fully attenuated profiles) tend to  
229 be associated with stronger visible radiances. However, radiances associated with optically-  
230 thick clouds at ~6 km along the orbit segment are similar to some of the radiances for  
231 optically-thin cloud profiles elsewhere. The lack of a complete correspondence between  
232 radiance and IABS is indicative of the different nature of the lidar backscatter signal and the  
233 visible reflectance. Multiple scattering affects the signals in different ways. In addition, it  
234 is likely that beam filling in this broken cloud field affects visible radiance and lidar signals  
235 differently. The clear evidence in this example for small and often optically-thin cloud  
236 elements on the scale of the lidar footprint compels us to understand how optically-thin  
237 clouds are related to cloud horizontal scales globally and regionally.

## 238 239 2.2 HSRL Dataset

240 In addition to a near-global examination of CALIOP’s view of optically-thin marine  
241 low clouds, we also compare collocated CALIOP cloud fractions and cloud lengths with  
242 those estimated using the airborne NASA LaRC High Spectral Resolution Lidar (HSRL)  
243 [*Hair et al.*, 2008]. Although HSRL is typically employed for aerosol studies and/or  
244 instrument validation (e.g., *Kacenenbogen et al.* [2011], *Rogers et al.* [2011]) for this  
245 work we use derived HSRL altitudes of the ocean surface and cloud top. Cloud data for the  
246 comparison are drawn from four daytime and four nighttime spatially and temporally  
247 matched HSRL underflights of the CALIOP orbit track over the tropical and subtropical  
248 western Atlantic (Table 1), and we include only data for flight segments with aircraft  
249 altitude greater than 8 km. HSRL cloud screening is based upon a wavelet transform  
250 technique [*Su et al.*, 2008]. HSRL and CALIOP do not have the same footprint dimensions  
251 due to different sampling configurations. The HSRL sample area is composed of 60 m  
252 along-track averages of contiguous 8 m single-shot footprints, yielding a sample area  
253 8 mx60 m in the across-track and along-track directions, respectively. The 60 m along-  
254 track dimension results from 2 Hz sampling at aircraft ground-speed of 120 m s<sup>-1</sup> (mean  
255 value for data analyzed herein). However, CALIOP 90 m diameter footprints are  
256 horizontally separated by 335 m in the along-track direction. Henceforth, we refer to an  
257 instrument footprint, pixel size or sample area as a field of view (FOV) for clarity of  
258 notation. HSRL data have 30 m vertical resolution. In keeping with the CALIOP analysis,  
259 we identify cloudy HSRL profiles as those with CTH less than or equal to 3 km. We use

260 the HSRL ocean surface altitude parameter to determine whether a surface return was  
261 detected and thereby discriminate between low and high optical depth clouds, i.e., whether  
262 or not the signal was fully-attenuated by the cloud before reaching the surface. For the  
263 ocean data used, we expect the surface altitude for a cloud profile to be  $\sim 0$  m if the signal is  
264 not fully attenuated. However, the surface altitude detection algorithm is not fully  
265 developed (HSRL Data Summary, unpublished material, 2009) and the detected surface  
266 altitude departs somewhat from zero, even for cloud-free profiles over the ocean.  
267 Therefore, we define a cloud profile with surface altitude less than or equal to 85 m as  
268 optically-thin. The upper limit of 85 m is the mean surface altitude plus three standard  
269 deviations observed for cloud-free HSRL profiles over the ocean. For the HSRL-CALIOP  
270 comparisons, we include only HSRL and CALIOP data that achieve temporal coincidence  
271 within  $\pm 15$  minutes. Both datasets are high-cloud-screened using the CALIPSO VFM  
272 described in Section 2.1

273

## 274 2.3 Analysis Methods

275 After screening for high cloud, CALIOP full-resolution data are binned into  $5^\circ \times 5^\circ$   
276 latitude-longitude gridboxes. Gridding the data allows us to examine the geographical  
277 variations of the various cloud features observed. We define a transect as being the portion  
278 of ground track for each CALIPSO orbit within a gridbox. Separate analysis methods are  
279 applied to these binned data to produce cloud cover statistics, and cloud length distributions.

280

### 281 2.3.1 Cloud Cover Analysis Method

282 Monthly cloud cover values calculated from counts of full-resolution clear profiles  
283 and low-cloud profiles within a gridbox (Table 2) are averaged to produce two-year mean  
284 low-cloud cover values for each gridbox. Domain-wide cloud cover is the mean of all  
285 gridbox values.

286 Minima in high-cloud-screened CALIOP profile counts (Fig. 4a) exhibit a similar  
287 pattern to ISCCP – D2 cloud amount maxima for mid plus high clouds (available from  
288 <http://isccp.giss.nasa.gov/products/browsed2.html>). ISCCP annual mean data (not shown)  
289 indicate maximum high cloud amounts in the western equatorial Pacific and northern Indian  
290 Ocean (amounts  $> 40\%$ ), and mid-level cloud amount maxima (30-40%) in the Northern  
291 and Southern Hemisphere storm-track regions.

292

### 293 2.3.2 Cloud Length Distribution Method

294 Cloud length is calculated for individual cloud entities. We define a cloud entity as  
295 a series of consecutive marine low-cloud profiles, bounded on each end by a clear profile.  
296 The cloud length data are binned into logarithmically spaced size bins, with the lower size  
297 limit set by instrument resolution, and the upper limit defined by the sampling method  
298 discussed below. An additional criterion is applied to the size distribution data that is not  
299 required for cloud cover analysis: we require a transect to comprise at least 100 km of  
300 consecutive high-cloud-screened profiles, thus ensuring a long enough transect length to  
301 capture most of the contributions to cloud cover from clouds of different sizes. A transect  
302 may be less than 100 km, and therefore excluded from the cloud length dataset, because a) it  
303 “clips” the edge of a  $5^\circ \times 5^\circ$  gridbox, b) high cloud screening eliminates a portion of the  
304 transect so that the remaining segment lengths are shorter than 100 km, or c) a portion of  
305 transect is land. However, even transects longer than 100 km have a maximum length  
306 ( $\sim 550$  km) imposed by the  $5^\circ \times 5^\circ$  gridding. For every cloud length calculated we retain the  
307 length of the transect containing the observed cloud, and use this information to formulate a

308 simple correction to the size distribution (see *Wood and Field* [2011] for details) to account  
309 for this upper limit cut-off as follows,

$$310 \quad n_{corr}(L)dL = n(L)dL[L_{tran}/(L_{tran} - L_{min})] \quad (1)$$

311 where  $n_{corr}(L)$  is the corrected cloud count,  $L_{tran}$  is the mean of the binned transect lengths,  
312  $L_{min}$  is the bin minimum cloud length,  $n(L) dL$  is the number of clouds sampled with sizes  
313 between  $L$  and  $L + dL$ . To help elucidate the contributions of optically-thin and optically-  
314 thick profiles to clouds of a given size, we construct size distributions for “majority  
315 optically-thick” entities (contiguous cloudy profiles for which over 90% of contributing  
316 profiles are optically-thick) and “majority optically-thin” entities (over 90% of contributing  
317 profiles are optically-thin). The threshold of 90% was chosen to accommodate observed  
318 variability in cloud optical depth [*Wood and Taylor*, 2002; *Rossow et al.*, 2002], while  
319 ensuring that the clouds analyzed were predominantly optically thick or optically thin,  
320 respectively. The cloud length dataset is a subset of the dataset described in Section 2.3.1,  
321 and is comprised of approximately three times fewer profiles (cf. Fig. 4a and 4b) because of  
322 the minimum transect length restriction.

323

324

## 3.0 Results

### 3.1 Cloud Cover

High-cloud-screened marine low-cloud cover ( $f_{\text{cld}}$ ) is at a minimum throughout the deep tropics, and is at a maximum in the northern hemisphere and southern hemisphere storm-track regions, with additional maxima in the stratocumulus regions on the eastern side of the oceanic subtropical highs (Fig. 5a). Values range from 0.10 to greater than 0.80, with a two-year, domain mean of 0.50. The spatial pattern of cloud cover is similar to climatologies from surface observations (Warren et al., 2010, Climatic atlas of clouds over land and ocean, available from <http://www.atmos.washington.edu/CloudMap>) and remote sensing [Rossow and Schiffer, 1999]. Optically-thin low clouds are present in 10-30% of high-cloud-screened profiles over the non-polar oceans, with no strong geographic variation ( $f_{\text{thin}}$ ; Fig. 5b). When expressed as the fraction of marine low-cloud profiles that are optically-thin ( $f_{\text{thin,cld}}$ ), a clear spatial pattern emerges, with maxima in the trade wind cumulus regions, and minima in persistent stratocumulus and storm-track regions (Fig. 5c). The two-year mean value of  $f_{\text{thin,cld}}$  is 0.45 demonstrating that almost half of all marine low clouds are optically thin. While 0.45 is higher than that from MODIS (Fig. 1), both sensors indicate a major role for optically-thin clouds in determining global cloud cover.

A remarkably strong negative correlation between  $f_{\text{thin,cld}}$  and  $f_{\text{cld}}$  exists at both monthly and annual timescales (Fig. 6). In the trade wind regions, where marine low-cloud cover is lower than 0.25 [Fig. 5a, and Medeiros et al., 2010], it is remarkable that consistently greater than 80% of the observed clouds are optically thin. In regions typical of marine stratocumulus, where nighttime marine low-cloud cover exceeds 0.80,  $f_{\text{thin,cld}}$  is consistently lower, with values of 0.30 or less being typical. The presence of a large amount of marine low cloud that is optically-thin on the scale of the relatively small CALIOP footprint is surprising and warrants further investigation. We devote much of this paper to analyzing these optically-thin clouds. However, we should note that this result is consistent with the TOA albedo from Clouds and the Earth's Radiant Energy System (CERES) data (available from <http://eos.atmos.washington.edu/cgi-bin/ceres/disp.pl?ceres.alb.ann.d>) of the trade wind regions being quite low (typically 0.16) despite cloud coverage in the trade wind regions of 0.20–0.25. Given observed trade wind TOA albedo and cloud cover, and using a simple cloud albedo – cloud optical depth relation [Lacis and Hansen, 1974] and a simple TOA albedo model (see Section 4), we estimate that cloud optical depth  $\tau_{\text{cld}}$  would need to be 2.7–3.0 to produce the observed TOA albedo, if all clouds have the same  $\tau_{\text{cld}}$ . While the prevalence of optically-thin clouds may appear surprising, it is consistent with regional albedo observations.

### 3.2 Cloud Length Distribution

Here, we examine the properties of marine low clouds, and especially optically-thin clouds, as a function of their horizontal scale. Results presented in this section are derived from analysis of cloud entities. We find that cloud length distributions approximately follow a power law fit ( $n(L) = \alpha L^{-\beta}$ ), where  $n(L)$  is the number of clouds of size  $L$  to  $L+dL$ ,  $L$  is the bin center, and  $\alpha$  and  $\beta$  are constants, consistent with previous studies [Zhao and Di Girolamo, 2007; Wood and Field, 2011; Benner and Curry, 1998]. For  $\beta = 2$ , each logarithmic size interval (e.g., 100m-1km, 1km-10km, etc.) contributes equally to total cloud cover and  $\beta > (<) 2$  implies that the distribution is weighted towards smaller (larger) clouds [Wood and Field, 2011].

For the CALIOP dataset comprising all low clouds, the size distribution is well fitted by a single power law with  $\beta = 2.03 \pm 0.06$  (error at  $2\sigma$  level). This is higher than the value

373 of  $1.66 \pm 0.04$  reported by *Wood and Field* [2011] for a near-global cloud dataset comprising  
374 aircraft and satellite data. One explanation may be differences in cloud sampling at the low  
375 and high ends of the size distribution. Although *Wood and Field* [2011] aircraft data have a  
376 higher resolution than CALIOP,  $\sim 100$  m compared to 335 m, respectively, the aircraft did  
377 not sample regions where small cumulus is the dominant cloud type [*Wood and Field*, 2000,  
378 2011]. Further, CALIOP may undersample large clouds as the upper cloud length limit for  
379 this study is  $\sim 550$  km compared to  $\sim 4000$  km in *Wood and Field* [2011]. Thus, our dataset  
380 tends to sample smaller clouds on average, and smaller clouds tend to have larger exponents  
381 (see below, and Fig. 12 in *Wood and Field* [2011]).

382 “Majority optically-thin” clouds (defined in Section 2.3.2) constitute nearly all of the  
383 clouds smaller than 4 km, whereas clouds larger than 35 km have a greater contribution  
384 from optically-thick profiles (Fig. 7). This highlights the increasing importance of  
385 optically-thin clouds as clouds become smaller.

386 To further explore differences in cloud length distributions, we focus on three  
387 specific cloud regimes (Table 3): 1) a region of persistent stratocumulus off the California  
388 coast; 2) a tropical Pacific stratocumulus to cumulus (Sc-Cu) transition region; 3) a trade  
389 wind cumulus region of the tropical Pacific. Cloud length distributions for the regions of  
390 stratocumulus and Sc-Cu transition follow a  $\beta < 2$  power law fit with values of  $1.81 \pm 0.08$   
391 and  $1.68 \pm 0.09$ , respectively, indicating that large clouds dominate cloud cover (Fig. 7,  
392 inset). As expected, the trade wind cloud length distribution has a  $\beta$  greater than 2  
393 ( $2.49 \pm 0.25$ ) and is dominated by clouds with lengths less than 2 km, consistent with prior  
394 studies of trade wind cumulus [*Benner and Curry*, 1998; *Zhao and Di Girolamo*, 2007].  
395 Since trade wind regions consist largely of optically-thin clouds (Fig. 5c), these results paint  
396 an emerging picture of small marine low clouds being optically-thin. This has implications  
397 for the remote sensing of cloud and aerosol properties given current sensor resolution for  
398 most passive satellite instruments used to construct cloud climatologies.

399

### 400 3.2.1 Cloud Length at Median Cloud Cover

401 We define a cloud length at median cloud cover ( $L_{50}$ ) for each of the three cloud  
402 categories (marine low cloud, “majority optically-thin” cloud, “majority optically-thick”  
403 cloud) such that clouds with lengths up to  $L_{50}$  contribute 50% to cloud cover for that  
404 category. Regions with the smallest coverage of low cloud are associated with smaller  $L_{50}$   
405 values (cf. Fig. 5a and Fig. 8a), with  $L_{50} < 2$  km in trade wind cumulus regions and  
406  $L_{50} > 30$  km in stratocumulus regions. Our results for the trades are qualitatively consistent  
407 with a recent study in the tropical western Atlantic trade wind region reporting  $L_{50}$  of 2 km  
408 using high-resolution passive visible satellite data [*Zhao and Di Girolamo*, 2007].  
409 Additionally, we observe an inverse relationship between marine low cloud  $L_{50}$  and  $\beta$  (not  
410 shown) also noted by *Wood and Field*, [2011], indicating that variations in  $\beta$  explain almost  
411 one-third of the observed geographical variance in  $L_{50}$ . Instances of horizontally extensive  
412 (greater than 10 km) “majority optically-thin” clouds are observed (Fig. 7) but with an order  
413 of magnitude smaller frequency than “majority optically-thick” clouds at these lengths.  
414 Most notably, poleward of  $40^\circ$  in both hemispheres, “majority optically-thin”  $L_{50}$  ranges  
415 from 5 km to greater than 10 km (Fig. 8b). Cloud cover in these regions is mainly Sc  
416 (Warren et al., 2010, Climatic atlas of clouds over land and ocean  
417 <http://www.atmos.washington.edu/CloudMap/WebO/index.html>). These relatively large  $L_{50}$   
418 values are consistent with the dimensions of clouds in shallow open cellular convection  
419 [*Wood and Hartmann*, 2006]. “Majority optically-thick” and “majority optically-thin”  $L_{50}$

420 maps support our earlier statement that optically-thin marine low clouds are in general,  
421 smaller in size than optically-thick marine low clouds (Fig. 8b and 8c).

### 422 3.2.2 Cloud Length Distribution from HSRL

423 The horizontal separation between consecutive CALIOP pulses is 0.335 km but  
424 many small cumuli are only a few tens to hundreds of meters in size [*Plank 1969; Hozumi*  
425 *et al.*, 1982; *Wielicki and Welch*, 1986; *Koren et al.*, 2008]. The inter-profile region not  
426 sampled by CALIOP may be populated with one or more small clouds with clear gaps, or  
427 may be entirely clear. For our cloud length estimates, if two consecutive profiles are  
428 flagged as cloudy we assume that the intervening atmosphere is cloudy, and this could lead  
429 to CALIOP overestimating cloud lengths. To examine the extent to which this may occur,  
430 we compare spatially matched CALIOP and HSRL cloud length data, noting that HSRL  
431 footprints are contiguous. Data are from the tropical western Atlantic trade wind region  
432 (Section 2.2 for dataset description).

433 A minimum transect length criterion of 30 km is applied to the HSRL-CALIOP  
434 comparison dataset for estimation of cloud lengths. This is shorter than for our analysis  
435 from CALIPSO alone (100 km), because the  $\pm 15$  minute temporal matching limits transect  
436 lengths to less than 200 km, and we find that a minimum length requirement of 100 km  
437 greatly reduces the data available for comparison in what is already a relatively small  
438 dataset of collocated data. However, this work (Fig. 8) and other studies indicate that  
439 characteristic cloud lengths in the tropical western Atlantic region are substantially smaller  
440 than 30 km [*Benner and Curry*, 1998; *Wielicki and Welch*, 1986; *Zhao and Di Girolamo*,  
441 2007], and so our reduced transect length threshold for the HSRL/CALIOP size distribution  
442 comparison should not produce large problems.

443 The CALIOP and HSRL cloud length distributions are broadly similar for sizes  
444 larger than 1 km, but CALIOP samples more large clouds (Fig. 9). This is because HSRL is  
445 able to sample smaller clouds than CALIOP can resolve and such clouds comprise 55% of  
446 the clouds sampled by HSRL. Sub-sampling HSRL as CALIOP (i.e., we subsample HSRL  
447 profiles every  $\sim 335$  m) shifts the HSRL cloud length distribution to larger sizes, closely  
448 matching the CALIOP data (Fig. 9), suggesting that the size distribution difference at cloud  
449 lengths smaller than 1 km is primarily due to sparse sampling by CALIOP rather than FOV  
450 size differences (HSRL FOV is 60 m; CALIOP FOV is 90 m). CALIOP and HSRL mean  
451  $f_{\text{cld}}$  values (0.24 and 0.31, respectively) agree to within 2 standard deviations, but  $f_{\text{thin,cld}}$   
452 values do not agree, being 0.89 (CALIOP) and 0.61 (HSRL). This effect cannot be due to  
453 the relative along-track sparsity of CALIPSO sampling. Differences in HSRL and CALIOP  
454 cloud detection algorithms may be partly responsible for the  $f_{\text{thin,cld}}$  disparity (e.g., the HSRL  
455 layer detection method may not be well suited to detecting attenuated surface returns).  
456 Another likely cause is the narrow HSRL FOV (8 m across-track compared with 60 m  
457 along-track), giving rise to the possibility that HSRL will miss some small clouds that lie on  
458 either side of the HSRL FOV, but which lie inside the CALIOP FOV (see schematic in Fig.  
459 9, inset). Since smaller clouds tend to be optically-thin, the clouds missed by HSRL are  
460 likely to be optically-thin.

461 The HSRL size distribution in Fig. 9 shows a roll-off at cloud lengths below  
462 about 0.5 km. This may be a hint that there exists a lower limitation to the cloud length of  
463 the smallest clouds, i.e., a scale break in the cloud length distribution. We note that this  
464 length is consistent with the typical length scale of turbulent eddies in the boundary layer.  
465 Without further evidence from observations and models we cannot be sure if the scale break  
466 is a real physical feature or reflects sampling a mixture of distributions on different days,  
467 but very high-resolution large eddy simulations could be used to examine this. We note that

468 one high-resolution Landsat satellite study [Sengupta et al., 1990] reported a scale break in  
 469 cumulus cloud length distribution at 1 km, with larger (smaller)  $\beta$  observed for larger  
 470 (smaller) clouds. The scale break was observed at a larger size (1 km) than in the HSRL  
 471 data shown here ( $\sim 500$  m). Such a scale break is not a universal feature in trade cumulus  
 472 [Zhao and Di Girolamo, 2007]. Further studies with high-resolution airborne lidar would  
 473 be helpful to understand any physical limits to the population of the smallest clouds.

474 Cloud lengths at median  $f_{\text{cld}}$  for the matched HSRL-CALIOP data indicate that  $L_{50}$   
 475 is 1.3 km for HSRL and 1.7 km for CALIOP. The  $L_{50}$  value from two years of CALIOP  
 476 data for this region is 1.3 km, close to the matched dataset value. This is also consistent with  
 477 high resolution airborne imaging radiometer measurements in other trade wind regions such  
 478 as the Indian Ocean [McFarquhar et al. 2004]. We therefore believe that the conclusions we  
 479 derive (see Section 6) from the HSRL-CALIOP comparison are somewhat representative of  
 480 trade cumulus regions throughout the tropics.

### 481 3.2.3 Cloud Length Distribution and Optically-thin Cloud Cover Variation

482 Another approach to examine how optically-thin profiles are distributed among  
 483 clouds of different sizes is to bin, as a function of cloud length  $L$ , the fraction of cloudy  
 484 profiles that are optically thin ( $f_{\text{thin},\text{cld}}^L$ ). For this we use the entire CALIOP cloud dataset  
 485 (1.7 million cloud entities). Surprisingly, the function  $f_{\text{thin},\text{cld}}^L$  is almost region-invariant  
 486 (Fig. 10). As cloud length increases, the fraction of optically-thin profiles comprising that  
 487 cloud decreases. Clouds of a given horizontal size possess a similar fraction of optically-  
 488 thin elements regardless of whether they occur in regions dominated by trade wind cumulus  
 489 or stratocumulus. Approximately 50% of profiles comprising clouds of size  $L \sim 15$  km, as  
 490 determined by CALIOP, are optically thin. Clouds smaller than 2 km are almost entirely  
 491 optically thin, whereas clouds larger than 100 km are less than 20% optically thin. A  
 492 similar relationship is observed in the HSRL data (not shown), but with the HSRL  $f_{\text{thin},\text{cld}}^L$   
 493 curve shifted to smaller cloud lengths for reasons discussed in the preceding section.

494 The near region-invariant nature of  $f_{\text{thin},\text{cld}}^L$  can be explained if each region  
 495 essentially experiences low clouds of all types, from small cumulus to relatively large sheets  
 496 of stratocumulus. Visual inspection of visible imagery suggests that this is a reasonable  
 497 supposition. Surprisingly, each of these cloud types appears to have very similar  
 498 characteristics in terms of its optical thickness distribution, regardless of location. Clearly,  
 499 however, the frequency with which each region experiences clouds of different sizes varies  
 500 dramatically from region to region. This result leads us to ask whether geographical  
 501 variations in the fraction of clouds that are optically thin might be largely explained by  
 502 regional differences in the cloud size distribution.

503 To test this, we make an estimate of the optically-thin fraction of marine low clouds  
 504 ( $\tilde{f}_{\text{thin},\text{cld}}$ ) using a single  $f_{\text{thin},\text{cld}}^L$  function based on the entire non-polar ocean dataset  
 505 (Fig. 10) together with the  $5^\circ \times 5^\circ$  binned two-year cloud length distributions  $n(L)$  as,

$$506 \quad \tilde{f}_{\text{thin},\text{cld}} = \frac{1}{\int_{L_{\text{min}}}^{L_{\text{max}}} n(L) dL} \int_{L_{\text{min}}}^{L_{\text{max}}} f_{\text{thin},\text{cld}}^L L n(L) dL \quad (2)$$

507 where  $L_{\text{max}/\text{min}}$  is the upper/lower CALIOP size range limit and  $L$  is the bin center cloud  
 508 length. We find that over non-polar oceans and separately over the tropics, the cloud length  
 509 distribution alone explains three-quarters of the variance in  $f_{\text{thin},\text{cld}}$  ( $R^2 = 0.73$  and  $0.77$ ,  
 510 respectively). What this result tells us is that knowledge of how the marine low-cloud  
 511 length distribution of all clouds varies is sufficient to accurately predict the geographical  
 512 variation in the optically-thin fraction of clouds across most of the ocean. This is important  
 513

514 because it means that the occurrence of optically-thin clouds is inherently tied to the size  
515 distribution of all low clouds. Further, it suggests that the processes controlling optically-  
516 thin clouds should not be considered as independent of the processes controlling low clouds  
517 in general.

### 518 3.3 Cloud Vertical Structure

519 Knowing the vertical distribution of cloud condensate is important for radiative transfer  
520 calculations and for model parameterization of cloud overlap. Vertically-resolved (60 m)  
521 optically-thin and optically-thick CTH, optically-thin cloud base height (CBH), and  
522 optically-thin cloud-layer depth are presented for each of the three regions/regimes  
523 discussed previously (Section 3.2; Table 3). Note that CBH cannot be observed for  
524 optically-thick clouds as, by definition, the lidar signal is fully attenuated by the cloud. The  
525 dynamical processes associated with each cloud regime are as follows: The vertical  
526 development of marine stratocumulus is suppressed by subsidence associated with either the  
527 descending branch of the Hadley/Walker circulation or subsiding air masses associated with  
528 mid-latitude cyclones, depending upon location. As a result, the marine stratocumulus layer  
529 is capped by a strong temperature inversion, resulting in extensive low clouds [*Klein and*  
530 *Hartmann, 1993*]. As the stratocumulus air mass moves equatorward it is advected over  
531 warmer sea-surface temperatures (SSTs), surface latent heat fluxes increase, and the  
532 subsequent transition to a cumulus regime is explained by the deepening-warming  
533 mechanism proposed by *Wyant and Bretherton [1997]*. Consistent with a deepening marine  
534 boundary layer from stratocumulus to trade cumulus, both optically-thin and optically-thick  
535 CTH distributions peak at higher altitudes transitioning from a stratiform to a convective  
536 type cloud regime (Fig. 11a-11c).

537 In the stratocumulus region, optically-thick CTH peaks at twice the altitude  
538 observed for optically-thin clouds (1.2 km and 0.6 km, respectively). Although cloud cover  
539 in stratocumulus regions is dominated by optically-thick cloudy profiles (Fig. 5c), optically-  
540 thin cloudy profiles comprise greater than 95% of clouds up to 4 km in size (Fig. 10).  
541 Together, this suggests that small clouds typically have lower CTH than larger, optically-  
542 thicker clouds. Noting that estimates of stratocumulus depth range from 200-300 m in the  
543 subtropics and tropics [*Wood, 2011*], and that optically-thin cloud-layers preferentially  
544 reside in the lowest 1 km with mean depth of 200 m (Fig. 12a), we suggest that incidences  
545 of optically-thin CTH below 1 km in this stratocumulus region are separate cloud entities  
546 from the optically-thicker clouds with CTH greater than 1 km. Many of these likely reflect  
547 expansive collapsed boundary layers and fog layers (note occurrence of very low-cloud  
548 bases in Fig. 11d and data along the 1:1 line in Fig. 12a) that occur periodically in regions  
549 dominated by marine stratocumulus [*Ackerman et al., 1993; Christensen and Stephens,*  
550 *2011*].

551 As would be expected given the weakening of the capping inversion and deepening  
552 of the marine layer in the transition from stratocumulus to cumulus, optically-thin and  
553 optically-thick CTH distributions are shifted upwards relative to the stratocumulus regime  
554 (cf. Figs. 11a and 11b). Cloud cover remains largely optically-thick, but in contrast to the  
555 stratocumulus region the tops of these layers tend to occur at altitudes closer to the  
556 optically-thick layers, suggesting that the optically-thin clouds in the transition region may  
557 be largely associated with thinning stratocumulus and stratus layers and remnants of  
558 dissipating stratocumuli near the top of the boundary layer.

559 In the trade wind cumulus region (Fig. 11c and 11f), the modal heights for optically-  
560 thin and optically-thick cloudy profiles are again at different levels (0.8 and 2 km  
561 respectively). Optically-thin profiles dominate the cloud cover in this regime ( $f_{\text{thin, cld}}=0.84$ ),

562 and clouds are not horizontally extensive ( $L_{50}=0.96$  km). Despite having modes at distinct  
563 levels, optically-thin and optically-thick CTH are observed above 400 m up to a height  
564 of 3 km (Fig. 11c). Peak optically-thin CBH is  $\sim 600$  m (Fig. 11f). Detrained elements from  
565 optically-thick clouds may account for a large proportion of optically-thin clouds at  
566 altitudes greater than 1 km, but the mode in optically-thin clouds below 1 km probably  
567 consists of small cumuli and associated fragments of clouds with insufficient buoyancy to  
568 grow in stature. An aerial photograph of a trade wind cloud field (Fig. 13) highlights the  
569 wide range of horizontal and vertical cloud dimensions found in trade wind regions.  
570 Notable is the prevalence of tenuous, wispy clouds with limited vertical extent, along with  
571 optically thicker clouds, although the latter are less numerous. The bimodal nature of the  
572 combined optically-thin and optically-thick cumulus CTH distributions (sum of both  
573 distributions, not shown) agrees well with high-resolution (90 m) satellite observations of  
574 trade wind cumulus CTH [Genkova *et al.*, 2007] and with a large-eddy simulation (LES)  
575 cloud cover distribution for a non-precipitating trade cumulus case [Siebesma *et al.*, 2003].  
576 This encouragingly suggests that LES is a useful tool to examine macrophysical properties  
577 of optically-thin clouds that are the dominant cloud type in the trade wind cloud regimes.

578 The optically-thin cloud-layer thickness distributions (Fig. 12 insets) are similar for  
579 all three regions. The drop-off in frequency of occurrence of layers less than 180 m may be  
580 an artifact of the increase in the nighttime optical depth detection threshold for shallow  
581 layers, i.e. the threshold is 0.03 for layers greater than 180 m, and 0.1 for shallower layers  
582 (Section 2.1). We note, however, that optically-thin cloud layer depth may be  
583 overestimated by CALIOP VFM: if a cloud partially fills a 60 m height bin yet has an  
584 optical depth greater than 0.1, the entire 60 m bin will be classified as cloudy. This high  
585 bias will affect shallow layers more than deep layers, in a relative sense. For example,  
586 a 60 m optically-thin cloud partially filling two vertically adjacent 60 m bins, will register  
587 as 120 m in depth, an error of 100%, whereas, an optically-thin cloud-layer occupying four  
588 60 m bins, and partially-filling a fifth bin by 50% will have a layer depth relative error of  
589  $\sim 11\%$ . Conservatively assuming the maximum error of 60 m in layer depth yields a 25%  
590 potential overestimate of mean optically-thin cloud-layer depth. Irregular or  
591 inhomogeneous clouds e.g., trade cumulus, will be most prone to this error. A recent LES  
592 study investigating the effects of cumulus cloud irregularity on cloud overlap quantified the  
593 effect of partially filled model levels as a function of model vertical resolution [Neggers *et*  
594 *al.*, 2011]. Employing their formulation at CALIOP VFM cloud-layer boundary detection  
595 resolution (60 m) produces a high-bias estimate of 28%, which is very similar to our own  
596 estimate.

#### 597 598 **4.0 Discussion**

599 Our findings point to a prevalence of optically-thin low clouds over the oceans, even  
600 at the scale of individual CALIOP lidar profiles. While this might be expected at the scale  
601 of typical moderate resolution passive visible remote sensors used to construct cloud  
602 climatologies, our prior expectation was that most of the optically-thin cloudy pixels from  
603 these moderate-resolution passive sensors would be partially filled by small cumulus  
604 covering only a small fraction of the pixel, and that the clouds themselves might have  
605 considerably higher optical thicknesses than the pixel as a whole. This appears not to be the  
606 case, and it prompts us to ask whether even the CALIOP FOVs are partially filled with yet  
607 smaller clouds. Given that there are many single-profile cloud entities in the CALIOP  
608 dataset (Figs. 7 and 9), it is possible a significant fraction of these profiles are partially  
609 filled.

610 The cloud within any CALIOP partially filled field of view (pFOV) could be  
 611 optically-thick or optically-thin. Based upon CALIOP data alone, we cannot diagnose  
 612 whether or not a cloudy FOV is partially cloud-filled, nor can we determine if the cloud-  
 613 filled part is optically-thick or optically-thin. Classifying a pFOV as optically-thin when the  
 614 cloud intercepted is optically-thick cloud will lead to an overestimate of  $f_{\text{thin, cld}}$ . To  
 615 investigate the frequency of pFOVs, and whether or not the clouds partially filling the lidar  
 616 FOVs are optically-thick, we employ a simple bounded cascade fractal model.

#### 617 4.1 Bounded Cascade Model Description

618 The theoretical basis for, and an in-depth description of the model is presented by  
 619 *Marshak et al.* [1994] and its application to simulating one-dimensional cloud fields is  
 620 provided by *Wood and Field* [2011]. A brief description follows. A one-dimensional  
 621 multiplicative cascade broken cloud field is generated by initially assuming a uniform value  
 622 of a cloud property; here, optical depth. The field is bisected, and each half is scaled by a  
 623 factor  $1 \pm w_i$ , with the gaining half chosen at random. This is repeated for each half, and so  
 624 on, with weighting factors ( $w_i$ ) calculated for each step  $i$  as,

$$625 \quad w_i = \frac{(1-2p)}{2^{H(i-1)}} \quad (3)$$

626 where  $p$  ( $0 < p < 0.5$ ) controls the intermittency of the resulting optical depth field, and  $H$  ( $0$   
 627  $< H < \infty$ ) is a scaling factor representing the smoothness of the field. This simple model has  
 628 displayed skill in representing scaling properties of marine stratocumulus liquid water  
 629 content [*Marshak et al.*, 1997; *Davies et al.*, 1999], and regional-global cloud size statistics  
 630 derived from the cloud optical depth field [*Wood and Field*, 2011]. Specifically, it is shown  
 631 in the latter case that the power law exponent  $\beta$  in the cloud chord length distribution is  
 632 inversely related to the value of  $H$  controlling the optical depth field, such that smooth fields  
 633 (large  $H$ ) are associated with more large clouds and low  $\beta$ , and vice versa.

634 The model is used to generate 100 optical depth realizations, each with 17 cascade  
 635 steps and therefore  $2^{18}$  individual data points of horizontal scale 2.8 m and total length 700  
 636 km. This large range of scales is chosen so that we can (a) use the model to explore sub-  
 637 FOV variability within the CALIOP FOV (we have 32 model pixels in the 90 m FOV); (b)  
 638 have sufficiently long transects to provide good statistics and be approximately consistent  
 639 with the aggregation scale for our CALIOP analysis. Consecutive FOVs are separated by  
 640 335 m to be consistent with CALIOP sampling. We correct all modeled cloud length  
 641 distributions for sampling bias associated with clouds approaching the limited transect  
 642 length using Eq. (1).

643 The optical depth threshold for CALIOP cloud detection ( $\tau_{\text{cld, thresh}}=0.03$ ) is applied to  
 644 determine if a given model FOV is cloudy. If a surface is detected by CALIOP in the  
 645 presence of a cloud, it is not known *a priori* whether the cloud is optically thin, or if an  
 646 optically-thick cloud is partially filling the FOV. To determine if CALIOP would detect the  
 647 surface for any given cloudy model FOV (i.e., to determine if the FOV would be classified  
 648 as optically thin cloud), we compare the number of surface-returned photons ( $n$ ) that would  
 649 be obtained from a model-simulated cloudy FOV to the *minimum* number of surface-  
 650 returned photons ( $n_*$ ) required to detect the surface. If the surface returns  $N$  photons without  
 651 cloud attenuation, then from Beer's law,  $n = Ne^{-2\tau}$ . This ignores multiple scattering,  
 652 which is weak for optically-thin clouds. We know that the surface is *just* detected in the  
 653 presence of a uniform cloud with  $\tau_*=3$ , and so  $n_* = Ne^{-2\tau_*}$ , which gives  $n/n_* =$   
 654  $e^{-2(\tau_*-\tau)}$ . We compute  $n/n_*$  for every model pixel (2.8 m scale) and then take the mean  
 655 over the CALIOP FOV. If this FOV mean  $\overline{n/n_*} > 1$  then the surface would be detected and

656 such a cloudy FOV is classified as optically thin to best match our CALIOP observational  
657 definition.

658 We focus our investigation on our representative tropical Pacific trade wind region  
659 (Table 3) because the trades have the smallest marine low clouds (Fig. 8a) and are therefore  
660 the most prone to the partially filled FOV problem [Jones *et al.*, 2012]. Experimentation  
661 shows that model settings best simulating CALIOP observations are  $p = 0.0005$ ,  $H = 0.07$ ,  
662 with scene mean optical depth (also a specified input variable) set to 4.5, selected as  
663 follows. Because  $H$  is closely related to  $\beta$  (Fig. 15; Wood and Field, [2011]), we vary  $H$  to  
664 match CALIOP observed  $\beta$  (Fig. 14). We also vary  $p$  to match observed CALIOP trade  
665 wind region mean  $f_{\text{cld}}$  and  $f_{\text{thin,cld}}$  as closely as possible. Finally, scene mean optical depth is  
666 adjusted such that TOA albedo  $\alpha_{\text{TOA}}$  estimated from model  $\tau_{\text{cld}}$  values, falls within the range  
667 of CERES albedo values for this region. Calculation details are presented in full in Table 4.

668

#### 669 4.1.1 Model Simulation of Observed Cloud Properties

670 The model demonstrates skill in generating realistic representations of CALIOP  
671 observed cloud properties  $f_{\text{cld}}$  and  $f_{\text{thin,cld}}$ ,  $\beta$  and TOA albedo  $\alpha^{\text{TOA}}$  (Table 4). Consistency  
672 between observed (CALIOP) and predicted cloud size distributions suggests that optically-  
673 thin cloud behavior is broadly consistent with a power law scaling of optical depth: this  
674 scaling is an underlying assumption of the model. For comparison, we also sample the  
675 model output simulating MODIS and HSRL FOV (1 km and 60 m along-track dimensions,  
676 respectively). HSRL-simulated model output is uses the same  $\tau_{\text{cld,thresh}}=0.03$  and  $n/n_* > 1$   
677 thresholds used for CALIOP. We set MODIS  $\tau_{\text{cld,thresh}}$  to be 0.4 [Ackerman *et al.*, 2008],  
678 with  $\tau_{\text{thin,thresh}}$  as for CALIOP. Perhaps unexpectedly, model simulated MODIS  $f_{\text{cld}}$  is 1.5  
679 times *lower* than CALIOP and HSRL simulated  $f_{\text{cld}}$  (Table 4). Varying MODIS  $\tau_{\text{cld,thresh}}$   
680 within reasonable limits (0.2 to 0.6) produces model simulated MODIS  $f_{\text{cld}}$  values of 0.20  
681 and 0.14 respectively (i.e. still lower than the CALIOP value of 0.27). We therefore suggest  
682 that overestimation of cloud cover using passive visible imagers (e.g., Zhao and Di  
683 Girolamo, [2006]) may not necessarily always occur. Rather, the bias depends sensitively  
684 upon cloud length and the reflectance threshold applied.

685 As instrument resolution increases (FOV size decreases), each simulated size  
686 distribution is shifted to smaller sizes (Fig. 14). We note that the model-simulated HSRL  
687 cloud length distribution, since it is derived from a bounded cascade, does not reproduce the  
688 observed gradual roll-off to a flatter size distribution at cloud lengths smaller than  $\sim 500$  m  
689 seen in the tropical western Atlantic trade wind region (Fig. 9). Interestingly, simulated  
690 cloud length distributions for all instrument FOVs have similar  $\beta$ , a result also noted by  
691 Wood and Field [2011], whereby the cloud length distribution is not strongly sensitive to  
692 sensor resolution when the domain to FOV ratio approaches 1000 (i.e. when there is a  
693 sufficient range of scales). This ratio is much greater than 1000 for model clouds sampled  
694 like CALIOP and HSRL, and is  $\sim 1000$  in the case of MODIS, and so we would not expect  $\beta$   
695 to be very different across the sensors.

696 The HSRL has a pulse repetition frequency of 200 Hz, and so prior to averaging, the  
697 data have less than 1 m along-track sampling resolution for an aircraft flight speed  
698 of  $120 \text{ m s}^{-1}$ . Although onboard averaging means that we do not actually have observations  
699 at this resolution from HSRL, it is interesting to see whether size distributions using higher  
700 resolution cloud masks would be shifted to even smaller sizes. The model results suggest  
701 that observations at CALIOP and HSRL resolution both miss finer scale cloud length  
702 variations, and overestimate cloud lengths (Fig. 14). Additionally, even cloud cover  
703 estimates made at CALIOP and HSRL resolution may overestimate  $f_{\text{cld}}$  by almost 25% in

704 trade wind regions (Table 4). Although these model results require verification, especially  
705 the model behavior at small scales, we note that a four-fold overestimate in  $f_{\text{cld}}$  was  
706 estimated from high-resolution (15 m) trade Cu observations when sensor resolution was  
707 degraded to 1 km while cloud detection threshold was held constant [Dey *et al.*, 2008].  
708 Indeed, this suggests that the very concept of cloud cover may be so sensitive to sensor  
709 resolution and detection sensitivity in the trades that it serves no useful function in these  
710 regions. For lower resolution sensors (e.g., MODIS) such small-scale clouds likely produce  
711 cloud-contamination in clear FOVs [Koren *et al.*, 2007], biasing clear-sky retrievals.

712

#### 713 4.2 Partially Filled Fields of View

714 As previously mentioned, a CALIOP FOV may be misidentified as optically-thin if  
715 partially filled by an optically-thick cloud. The model is employed to identify simulated  
716 CALIOP cloudy FOVs that are partially cloud filled (pFOV), i.e., not all 32 optical depth  
717 values comprising the FOV exceed the threshold for cloud detection ( $\tau_{\text{cld,thresh}}$ ). As  
718 discussed in Section 4.1, if the FOV-mean ratio of surface-backscattered photons detected  
719 by the sensor ( $n$ ) to the minimum number required for surface detection ( $n_*$ ) is greater than  
720 one, the cloudy FOV is classified as optically thin. The mean  $n/n_*$  for the entire pFOV is  
721 calculated, along with the mean optical depth of the cloud-only segment. For a given pFOV  
722 with  $n/n_* > 1$  and mean optical depth of the cloud-only segment greater than 3, the pFOV  
723 is therefore misclassified as an optically-thin cloud whereas it is really a partially-filled  
724 FOV containing an optically thick cloud. The frequency of such misclassification is  
725 estimated. Results are binned by cloud length, and presented in Fig. 15.

726 The fraction of FOVs that are partially cloud filled is high, especially for  
727 smaller clouds where it exceeds 0.8. This fraction decreases to around 0.3 for larger clouds  
728 (Fig. 15, circles). CALIOP-observed cloud entities consisting of a single profile are  
729 therefore almost all likely to be broken clouds. The bin median pFOV cloud fraction, i.e.,  
730 fraction of a pFOV that is cloud, increases with cloud length, ranging from 0.50 to 0.85  
731 (Fig. 15, squares), consistent with fewer cloud edges being observed as cloud length  
732 increases. Despite the prevalence of partially-filled FOVs, the FOVs classified as optically  
733 thin typically do contain mainly optically-thin cloud (Fig. 15, diamonds). Summing the  
734 fraction of pFOV that are misclassified for each size bin, and weighting this value by the  
735 fraction of optically-thin FOVs that are partially filled, it can be concluded that the potential  
736 overestimate of  $f_{\text{thin,cld}}$  is at most 6% for this cloud regime. In addition, the effect of sensor  
737 resolution on  $f_{\text{cld}}$  can be quantified from the model output. The overestimate in  $f_{\text{cld}}$  caused  
738 by assuming that partially cloud-filled FOVs are 100% cloud filled is estimated at  $\sim 0.06$  in  
739 absolute terms. This value is similar to the ratio of model generated CALIOP and HSRL  
740 Full Res  $f_{\text{cld}}$  ( $0.27/0.22 = 1.23$ ), with the latter representing “true”  $f_{\text{cld}}$ .

741 Based upon the model output, which is designed to simulate typical trade wind  
742 cumulus conditions, it is concluded that the majority of cloudy CALIOP profiles are indeed  
743 partially cloud filled. Rapid fluctuations of cloud particle concentration on scales  $\sim 10$  m  
744 have been measured in situ [Pawlowska *et al.*, 1997], suggesting that this conclusion is  
745 physically realistic. These results have implications for the minimum sensor resolution  
746 required for cloud observations in regions of broken cloud, as simulated by the model. A  
747 recent study suggests that a resolution higher than 45 m is required to accurately determine  
748 the cloud cover to within 0.01 for trade cumulus regimes [Jones *et al.*, 2012]. Our study  
749 suggests that the CALIOP FOV size of 90 m leads to an overestimate of some 0.06 in cloud  
750 cover in regions of trade cumulus. This is consistent with Jones *et al.* [2012], whose results

751 (their Fig. 1) suggest a median overestimation of 0.05 for a sensor with a resolution  
752 of  $\sim 100$  m.

753

#### 754 4.3 Additional Applications of Cascade Model Output

##### 755 4.3.1 Are Optically-thin Clouds Adiabatic?

756 We analyze the model output to investigate the extent to which optically-thin clouds  
757 might be considered adiabatic, i.e., liquid water path ( $LWP$ ) values determined by  
758 temperature, pressure and cloud thickness alone. Liquid water path for an adiabatic cloud is  
759 proportional to cloud-layer depth ( $h$ ) squared, the adiabatic rate of increase of liquid water  
760 with respect to height (assumed to be  $2 \times 10^{-3} \text{ g m}^{-3} \text{ km}^{-1}$ ) [Albrecht *et al.*, 1990], and the  
761 adiabatic fraction ( $f_{\text{ad}}$ ), is defined as the ratio of observed  $LWP$  to adiabatic  $LWP$ , which  
762 here we assume to be 1. Using model simulated optically-thin cloud optical depth for a  
763 trade wind region as before, and assuming a cloud droplet effective radius ( $r_e$ ) of  $10 \mu\text{m}$ , we  
764 estimate  $LWP = 2/3\tau\rho_L r_e$  [Stephens, 1994], where  $\rho_L$  is liquid water density, and derive  
765 corresponding adiabatic cloud-layer depth,  $h$ . The simulated adiabatic layer depth  
766 distribution (not shown) indicates that depths are on the order of a few 10s of meters, with  
767 90% of the values being less than 90 m, much thinner than observed. In addition, if this  
768  $LWP$  is evenly distributed over a 60 m layer, we estimate that  $f_{\text{ad}} \sim 0.05$ , and ranges from  
769 0.02 to 0.09 when  $r_e$  is halved and doubled, respectively. Alternatively, if  $\tau_{\text{cld}}$  is 0.1, layer  
770 depth corresponding to an adiabatic  $LWP$  is 25 m. Therefore, we conclude that the  
771 optically-thin cloud-layers are not adiabatic, but are more likely to be diluted cloud  
772 fragments, and their  $LWP$  should not be estimated adiabatically. This is consistent with  
773 effects of entrainment on marine layer clouds (e.g., McFarlane and Grabowski [2007]).  
774 Stochastically-controlled entrainment events essentially dictate the properties of shallow  
775 trade cumulus clouds [Romps and Kuang, 2010].

776

##### 777 4.3.2 Optically-thin Cloud Contribution to TOA Cloud Albedo

778 We have shown that optically-thin clouds comprise almost one-half of all marine  
779 low clouds over non-polar oceans. To investigate the magnitude of their contribution to  
780 low-cloud  $\alpha^{\text{TOA}}$  in trade wind regions at CALIOP FOV resolution, model simulated optical  
781 depth is converted to albedo using a simple approximation,  $\alpha^{\text{cld}} = \tau_{\text{cld}} / (\tau_{\text{cld}} + 7.7)$  [Lacis and  
782 Hansen, 1974]), and  $\alpha^{\text{TOA}}$  is estimated for optically-thin clouds and all low clouds (Table 4).  
783 The ratio of these values is expressed per unit of low-cloud cover and per unit of optically-  
784 thin cloud cover, and the calculation is repeated at model full-resolution to investigate how  
785 optically-thick clouds classified as optically-thin clouds in pFOVs, affect the results.

786 In addition to the marine cumulus region described in Table 3, two other trade  
787 cumulus regions (of equal size) were examined; one in the tropical western Pacific and one  
788 in the tropical western Atlantic. For trade wind regions it can be concluded that with  
789 observed  $f_{\text{thin,cld}} > 0.90$ , and  $f_{\text{cld}} \sim 0.20 - 0.30$ , optically-thin low clouds can contribute 55-  
790 60% to the total low-cloud albedo. However, the range of values obtained from the full-  
791 resolution model simulation indicates that the true contribution of optically-thin clouds to  
792 low-cloud albedo is more than a factor of two smaller at 23% - 26%.

793

794

795

## 796 5.0 Summary and Concluding Remarks

797 Analysis of two years of CALIOP nighttime data indicates that 45% of marine low  
798 clouds with tops below 3 km, over the non-polar oceans, are classified as optically thin such  
799 that they do not fully attenuate the CALIOP lidar signal. Few observational data, other than  
800 this work, are available on the physical nature of these clouds. Key results from our  
801 analysis are as follows:

- 802 1) Over the non-polar oceans, optically-thin clouds comprise 0.25 – 0.75 of marine low  
803 clouds with cloud top height below 3 km, with a mean of 0.45;
- 804 2) The optically-thin fraction of marine low cloud varies inversely with marine low-  
805 cloud cover, and reaches a maximum ( $> 0.80$ ) in trade wind regions;
- 806 3) Although the optically-thin fraction of low clouds peaks in trade wind regions, the  
807 absolute frequency of occurrence does not strongly vary over the non-polar oceans;
- 808 4) In trade wind regions, according to our analysis of CALIOP observations, clouds  
809 smaller than 2 km contribute over 50% of the cloud cover;
- 810 5) Optically-thin marine low clouds are predominantly small clouds. The cloud length  
811 distribution of all clouds explains three-quarters of the geographical variance in the  
812 optically-thin fraction of marine low clouds;
- 813 6) Over the non-polar oceans, optically-thin low clouds have a mean cloud top height  
814 of 1.2 km and a mean thickness of 0.25 km. However, optically-thin cloud-layer  
815 depth may be overestimated by 25% due to cloud partially filling the 60 m vertical  
816 resolution CALIOP bins;
- 817 7) Collocated HSRL and CALIOP data for the tropical western Atlantic trade wind  
818 region reveal that many clouds are smaller than can be resolved at CALIOP  
819 resolution. This suggests that clouds, especially optically-thin low clouds, may exist  
820 globally on finer scales than reported herein;
- 821 8) A simple bounded cascade fractal model is used to represent a one-dimensional  
822 optical depth distribution for conditions representative of the trade wind regions.  
823 Analysis of model output, in conjunction with observations, suggests that:
  - 824 a) Optically-thin low-cloud properties are consistent with a power law scaling  
825 of optical depth;
  - 826 b) CALIOP fields of view partially filled with optically-thick low cloud  
827 produce a potential 6% overestimate in optically-thin fraction of marine low-  
828 cloud values in trade cumulus regions;
  - 829 c) Optically-thin cloud layer thicknesses are not adiabatic with greater than  
830 90% of simulated optically-thin adiabatic cloud depths less than 90 m,  
831 whereas observations indicate much greater thicknesses;
  - 832 d) In trade cumulus regions, we estimate that optically-thin low clouds  
833 contribute almost a quarter of the albedo contributed by low clouds, but a  
834 more sophisticated quantification using radiative transfer modeling is  
835 required to confirm this.

836 Remote sensing of clouds provides global information on cloud macrophysical and  
837 microphysical properties, both of which are important for radiative transfer calculations and  
838 verification of model representation of clouds. Differences in cloud properties retrieved by  
839 ISCCP and MODIS [*Pincus et al.*, 2010] and ISCCP, MODIS and MISR [*Marchand et al.*,  
840 2010] highlight the difficulties inherent in passive remote sensing of low clouds. *Pincus et*  
841 *al.* [2012] reported that ISSCP and MODIS views of clouds are consistent for all but the

842 optically-thinnest clouds, while sub-pixel scale and broken low-level clouds were identified  
843 by *Marchand et al.* [2010] as one of the reasons for discrepancies in cloud properties.

844 It is known that low clouds are the dominant contributor to the global albedo  
845 [*Hartmann et al.*, 1992]. Results from our work indicate that almost half of low clouds over  
846 the non-polar ocean are optically thin. Additionally, in trade wind regions, 50% of the low-  
847 cloud cover is comprised of clouds less than 2 km in length. Because such a large fraction  
848 of cloud is optically thin, even with CALIOP's 90 m footprint size, it is doubtful that the  
849 passive sensors are actually determining the coverage of the optically-thick or optically-thin  
850 clouds correctly. It is difficult therefore, to conclude that we yet have a precise  
851 understanding of the shape of the true cloud optical depth probability distribution function  
852 on a global scale, and especially in trade wind regions.

853 Although CALIOP data do not provide vertically-resolved high resolution off-track  
854 spatial context for the cloud fields sampled, such data in the along-track direction afford  
855 new insights on the nature of optically-thin clouds. Suggestions for design of future  
856 combined active and passive remote sensors would include an array of lasers providing a  
857 swath of vertically-resolved data with, ideally, horizontal resolution of a few tens of meters  
858 or less, in the across-track and along-track directions, and collocated visible imagery with  
859 the same horizontal resolution. Such a configuration would provide collocated profiles of  
860 cloud properties, and high-resolution TOA reflectance data to better observe these optically-  
861 thin clouds, that comprise almost half of all marine low clouds.

862  
863

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869

870 **References**

- 871 Ackerman, A.S., O.B. Toon, and P.V. Hobbs (1993) Dissipation of marine stratiform clouds  
872 and collapse of the marine boundary layer due to the depletion of cloud condensation  
873 nuclei by clouds, *Science*, 262, 226-229, doi:10.1126/science.262.5131.226.
- 874 Ackerman, S., R. E. Holz, R. Frey, E. W. Eloranta, B. C. Maddux, and M. McGill (2008),  
875 Cloud Detection with MODIS. Part II: Validation, *J. Atmos. Oceanic Technol.*, 25,  
876 1073-1086.
- 877 Albrecht, B. A., C. W. Fairall, D. W. Thomson, A. B. White, J. B. Snider, and W. H.  
878 Schubert (1990), Surface-based remote sensing of the observed and the Adiabatic liquid  
879 water content of stratocumulus clouds, *Geophys. Res. Lett.*, 17, 89-92,  
880 doi:10.1029/GL017i001p00089.
- 881 Albrecht, B. A., M. P. Jensen, and W. J. Syrett (1995), Marine boundary layer structure and  
882 fractional cloudiness, *J. Geophys. Res.*, 100, D7, 14,209-14,222.
- 883 Bender, F. A-M., R. J. Charlson, A. M-L. Ekman, and L. V. Leahy (2011), Quantification of  
884 monthly mean, regional scale albedo of marine stratiform clouds in satellite  
885 observations and GCMs, *J. Appl. Meteor. Clim.*, doi: 10.1175/JAMC-D-11-049.1
- 886 Benner, T. C., and J. A. Curry (1998), Characteristics of small tropical cumulus clouds and  
887 their impact on the environment, *J. Geophys. Res.*, 103, D22, 25,752-28,767.
- 888 Bucholtz, A. (1995), Rayleigh-scattering calculations for the terrestrial atmosphere, *Appl.*  
889 *Opt.*, 34, 2,765-2,773.
- 890 Charlson, R. J., A. S. Ackerman, F. A-M. Bender, T. L. Anderson, and Z. Liu (2007), On  
891 the climate forcing consequences of the albedo continuum between cloudy and clear air,  
892 *Tellus*, 59, 715-727, doi:10.1111/j.1600-0889.200700297.x.
- 893 Davies, A. B., A. Marshak, H. Gerber, and W. J. Wiscombe (1999), Horizontal structure of  
894 marine boundary layer clouds from centimeter to kilometer scales, *J. Geophys. Res.*,  
895 104, D6, 6,123-6,144.
- 896 Dey, S., L. Di Girolamo, and G. Zhao (2008), Scale effect on statistics of the macrophysical  
897 properties of trade wind cumuli over the tropical western Atlantic during RICO, *J.*  
898 *Geophys. Res.*, 113, D24214, doi: 10.1029/2008JD010295.
- 899 Efron, B., and G. Gong (1983), A Leisurely Look at the Bootstrap, the Jackknife, and  
900 Cross-Validation, *The American Statistician*, 37, 36-48.
- 901 Genkova, I., G. Seiz, P. Zuidema, G. Zhao, and L. Di Girolamo (2007), Cloud top height  
902 comparisons from ASTER, MISR, and MODIS for trade wind cumuli, *Remote Sens.*  
903 *Environ.*, 107, 211-222.
- 904 Garrett, T. J., and C. F. Zhao (2006), Increased Arctic cloud longwave emissivity associated  
905 with pollution from mid-latitudes, *Nature*, 440, 787-789.
- 906 Hair, J. W., C. A. Hostetler, A. L. Cook, D. B. Harper, R. A. Ferrare, T. L. Mack, W.  
907 Welch, L. R. Izquierdo, and F. E. Hovis (2008), Airborne High Spectral Resolution  
908 Lidar for profiling aerosol optical properties, *Appl. Opt.*, 47, 6,734-6,752.
- 909 Harrison, E. F., P. Minnis, B. R. Barkstrom, V. Ramanathan, R. D. Cess, and G. G. Gibson  
910 (1990), Seasonal variation of cloud radiative forcing derived from the Earth radiation  
911 budget experiment, *J. Geophys. Res.*, 95, 18687-18703.
- 912 Hartmann, D. L., M. E. Ockert-Bell, and M. L. Michelsen (1992), The Effect of Cloud Type  
913 on Earth's Energy Balance: Global Analysis, *J. Climate*, 5, 1281-1304, doi:  
914 10.1175/1520-0442(1992)005.
- 915 Hozumi, K., T. Harimaya, and C. Magono (1982), The Size Distribution of Cumulus Clouds  
916 as a Function of Cloud Amount, *J. Meteor. Soc. Japan*, 60, 691-699.

917 Hunt, W.H., D. M. Winker, M. A. Vaughan, K. A. Powell, P. L. Lucker, C. Weimer (2009),  
918 CALIPSO Lidar Description and Performance Assessment, *J. Atmos. Oceanic Technol.*,  
919 26, 1214–1228, doi: 10.1175/2009JTECHA1223.1.

920 Jones, A. L., L. Di Girolamo, and G. Zhao (2012), Reducing the resolution bias in cloud  
921 fraction from satellite derived clear-conservative cloud masks, *J. Geophys. Res.*, 117,  
922 D12201, doi:10.1029/2011JD017195, in press.

923 Kacenelenbogen, M., M. A. Vaughan, J. Redemann, R. M. Hoff, R. R. Rogers, R. A.  
924 Ferrare, P. B. Russell, C. A. Hostetler, J. W. Hair, and B. N. Holben (2011), An  
925 accuracy assessment of the CALIOP/CALIPSO version 2/version 3 daytime aerosol  
926 extinction product based on a detailed multi-sensor, multi-platform case study, *Atmos.*  
927 *Chem. Phys.*, 11, 3,981-4,000, doi:10.5194/acp-11-3981-2011.

928 Kay, J. E., Hillman, B., Klein, S., Zhang, Y., Medeiros, B., Gettelman, G., Pincus, R.,  
929 Eaton, B., Boyle, J., Marchand, R. and T. Ackerman (2012), Exposing global cloud  
930 biases in the Community Atmosphere Model (CAM) using satellite observations and  
931 their corresponding instrument simulators, *J. Climate*, doi:10.1175/JCLI-D-11-00469.1,  
932 in press.

933 Klein, S. A., and D. L. Hartmann (1993), The Seasonal Cycle of Low Stratiform Clouds, *J.*  
934 *Climate*, 6, 1,587-1,606.

935 Koren, I., L. A. Remer, Y. J. Kaufman, and Y. Rudich (2007), On the twilight zone between  
936 clouds and aerosols, *Geophys. Res. Lett.*, 34, L08805, doi:10.1029/2007GL029253.

937 Koren, I., L. Oreopoulos, G. Feingold, L. A. Remer, and O. Attaratz (2008), How small is a  
938 small cloud?, *Atmos. Chem. Phys.*, 8, 3,855-3,864.

939 Koren, I., G. Feingold, H. Jiang, and Orit Altaratz (2009), Aerosol effects on the inter-cloud  
940 region of a small cumulus cloud field, *Geophys. Res. Lett.*, 36, L14805,  
941 doi:10.1029/2009GL037424.

942 Lacis, A. A., and J. E. Hansen (1974), A Parameterization for the Absorption of Solar  
943 Radiation in the Earth's Atmosphere, *J. Atmos. Sci.*, 31, 118-133.

944 Marshak, A., A. Davis, R. Cahalan, and W. Wiscombe (1994), Bounded cascade models as  
945 nonstationary multifractals, *Phys. Rev. E*, 49(1), 55-69.

946 Marshak, A., A. Davis, W. Wiscombe, and R. Cahalan (1997), Scale invariance in liquid  
947 water distributions in marine stratocumulus. Part II: multifractal properties and  
948 intermittency issues, *J. Atmos. Sci.*, 54, 1,423-1,444.

949 Masonis, S. J., T. L. Anderson, and D. S. Covert (2003), A Study of the Extinction-to-  
950 Backscatter Ratio of Marine Aerosol during the Shoreline Environment Aerosol Study,  
951 *J. Atmos. Oceanic Technol.*, 20, 1,388-1,402.

952 McFarlane, S. A., and W. W. Grabowski (2007), Optical properties of shallow tropical  
953 cumuli derived from ARM ground-based remote sensing, *Geophys. Res. Lett.*, 34,  
954 L06808, doi:10.1029/2006GL028767.

955 McFarquhar, G. M., S. Platnick, L. Di Girolamo, H. Wang, G. Wind, and G. Zhao (2004),  
956 Trade wind cumuli statistics in clean and polluted air over the Indian Ocean from in situ  
957 and remote sensing measurements, *Geophys. Res. Lett.*, 31, L21105,  
958 doi:10.1029/2004GL020412.

959 Medeiros, B., L. Nuijens, C. Antoniazzi, and B. Stevens (2010), Low-latitude boundary  
960 layer clouds as seen by CALIPSO, *J. Geophys. Res.*, 115, D23207,  
961 doi:10.1029/2010JD014437.

962 Neggers, R. A. J., T. Heus, and A. P. Siebesma (2011), Overlap statistics of cumuliform  
963 boundary-layer cloud fields in large-eddy simulations, *J. Geophys. Res.*, 116, D21202,  
964 doi:10.1029/2011JD015650.

965 Pawlowska, H., J. L. Brenguier, and G. Salut (1997), Optimal Nonlinear Estimation for  
966 Cloud Particle measurements, *J. Atmos. Oceanic Technol.*, *14*, 88-104.

967 Pincus, R. S., S. Platnick, S. Ackerman, R. Hemler, and R. Hofmann (2012), Reconciling  
968 simulated and observed views of clouds: MODIS, ISCCP, and the limits of instrument  
969 simulators, *J. Climate*, doi:10.1175/JCLI-D-11-00267.1, in press.

970 Plank, V. G. (1969), The size distribution of cumulus clouds in representative Florida  
971 populations, *J. Appl. Meteorol.*, *8*, 46– 67.

972 Redemann, J., Q. Zhang, P. B. Russell, and J. M. Livingston (2009), Case studies of aerosol  
973 remote sensing in the vicinity of clouds, *J. Geophys. Res.*, *114*, D06209,  
974 doi:10.1029/2008JD010774.

975 Rogers, R. R., C. A. Hostetler, J. W. Hair, R. A. Ferrare, Z. Liu, M. D. Obland, D. B.  
976 Harper, A. L. Cook, K. A. Powell, M. A. Vaughan, and D. M. Winker (2011),  
977 Assessment of the CALIPSO Lidar 532 nm attenuated backscatter calibration using the  
978 NASA LaRC airborne High Spectral Resolution Lidar, *Atmos. Chem. Phys.*, *11*, 1,295-  
979 1,311, doi:10.5194/acp-11-1295-2011.

980 Romps, D. M., and Z. Kuang (2010), Nature versus nurture in shallow convection, *J. Atmos.*  
981 *Sci.*, *67*, 1655-1666.

982 Rossow, W. B., and R. A. Schiffer (1999), Advances in Understanding Clouds From  
983 ISCCP, *Bull. Am. Meteorol. Soc.*, *80(11)*, 2,261-2,287.

984 Rossow, W.B., C. Delo, and B. Cairns (2002), Implications of the observed mesoscale  
985 variations of clouds for the Earth's radiation budget, *J. Climate*, *15*, 557-585.

986 Sengupta, S. K., R. M. Welch, M. S. Navar, T. A. Berendes, and D. W. Chen (1990),  
987 Cumulus Cloud Field Morphology and Spatial Patterns Derived from High Spatial  
988 Resolution Landsat Imagery, *J. App. Meteor.*, *29*, 1,245-1,267.

989 Siebesma, A. P., C. S. Bretherton, A. Brown, A. Chlond, J. Cuxart, P. G. Duynkerke, H.  
990 Jiang, M. Khairoutdinov, D. L., Lewellen, C-H. Moeng, E. Sanchez, B. Stevens, and D.  
991 E. Stevens (2003), A Large Eddy Simulation Intercomparison Study of Shallow  
992 Cumulus Convection, *J. Atmos. Sci.*, *60(10)*, 1,201-1,219.

993 Smirnov, A., B. N. Holben, Y. J. Kaufman, O. Dubovik, T. F. Eck, I. Slutsker, C. Pietras,  
994 and R. N. Halthore (2002), Optical Properties of Atmospheric Aerosol in Maritime  
995 Environments, *J. Atmos. Sci.*, *59*, 501-523.

996 Stephens. G. L. (1994), Remote Sensing of the Lower Atmosphere. Oxford University  
997 Press, 523 pp.

998 Su, W., G. L. Schuster, N. G. Loeb, R. R. Rogers, R. A. Ferrare, C. A. Hostetler, J. W. Hair,  
999 and Michael D. Obland (2008), Aerosol and cloud interactions observed from high  
1000 spectral resolution lidar data, *J. Geophys. Res.*, *113*, D24202,  
1001 doi:10.1029/2008JD010588.

1002 Tackett, J. L., and L. Di Girolamo (2009), Enhanced aerosol backscatter adjacent to tropical  
1003 trade wind clouds revealed by satellite-based lidar, *Geophys. Res. Lett.*, *36*, L14804,  
1004 doi:10.1029/2009GL039264.

1005 Turner, D. D, A. M. Vogelmann, R. T. Austin, J. C. Barnard, K. Dacy-Pereira, J. C. Chiu, S.  
1006 A. Clough, C. Flynn, M. M. Kahiyyer, J. Liljegren, K. Johnson, B. Lin, C. Long, A.  
1007 Marshak, S. Y. Matrosov, S. A. McFarlane, M. Miller, Q. Min, P. Minnis, W. O'Hirok,  
1008 Z. Want, and W. Wiscombe (2007), Thin Liquid Water Clouds – Their importance and  
1009 our challenge, *Bull. Am. Meteorol. Soc.*, *88(2)*, 177-190, doi:10.1175/BAMS-88-2-177.

1010 Twohy, C. H., J. A. Coakley Jr., and W. R. Tahnk (2009), Effect of changes in relative  
1011 humidity on aerosol scattering near clouds, *J. Geophys. Res.*, *114*, D05205,  
1012 doi:1029/2008JD010991.

1013 Várnai, T, and A. Marshak (2009), MODIS observations of enhanced clear sky reflectance  
1014 near clouds, *Geophys. Res. Lett.*, *36*, L06807, doi:10.1029/2008GL037089.

1015 Várnai, T, and A. Marshak (2011), Global CALIPSO Observations of Aerosol Changes  
1016 Near Clouds, *IEEE Geosci. Remote Sens. Lett.*, *8(1)*, 19-23.

1017 Vaughan, M. A., D. M. Winker, and K. A. Powell (2005), Part 2: Feature Detection and  
1018 Layer Properties Algorithms. CALIOP Algorithm Theoretical Basis Document PC-  
1019 SCI-202.01, 87 pp. [Available online at [http://www-  
1020 calipso.larc.nasa.gov/resources/pdfs/PC-SCI-202\\_Part2\\_rev1x01.pdf](http://www-calipso.larc.nasa.gov/resources/pdfs/PC-SCI-202_Part2_rev1x01.pdf)].

1021 Vaughan, M. A., K. A. Powell, R. E. Kuehn, S. A. Young, D. M. Winker, C. A. Hostetler,  
1022 W. H. Hunt, Z. Liu, M. J. McGill, and B. J. Getzewich (2009), Fully Automated  
1023 Detection of Cloud and Aerosol Layers in the CALIPSO Lidar Measurements, *J. Atmos.*  
1024 *Oceanic Technol.*, *26*, 2,034-2,050, doi: 10.1175/2009JTECHA1228.1.

1025 Wen, G., A. Marshak, R. E. Cahalan, L. A. Remer, and R. G. Gleidman (2007), 3-D  
1026 aerosol-cloud radiative interaction observed in collocated MODIS and ASTER images  
1027 of cumulus cloud fields, *J. Geophys. Res.*, *112*, D13204, doi:10.1029/2006JD008267.

1028 Wielicki, B. A., and R. M. Welch (1986), Cumulus Cloud Properties Derived Using Landsat  
1029 Satellite Data, *J. Clim. Appl. Meteorol.*, *25*, 261-276.

1030 Winker, D. M., M. A. Vaughan, A. Omar, Y. Hu, K. A. Powell, Z. Liu, W. H. Hunt, and S.  
1031 A. Young (2009), Overview of the CALIPSO Mission and CALIOP Data Processing  
1032 Algorithms, *J. Atmos. Oceanic Technol.*, *26*, 2310–2323, doi:  
1033 10.1175/2009JTECHA1281.1.

1034 Winker, D. M., J. Pelon, J. A. Coakley, S. A. Ackerman, R. J. Charlson, P. R. Colarco, P.  
1035 Flamant, Q. Fu, R. M. THE CALIPSO MISSION A Global 3D View of Aerosols and  
1036 Clouds, *Bull. Am. Meteorol. Soc.*, *91*, 1,211-1,229, doi:10.1175/2010BAMS3009.1.

1037 Wood, R., and J. P. Taylor (2001), Liquid water path variability in unbroken marine  
1038 stratocumulus cloud, *Quart. J. Roy. Meteor. Soc.*, *127*, 2,635-2,662.

1039 Wood, R., and C. S. Bretherton (2004), Boundary layer Depth, Entrainment, and  
1040 Decoupling in the Cloud-Capped Subtropical and Tropical marine Boundary Layer, *J.*  
1041 *Climate*, *17*, 3,576-3,588.

1042 Wood, R., and D. L. Hartmann (2006), Spatial Variability of Liquid Water Path in Marine  
1043 Low Cloud: The Importance of Mesoscale Cellular Convection, *J. Climate*, *19*, 1,748-  
1044 1,764. Wood, R., and P. R. Field (2011), The Distribution of Cloud Horizontal Sizes, *J.*  
1045 *Climate*, *24*, 4,800–4,816, doi: 10.1175/2011JCLI4056.1.

1046 Wood, R. (2011) Stratocumulus Clouds, *Mon. Wea. Rev.*, doi:10.1175/MWR-D-11-  
1047 00121.1, in press.

1048 Wyant, M. C., C. S. Bretherton, H. A. Rand, and D. E. Stevens (1997) Numerical  
1049 Simulations and a Conceptual Model of the Stratocumulus to Trade Cumulus Transition,  
1050 *J. Atmos. Sci.*, *54*, 168–192.

1051 Zhang, J., and J. S. Reid (2005), An analysis of potential cloud artifacts in MODIS over  
1052 ocean aerosol optical thickness products, *Geophys. Res. Lett.*, *32*, L15803,  
1053 doi:10.1029/2005GL023254.

1054 Zhao, G., and L. Di Girolamo (2006), Cloud fraction errors for trade wind cumuli from  
1055 EOS-Terra instruments, *Geophys. Res. Lett.*, *33*, L20802, doi:10.1029/2006GL028088.

1056 Zhao, G., and L. Di Girolamo (2007), Statistics on the macrophysical properties of trade  
1057 wind cumuli over the tropical western Atlantic, *J. Geophys. Res.*, *112*, D10204,  
1058 doi:10.1029/2006JD007371.

1059 Zuidema, P., D. Leon, A. Pazmany, and M. Cadeddu (2012), Aircraft millimeter-wave  
 1060 passive sensing of cloud liquid water and water vapor during VOCALS-REx, *Atmos.*  
 1061 *Chem. Phys.*, 12, 355-369, doi:10.5194/acp-12-355-2012.

1062 **TABLES**

1063 **Table 1.** CALIOP – HSRL matched dataset description. <sup>a</sup>

<b>Date</b>	<b>Location<sup>b</sup></b>	<b>Overpass Time<sup>c</sup></b>	<b>Cloud Type<sup>d</sup></b>
26 Jan 2007	36.88°N, 75.52°W	07:11	-
23 May 2007	32.41°N, 77.07°W	18:32	Cumulus
26 Jan 2008	12.43°N, 60.26°W	17:36	Cumulus
11 Aug 2010	33.59N, 65.09W	17:43	Cumulus/Stratocumulus
18 Aug 2010	22.38°N, 63.95°W	17:48	Cumulus
22 Aug 2010	13.71°N, 69.04°W	06:39	-
24 Aug 2010	21.76°N, 64.09°W	06:15	-
26 Aug 2010	20.67N, 61.25W	06:02	-

1064 <sup>a</sup> Transect lengths range from 34 km to 215 km, with an 8-day mean of 126 km.

1065 <sup>b</sup> Segment mid-point (latitude, longitude).

1066 <sup>c</sup> Times are UTC.

1067 <sup>d</sup> Cloud type information from inspection of MODIS Aqua L1B Granule Images (daytime only),  
 1068 ([http://modis-atmos.gsfc.nasa.gov/IMAGES/index\\_myd021km.html](http://modis-atmos.gsfc.nasa.gov/IMAGES/index_myd021km.html)).

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1070

1071

1072 **Table 2.** Cloud cover<sup>a</sup> definitions for this work.

<b>Parameter<sup>b</sup></b>	<b>Definition</b>	<b>Equation<sup>c</sup></b>
Low-cloud cover ( $f_{cld}$ )	The fraction of all high-cloud-screened clear and cloud profiles within a gridbox that are clouds.	$\frac{\sum N_{cloud}}{\sum(N_{clear} + N_{cloud})}$
Optically-thin low-cloud cover ( $f_{thin}$ )	The fraction of all high-cloud-screened clear and cloud profiles within a gridbox that are optically-thin.	$\frac{\sum N_{thin}}{\sum(N_{clear} + N_{cloud})}$
Optically-thin fraction of low cloud ( $f_{thin,cld}$ ) i.e. $f_{thin}/f_{cld} = f_{thin,cld}$	The fraction of high-cloud-screened cloud profiles within a gridbox that are optically-thin.	$\frac{\sum N_{thin}}{\sum N_{cloud}}$

1073 <sup>a</sup> Unless otherwise stated, all gridbox cloud cover values are two-year averages of monthly calculated values.

1074 <sup>b</sup> All parameters are calculated at full-resolution.

1075 <sup>c</sup>  $N$  is the number of full-resolution profiles within a gridbox that are: cloudy ( $N_{cloud}$ ), cloudy and optically-thin  
1076 ( $N_{thin}$ ), clear ( $N_{clear}$ ).

1077

1078 **Table 3.** Stratocumulus (Sc), Cumulus (Cu), and transition from Sc to Cu regime  
 1079 (Sc-Cu), region details: location, low-cloud cover ( $f_{cld}$ ), optically-thin fraction of low  
 1080 cloud ( $f_{cld,thin}$ ).

Cloud Type	Latitude	Longitude	$f_{cld}^{a,b}$	$f_{thin,cld}^{a,b}$
Sc	20°N - 30°N	120°W - 130°W	0.79	0.36
Sc-Cu	10°N - 15°N	130°W - 140°W	0.70	0.37
Cu	10°S - 20°S	150°W - 160°W	0.26	0.84

1081 <sup>a</sup>  $f_{cld}$  and  $f_{cld,thin}$  are two-year means of monthly values. N = 96 (24 months x 4 gridbox values)  
 1082 <sup>b</sup>  $f_{cld}$  and  $f_{cld,thin}$  sampling uncertainties are 0.02 – 0.04 (95% confidence level), and are estimated from the  
 1083 bootstrap method [Efron and Gong, 1983].  
 1084

1085 **Table 4.** Cumulus cloud-field observed and modeled parameters.

Instrument	$f_{cld}$	$f_{thin,cld}$	$\beta^a$	$\alpha_{TOA}^b$
<i>Observed:</i>				
CALIOP	0.26(0.11) <sup>c</sup>	0.84(0.13) <sup>c</sup>	2.66±0.43	
CERES				0.16
MODIS				0.11
<i>Model<sup>d</sup></i>				
CALIOP	0.27	0.93	2.42±0.62	0.16
HSRL	0.26	0.93	1.93±0.12	-
HSRL full-res <sup>c</sup>	0.22	0.79	2.13±0.14	-
MODIS	0.16	0.95	1.97±0.19	-

1086 <sup>a</sup> Cloud length distribution power law exponent±95% CI.

1087 <sup>b</sup> Top-of-atmosphere albedo ( $\alpha^{TOA}$ ) values are used to constrain model mean (clear and cloudy) optical depth.

1088 CERES shortwave TOA albedo (<http://eos.atmos.washington.edu/cgi-bin/ceres/disp.pl?ceres.alb.ann.d>) and

1089 MODIS reflectance ( $R_{MODIS}$ ) estimated as  $R_{MODIS} = R_{cld}f_{cld} + R_{clr}f_{clr}$ , where  $R_{cld}$  is mean visible cloud

1090 reflectance [Wood and Field, 2011] at 1 km cloud length (minimum observable MODIS cloud length),  $R_{clr}$  is a

1091 typical visible clear-sky reflectance (0.08), and  $f_{cld}$  and  $f_{clr}$  weights are CALIOP observed cloud and clear

1092 fractions, for the tropical Pacific trade wind cumulus region [15°S, 155°W; center latitude, longitude]. We

1093 estimate model TOA albedo using model  $\tau_{cld}$  and a simplified cloud albedo equation based upon a two-stream

1094 approximation and assumed asymmetry parameter of 0.85 ( $\alpha^{cld} \approx \tau_{cld} / (\tau_{cld} + 7.7)$ ) [Lacis and Hansen, 1974].

1095 Cloud albedo is calculated for every data point with  $\tau_{cld}$  greater than 0.03 in the modeled cloud field, and these

1096 data are input to a simple TOA albedo model ( $\alpha^{TOA} = R_{clr} + [(1 - R_{clr})^2 * (R_{cld} + R_{cld}^*) / 1 - R_{clr} * (R_{cld} + R_{cld}^*)]$ ),

1097 where  $R_{cld}^* = R_{sfc} * [(1 - R_{cld})^2 / 1 - (R_{cld} * R_{sfc})]$ , with assumed ocean surface albedo ( $R_{sfc}$ ) of 0.05.

1098 <sup>c</sup> CALIOP two-year cloud cover mean (standard deviation) of monthly values. N = 96 (24 months x 4

1099 gridbox values).

1100 <sup>d</sup> Cloud Optical depth threshold for CALIOP and HSRL is  $\tau$  greater than 0.03 (*Chepfer, H., G. Cesana, D.*  
1101 *Winker, B. Getzewich, and M. Vaughan, Comparison of two different cloud climatologies derived from*  
1102 *CALIOP-Level 1 observations: the CALIPSO-ST and the CALIPSO-GOCCP, submitted to Journal of*  
1103 *Atmospheric and Oceanic Technology, 2012*). and for MODIS is  $\tau$  greater than 0.4 [*Ackerman et al., 2008*].  
1104 Optically-thin cloud is defined as a cloudy FOV where the number of photos required for surface detection ( $n$ )  
1105 is greater than the minimum number ( $n^*$ ) required i.e.,  $n/n^* > 1$ .  
1106 <sup>e</sup> Highest resolution HSRL achieved given the HSRL PRF = 200 Hz and assuming aircraft velocity =  $120 \text{ m s}^{-1}$   
1107 ( $\sim 1 \text{ m}$ ).  
1108

1109 **FIGURE CAPTIONS**

1110 **Figure 1.** Cumulative fraction of MODIS L3 liquid water cloud optical depth ( $\tau_{\text{clid}}$ ) for 12  
1111 months of ocean daytime data spanning  $\pm 45^\circ$  latitude. Cloud albedo is approximated as  
1112  $(\tau_{\text{clid}}/\tau_{\text{clid}}+7.7)$  [*Lacis and Hansen, 1974*]. MODIS data suggest approximately one-third of  
1113 low clouds over the oceans have optical depth less than 3 (optically-thin cloud definition,  
1114 this study). In addition, such clouds may contribute up to 15% to total low-cloud albedo.  
1115 Inset figure describes relative scales of CALIOP receiver footprint at single-shot resolution  
1116 and MODIS pixel at cloud product resolution. It should be noted that at most, 3 CALIPSO  
1117 footprints will traverse a given MODIS (on Aqua) pixel.

1118

1119 **Figure 2.** a) A bimodal distribution of 2 years of ocean nighttime lidar integrated attenuated  
1120 backscatter signal at 532 nm (integrated from 0 to 3 km altitude), red dashed line is  
1121 Rayleigh IABS from CALIOP molecular number density profile for a tropical Pacific  
1122 transect. The modes represent clear and cloudy atmospheric regions, as classified by the  
1123 CALIOP Vertical Feature Mask cloud-aerosol discrimination algorithm; b) marine low-  
1124 cloud data suggests optically-thin cloud profiles (as defined in Section 2) populate not only  
1125 the clear-cloud transition region, but also the broader marine low-cloud field.

1126

1127 **Figure 3.** a) A two-dimensional (5 km x 10 km), and, b) one-dimensional (10 km) view of a  
1128 high-cloud-screened broken marine low-cloud scene, generated from CALIOP data along a  
1129 10 km daytime orbit segment over the northwest tropical Atlantic [ $32^\circ\text{N}$ ,  $77^\circ\text{W}$ ] on 27 May  
1130 2007: a) Wide Field Camera (WFC) band-averaged radiance (620-670 nm) at 125 m  
1131 horizontal resolution with CALIOP ground track position superimposed. The WFC data are  
1132 continuous, along-track radiances. Symbols along the CALOP track indicate clear (yellow  
1133 square), optically-thin cloud (white circle), and optically-thick cloud (black circle), b)  
1134 Collocated WFC radiance (black) and Level 1 532 nm attenuated backscatter integrated  
1135 from 0 to 20 km (IABS) (pink) for the same segment as in a). Symbols identifying clear  
1136 profiles, and cloud profiles flagged as optically-thin or optically-thick are also indicated in  
1137 b).

1138

1139 **Figure 4.** Number of data points comprising a) the primary dataset for this analysis, i.e., the  
1140 high-cloud-screened non-polar ocean dataset, and b) the high-cloud-screened cloud length  
1141 dataset. Data for cloud length calculations are a subset of a), requiring a continuous transect  
1142 at least 100 km in length. Data are  $5^{\circ} \times 5^{\circ}$  resolution. Total number of data points: a)  $1.4 \times$   
1143  $10^8$ , b)  $3.9 \times 10^7$ .

1144

1145 **Figure 5.** Two-year mean ocean nighttime a) marine low-cloud cover, b) optically-thin  
1146 marine low-cloud cover, and c) optically-thin fraction of marine low cloud (see Table 2 for  
1147 fraction definitions). Data are  $5^{\circ} \times 5^{\circ}$  resolution. Corresponding domain-averaged values  
1148 are a) 0.50 (0.25), b) 0.23 (0.09), and c) 0.45 (0.28), mean (standard deviation). For this  
1149 calculation gridbox values are weighted by the fraction of each gridbox that is ocean, and  
1150  *$f_{thin\_cld}$  values are also weighted by  $f_{cld}$ .*

1151

1152 **Figure 6.** Joint distribution of monthly nighttime gridded optically-thin fraction of marine  
1153 low cloud, as function of marine low-cloud cover, over-plotted with two-year mean gridbox  
1154 values (grey). Both monthly and two-year averaged data exhibit a strong negative  
1155 correlation. The correlation coefficient for the monthly data is -0.83, significant at the 95%  
1156 confidence level.

1157

1158 **Figure 7.** Cloud length distribution based upon two years of gridded data for marine low  
1159 cloud, “majority optically-thin” cloud, and “majority optically-thick” cloud. Each  
1160 “Majority” category includes only clouds comprising greater than 90% cloud profiles  
1161 belonging to either the optically-thin or optically-thick category. Size bins are  
1162 logarithmically spaced and ordinate axis is normalized frequency.  $\beta = 2$  line shown for  
1163 reference. Inset: marine low-cloud length distribution for the following regions [region  
1164 center latitude, longitude]: California stratocumulus region (Sc) [ $25^{\circ}\text{N}$ ,  $125^{\circ}\text{W}$ ], tropical  
1165 Pacific stratocumulus to cumulus transition (Sc-Cu Trans) [ $15^{\circ}\text{N}$ ,  $135^{\circ}\text{W}$ ], and tropical  
1166 Pacific trade wind cumulus (Cu) [ $15^{\circ}\text{S}$ ,  $155^{\circ}\text{W}$ ]. Abscissa values are cloud length bin  
1167 centers.

1168

1169 **Figure 8.** Cloud length at median cloud cover ( $L_{50}$ ) for a) marine low cloud, b) “majority  
1170 optically-thin” cloud, c) “majority optically-thick” cloud, from two years of ocean nighttime  
1171 data at  $5^\circ \times 5^\circ$  resolution. Low (high) cloud cover values are dominated by small (large)  
1172 clouds. “Majority optically-thin” clouds dominate the cloud type in regions of low cloud  
1173 cover, e.g. trade wind cumulus regions. Each “majority” category includes only clouds  
1174 comprising greater than 90% cloud profiles belonging to either the optically-thin or  
1175 optically-thick category. White over ocean regions in c) indicates that no “majority  
1176 optically-thick” clouds observed.

1177

1178 **Figure 9.** HSRL (black) and CALIOP (pink) marine low-cloud length distributions with  
1179 bootstrap derived 95% confidence intervals [Efron and Gong, 1983]. Dashed line is  $\beta = 2$ ,  
1180 shown for reference. Where cloud size bins overlap, HSRL samples fewer clouds than  
1181 CALIOP can resolve. HSRL samples more clouds at smaller sizes than CALIOP. Sub-  
1182 sampling HSRL (blue stars) produces good agreement between HSRL and CALIOP size  
1183 distributions. CALIOP and HSRL mean  $f_{\text{cld}}$  values agree to within 2 standard deviations.  
1184 However, mean  $f_{\text{thin,cld}}$  values do not agree, with HSRL  $f_{\text{thin,cld}}$  lower than CALIOP  $f_{\text{thin,cld}}$ .  
1185 Inset schematic describes suggested mechanism for  $f_{\text{thin,cld}}$  discrepancy, i.e., it is the result of  
1186 differences in instrument sampling methods: blue circles represent CALIOP 90 m FOVs,  
1187 and orange bar shows the relative size of the HSRL FOV compared to the CALIOP FOV.

1188

1189 **Figure 10.** Optically-thin fraction of marine low-cloud length for non-polar oceans, and  
1190 three regions with different dominant cloud types [region center latitude, longitude]:  
1191 California stratocumulus region (Sc) [ $25^\circ\text{N}$ ,  $125^\circ\text{W}$ ], tropical Pacific stratocumulus to  
1192 cumulus transition (Sc-Cu Trans) [ $15^\circ\text{N}$ ,  $135^\circ\text{W}$ ], and tropical Pacific trade wind cumulus  
1193 (Cu) [ $15^\circ\text{S}$ ,  $155^\circ\text{W}$ ]. Increasing optically-thin fraction as cloud length decreases is  
1194 observed across all regions, despite different cloud types. Vertical bars are 95% confidence  
1195 intervals calculated from the bootstrap method [Efron and Gong, 1983]. Abscissa values  
1196 are cloud length bin centers.

1197

1198 **Figure 11.** Vertical distribution of two-year mean optically-thick (red) and optically-thin  
1199 (black) cloud top height (top row) and, optically-thin cloud base height (bottom row) for 3

1200 cloud regimes [region center latitude, longitude]: a) and d) California stratocumulus region  
1201 (Sc) [25°N, 125°W], b) and e) tropical Pacific stratocumulus to cumulus transition (Sc-Cu  
1202 Trans) [15°N, 135°W], and c) and f) tropical Pacific trade wind cumulus (Cu) [15°S,  
1203 155°W]. Layer boundaries are identified by CALIPSO feature detection algorithm at full  
1204 horizontal resolution, and 60m vertical resolution.

1205

1206 **Figure 12.** Normalized density distributions of optically-thin cloud top height and layer  
1207 depth for 3 cloud types [region center latitude, longitude]: a) California stratocumulus  
1208 region [25°N, 125°W], b) tropical Pacific stratocumulus to cumulus transition (Sc-Cu  
1209 Trans) [15°N, 135°W], and c) tropical Pacific trade wind cumulus (Cu) [15°S, 155°W].  
1210 Inset figures are layer depth distributions for each region, overlaid with a gamma  
1211 distribution (red). Text is region mean (standard deviation) layer depth (km), and number  
1212 of full-resolution optically-thin cloud profiles for each region. Layer depth histogram is  
1213 normalized such that the area under the curve = 1.

1214

1215 **Figure 13.** An aerial photograph of a typical trade wind cumulus cloud field observed in the  
1216 Western Caribbean during the Rain in Cumulus over Ocean (RICO) field program.  
1217 Apparent is the wide range of cloud depths and horizontal extents. (Image courtesy of B.  
1218 Stevens).

1219

1220 **Figure 14.** Two years of CALIOP tropical Pacific trade wind region observations (black  
1221 dots)  $\pm$  95% confidence intervals from bootstrap method [*Efron and Gong, 1983*], and  
1222 cascade model output simulating this cloud regime, sampled with CALIOP FOV (black  
1223 circles), MODIS FOV (red circles), HSRL FOV (blue circles), and HSRL Full Res i.e.,  
1224 highest resolution sampling of model output (blue squares). Size bins are logarithmically  
1225 spaced and ordinate axis is normalized frequency.  $\beta = 2$  line (dashed line) shown for  
1226 reference. Cascade model output from 100 runs with  $H = 0.07$  and  $p = 0.0005$ . Abscissa  
1227 values are cloud length bin centers.

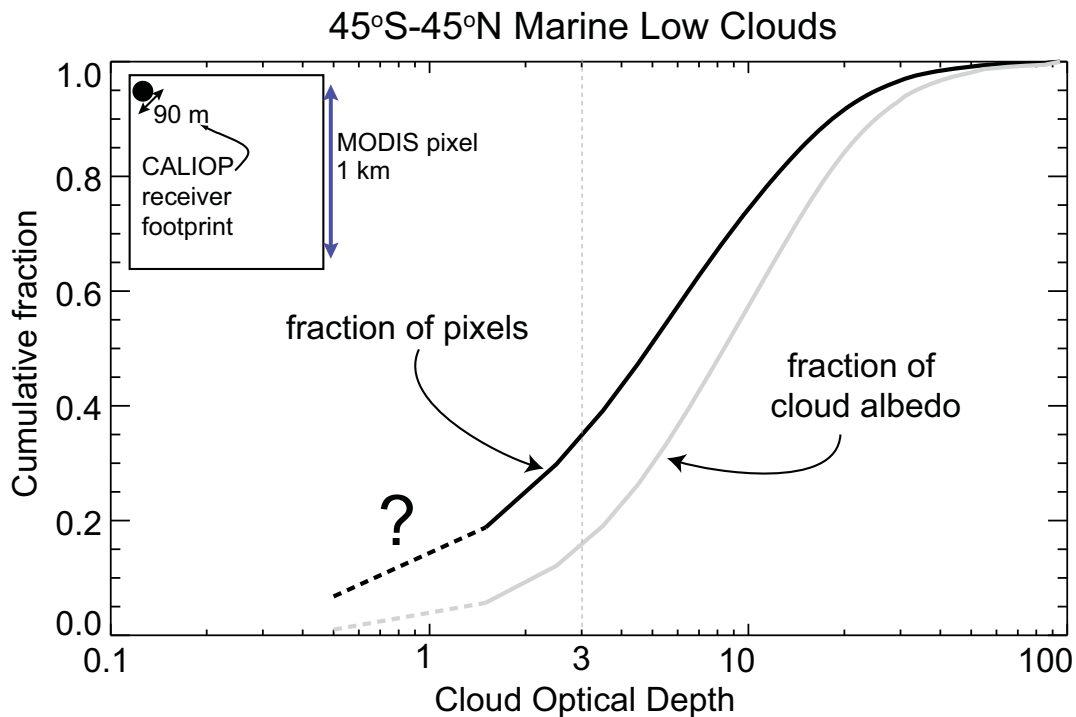
1228

1229 **Figure 15** Cloudy fields of view (FOVs) and partially cloud-filled FOVs (pFOVs) values  
1230 binned by cloud length, as estimated from model simulated optical depth ( $\tau$ ) values sampled  
1231 at CALIOP FOV resolution. Displayed for each cloud length bin is the fraction of all cloud  
1232 FOVs that are less than 100% cloudy, i.e., pFOV (filled circles); the median cloud fraction  
1233 for pFOVs (filled squares); the fraction of optically-thin pFOVs incorrectly classified as  
1234 optically-thin, i.e., a FOV with mean  $n/n_*$  and mean  $\tau$  of the cloudy-only segment greater  
1235 than 3 (filled diamonds). Cascade model output is from 100 runs with  $H = 0.07$  and  $p =$   
1236  $0.0005$ .

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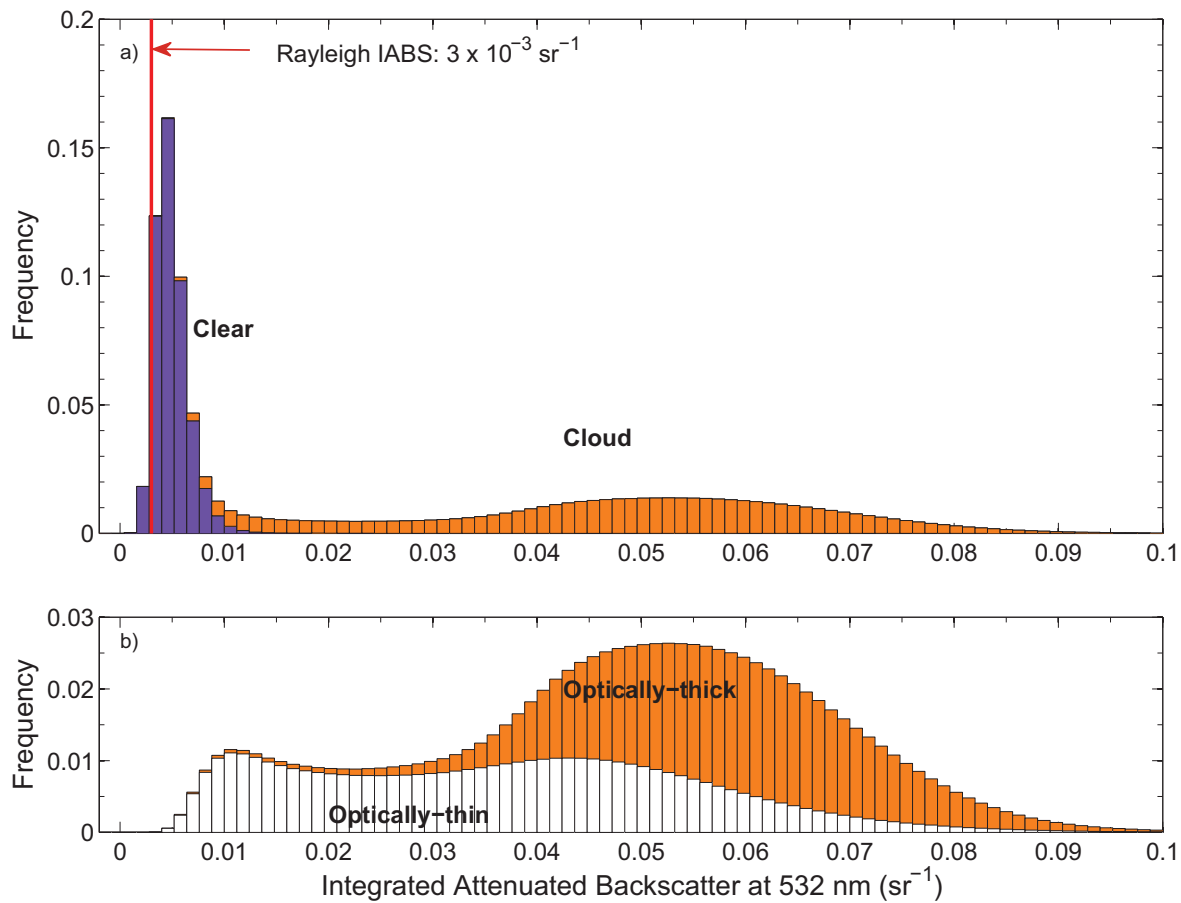
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1241

1242 **Figure 1.** Cumulative fraction of MODIS L3 liquid water cloud optical depth ( $\tau_{\text{cl}d}$ ) for 12  
 1243 months of ocean daytime data spanning  $\pm 45^\circ$  latitude. Cloud albedo is approximated as  
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 1247 Inset figure describes relative scales of CALIOP receiver footprint at single-shot resolution  
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 1249 footprints will traverse a given MODIS (on Aqua) pixel.

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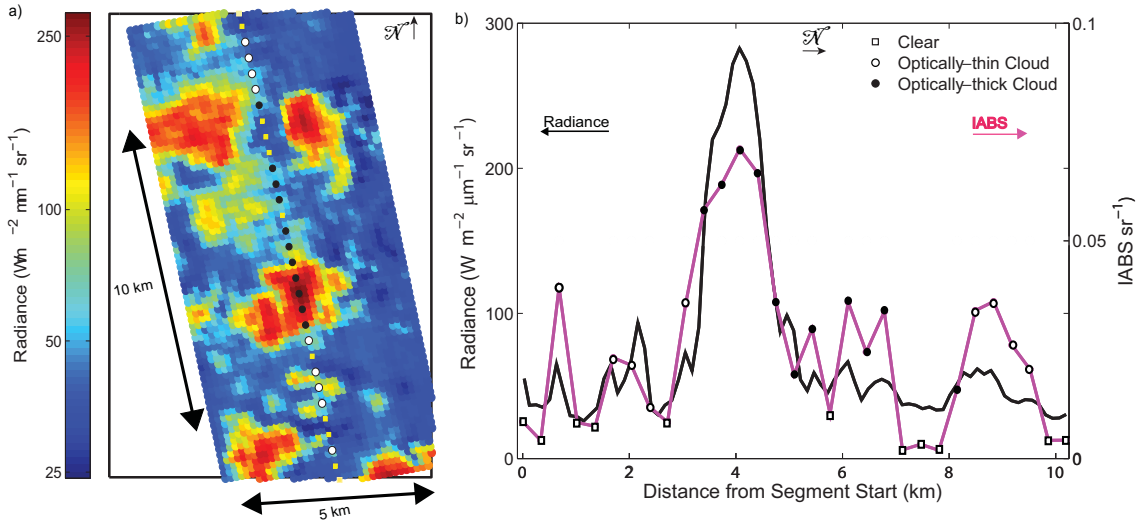


1251

1252 **Figure 2.** a) A bimodal distribution of 2 years of ocean nighttime lidar integrated  
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 1255 transect. The modes represent clear and cloudy atmospheric regions, as classified by the  
 1256 CALIOP Vertical Feature Mask cloud-aerosol discrimination algorithm; b) marine low-  
 1257 cloud data suggests optically-thin cloud profiles (as defined in Section 2) populate not only  
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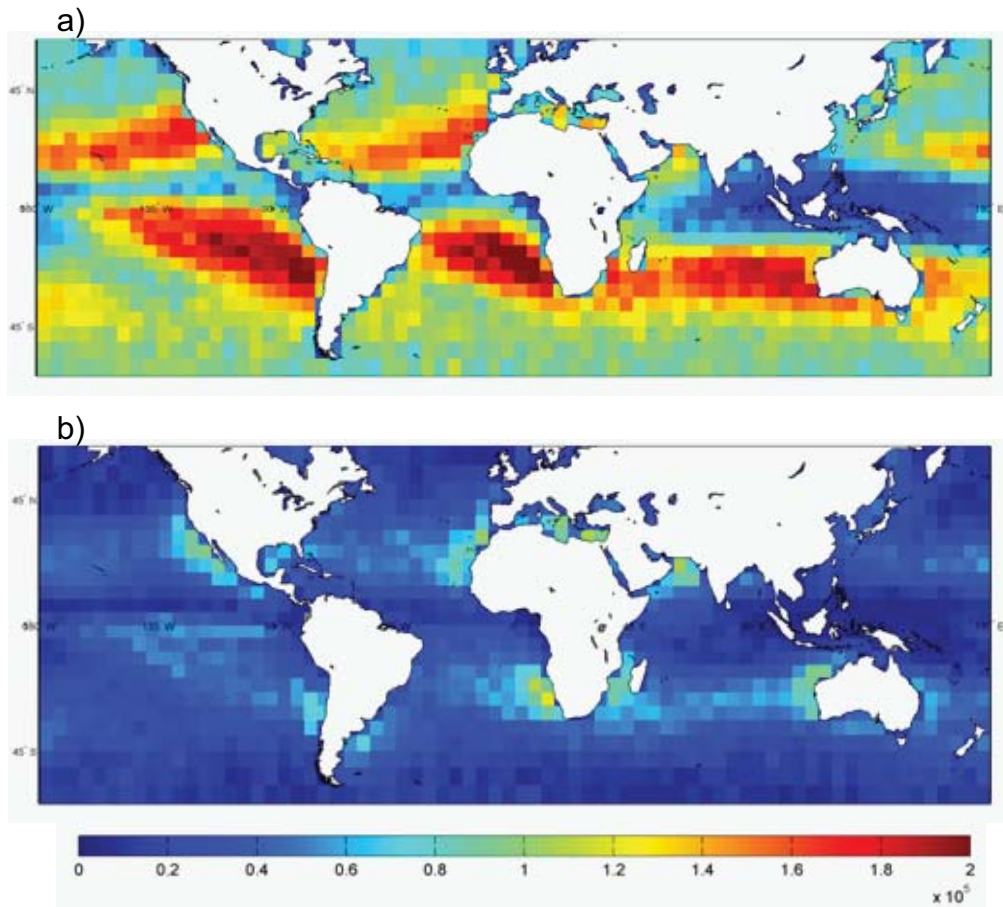
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Orbit Segment: Latitude 32°N, Longitude 77°W



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**Figure 3.** a) A two-dimensional (5 km x 10 km), and, b) one-dimensional (10 km) view of a high-cloud-screened broken marine low-cloud scene, generated from CALIOP data along a 10 km daytime orbit segment over the northwest tropical Atlantic [32°N, 77°W] on 27 May 2007: a) Wide Field Camera (WFC) band-averaged radiance (620-670 nm) at 125 m horizontal resolution with CALIOP ground track position superimposed. The WFC data are continuous, along-track radiances. Symbols along the CALOP track indicate clear (yellow square), optically-thin cloud (white circle), and optically-thick cloud (black circle), b) Collocated WFC radiance (black) and Level 1 532 nm attenuated backscatter integrated from 0 to 20 km (IABS) (pink) for the same segment as in a). Symbols identifying clear profiles, and cloud profiles flagged as optically-thin or optically-thick are also indicated in b).

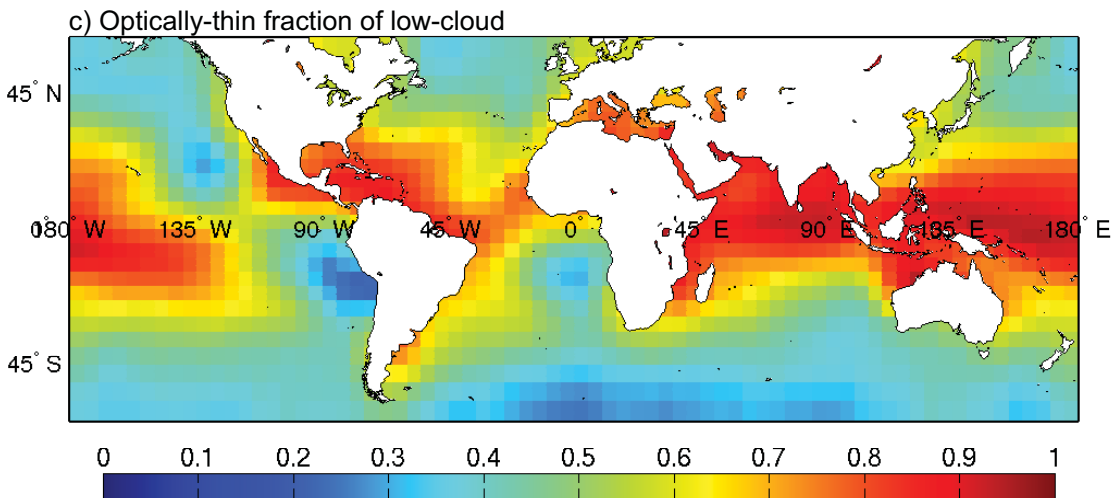
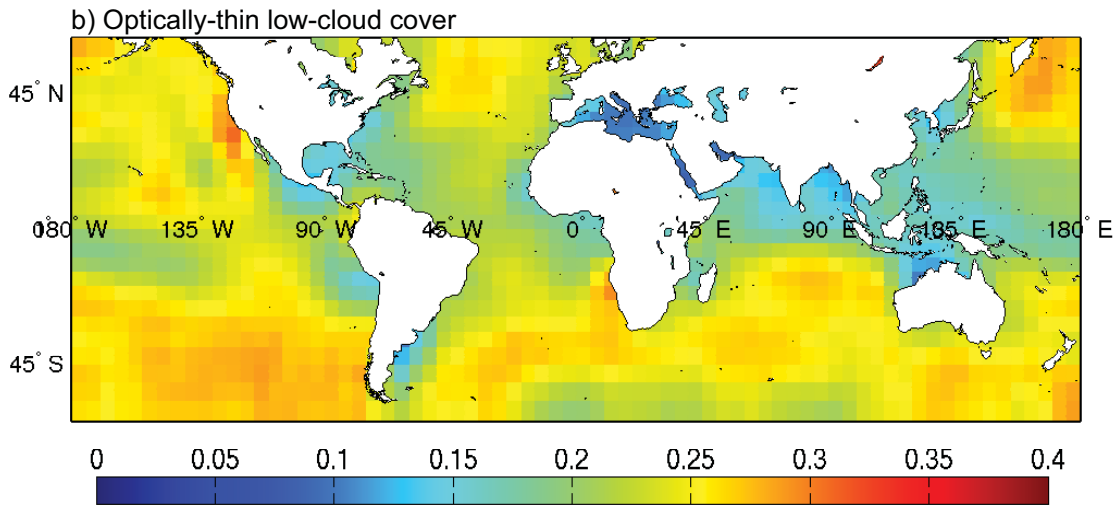
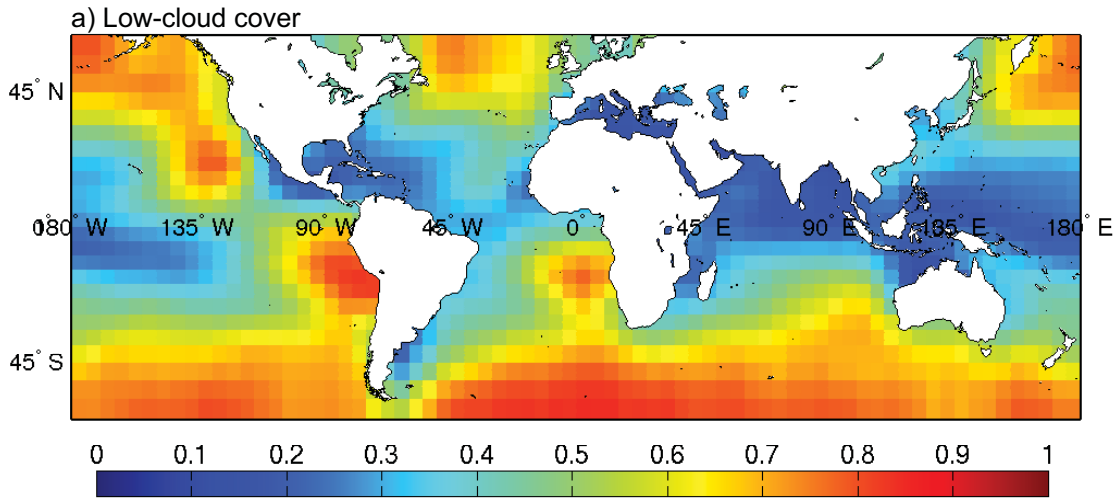


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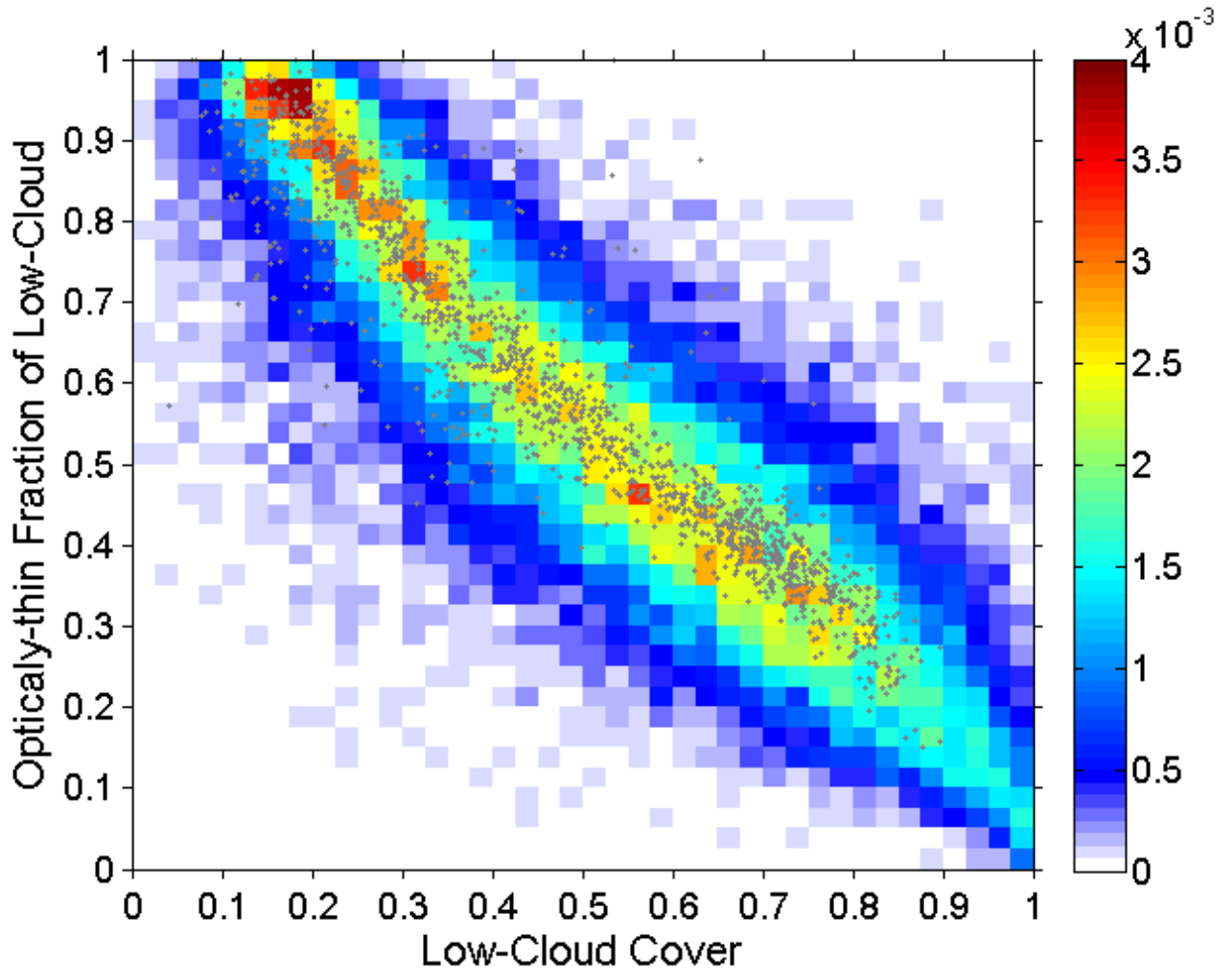
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 1277 high-cloud-screened non-polar ocean dataset, and b) the high-cloud-screened cloud length  
 1278 dataset. Data for cloud length calculations are a subset of a), requiring a continuous transect  
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 1280  $10^8$ , b)  $3.9 \times 10^7$ .

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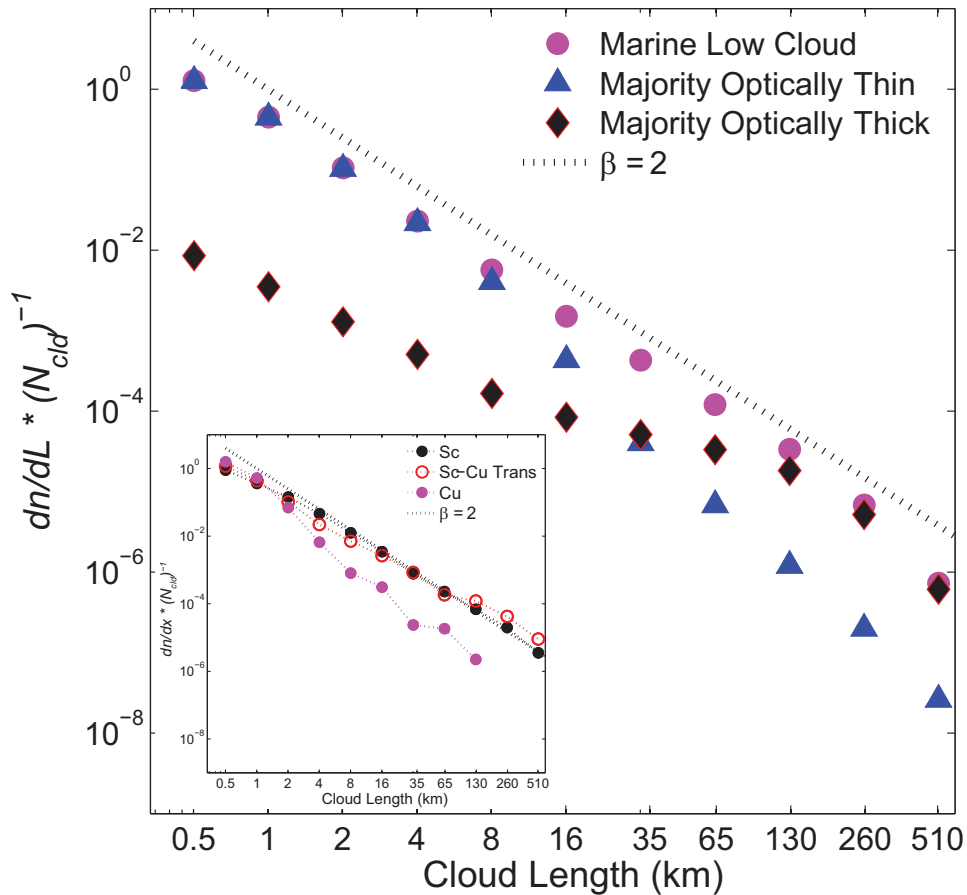
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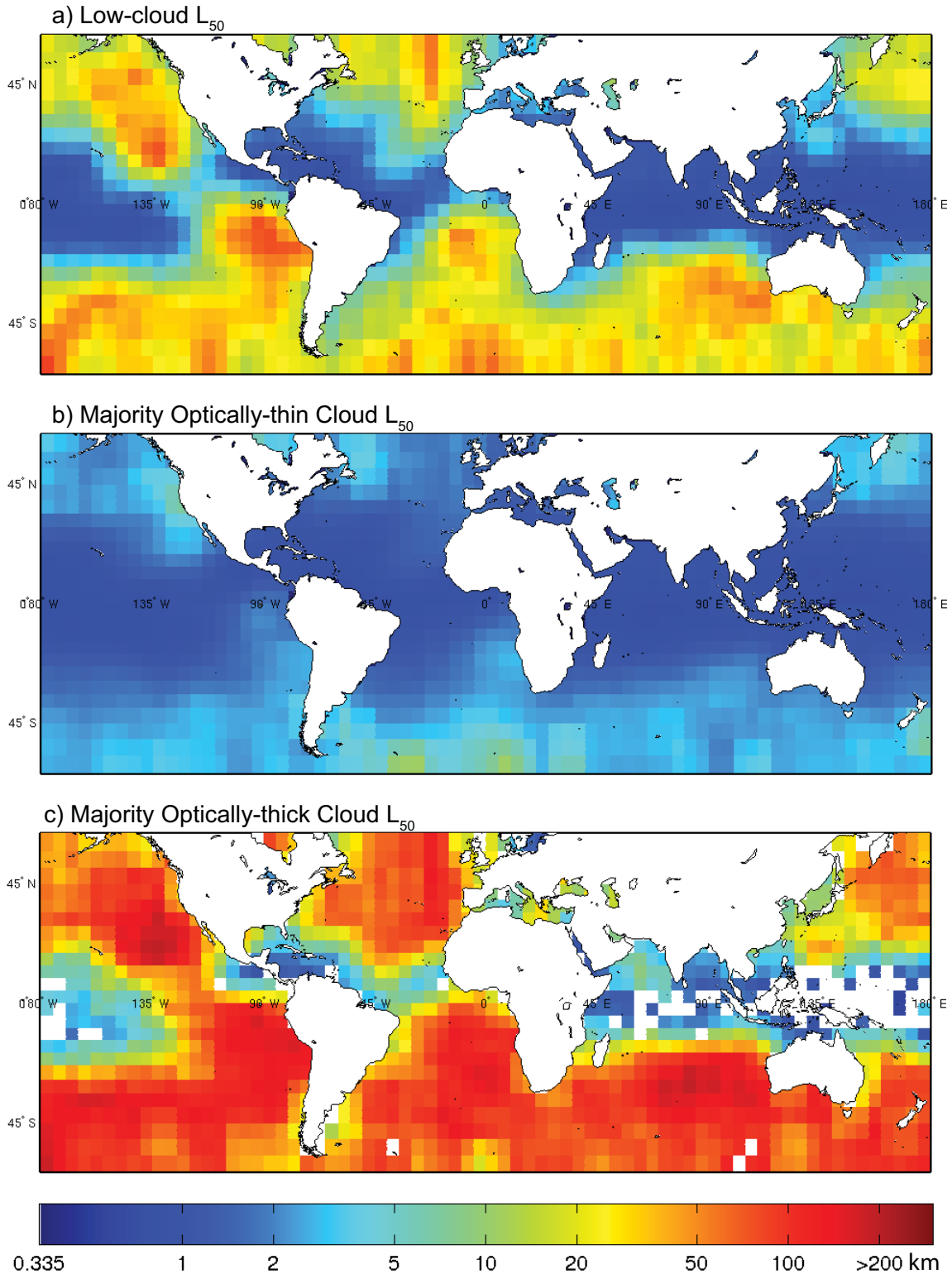
1284 **Figure 5.** Two-year mean ocean nighttime a) marine low-cloud cover, b) optically-thin  
1285 marine low-cloud cover, and c) optically-thin fraction of marine low cloud (see Table 2 for  
1286 fraction definitions). Data are  $5^\circ \times 5^\circ$  resolution. Corresponding domain-averaged values  
1287 are a) 0.50 (0.25), b) 0.23 (0.09), and c) 0.45 (0.28), mean (standard deviation). For this  
1288 calculation gridbox values are weighted by the fraction of each gridbox that is ocean, and  
1289  *$f_{thin\_cld}$  values are also weighted by  $f_{cld}$ .*



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 1291 **Figure 6.** Joint distribution of monthly nighttime gridded optically-thin fraction of marine  
 1292 low cloud, as function of marine low-cloud cover, over-plotted with two-year mean gridbox  
 1293 values (grey). Both monthly and two-year averaged data exhibit a strong negative  
 1294 correlation. The correlation coefficient for the monthly data is -0.83, significant at the 95%  
 1295 confidence level.  
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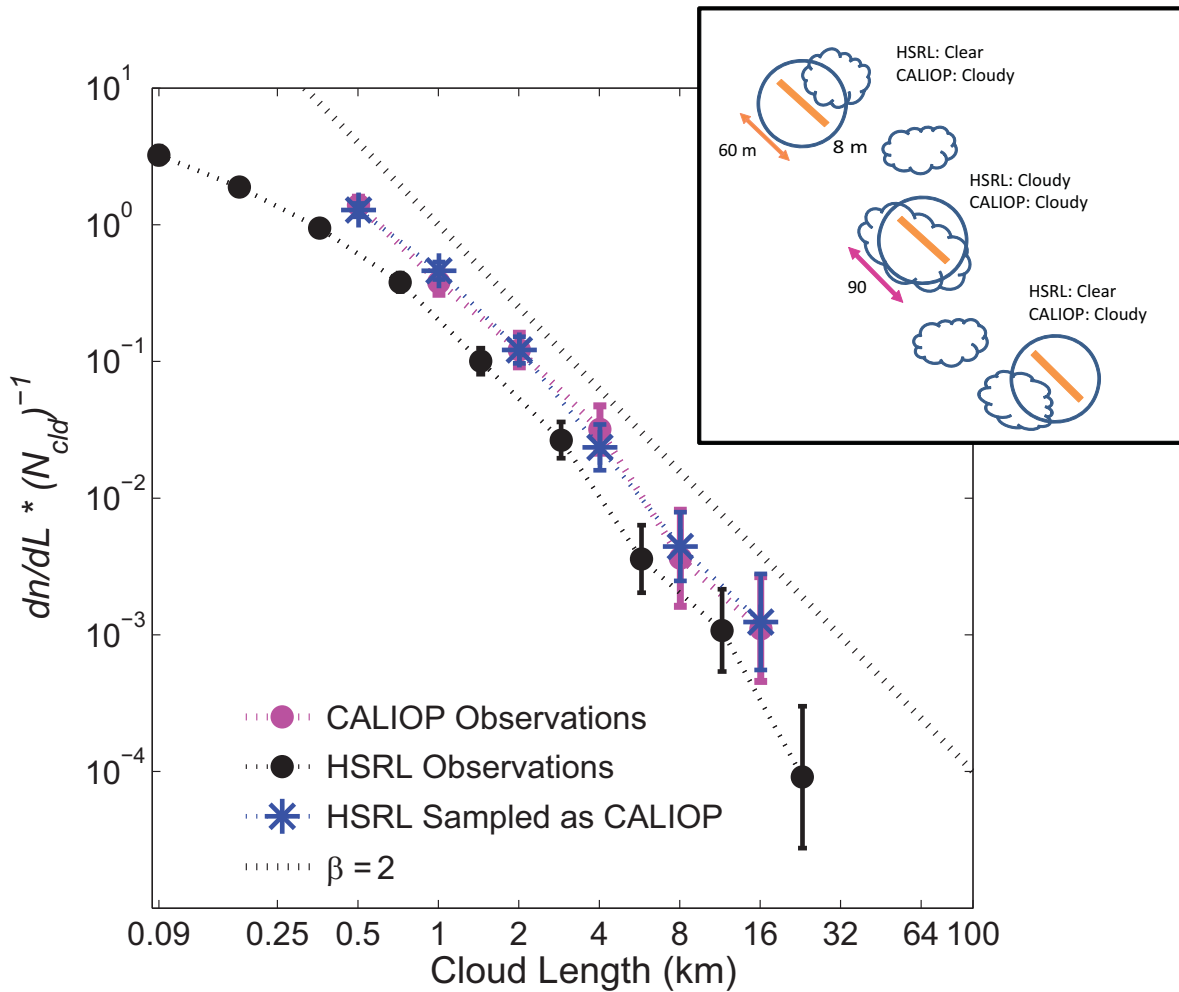


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 1298 **Figure 7.** Cloud length distribution based upon two years of gridded data for marine low  
 1299 cloud, “majority optically-thin” cloud, and “majority optically-thick” cloud. Each  
 1300 “Majority” category includes only clouds comprising greater than 90% cloud profiles  
 1301 belonging to either the optically-thin or optically-thick category. Size bins are  
 1302 logarithmically spaced and ordinate axis is normalized frequency.  $\beta = 2$  line shown for  
 1303 reference. Inset: marine low-cloud length distribution for the following regions [region  
 1304 center latitude, longitude]: California stratocumulus region (Sc) [25°N, 125°W], tropical  
 1305 Pacific stratocumulus to cumulus transition (Sc-Cu Trans) [15°N, 135°W], and tropical  
 1306 Pacific trade wind cumulus (Cu) [15°S, 155°W]. Abscissa values are cloud length bin  
 1307 centers.

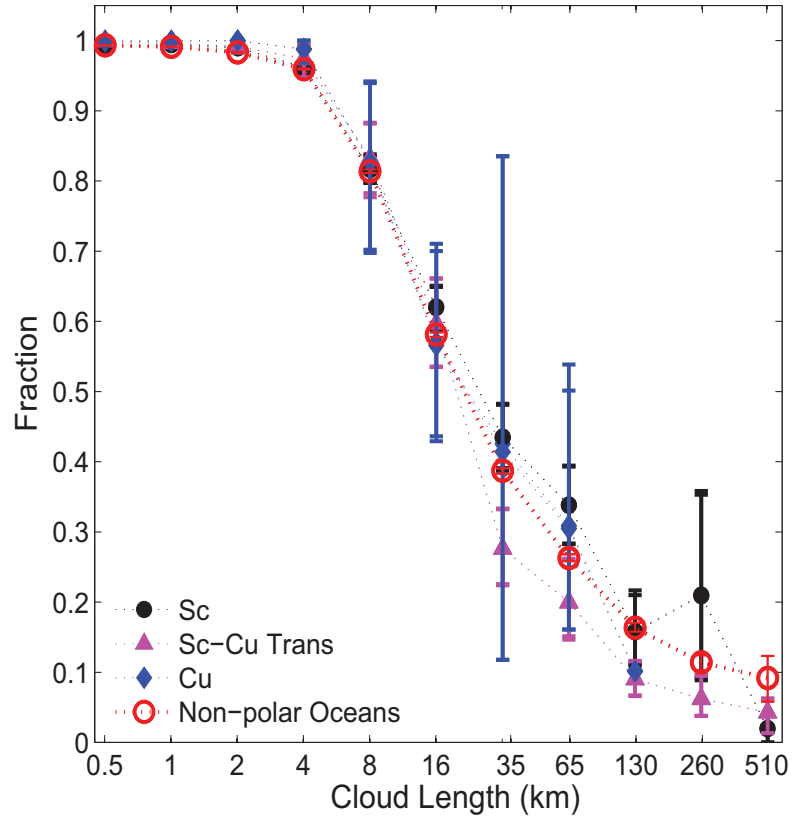


1308  
 1309 **Figure 8.** Cloud length at median cloud cover ( $L_{50}$ ) for a) marine low cloud, b) “majority  
 1310 optically-thin” cloud, c) “majority optically-thick” cloud, from two years of ocean nighttime  
 1311 data at 5°x5° resolution. Low (high) cloud cover values are dominated by small (large)  
 1312 clouds. “Majority optically-thin” clouds dominate the cloud type in regions of low cloud  
 1313 cover, e.g. trade wind cumulus regions. Each “majority” category includes only clouds

1314 comprising greater than 90% cloud profiles belonging to either the optically-thin or  
1315 optically-thick category. White over ocean regions in c) indicates that no “majority  
1316 optically-thick” clouds observed.

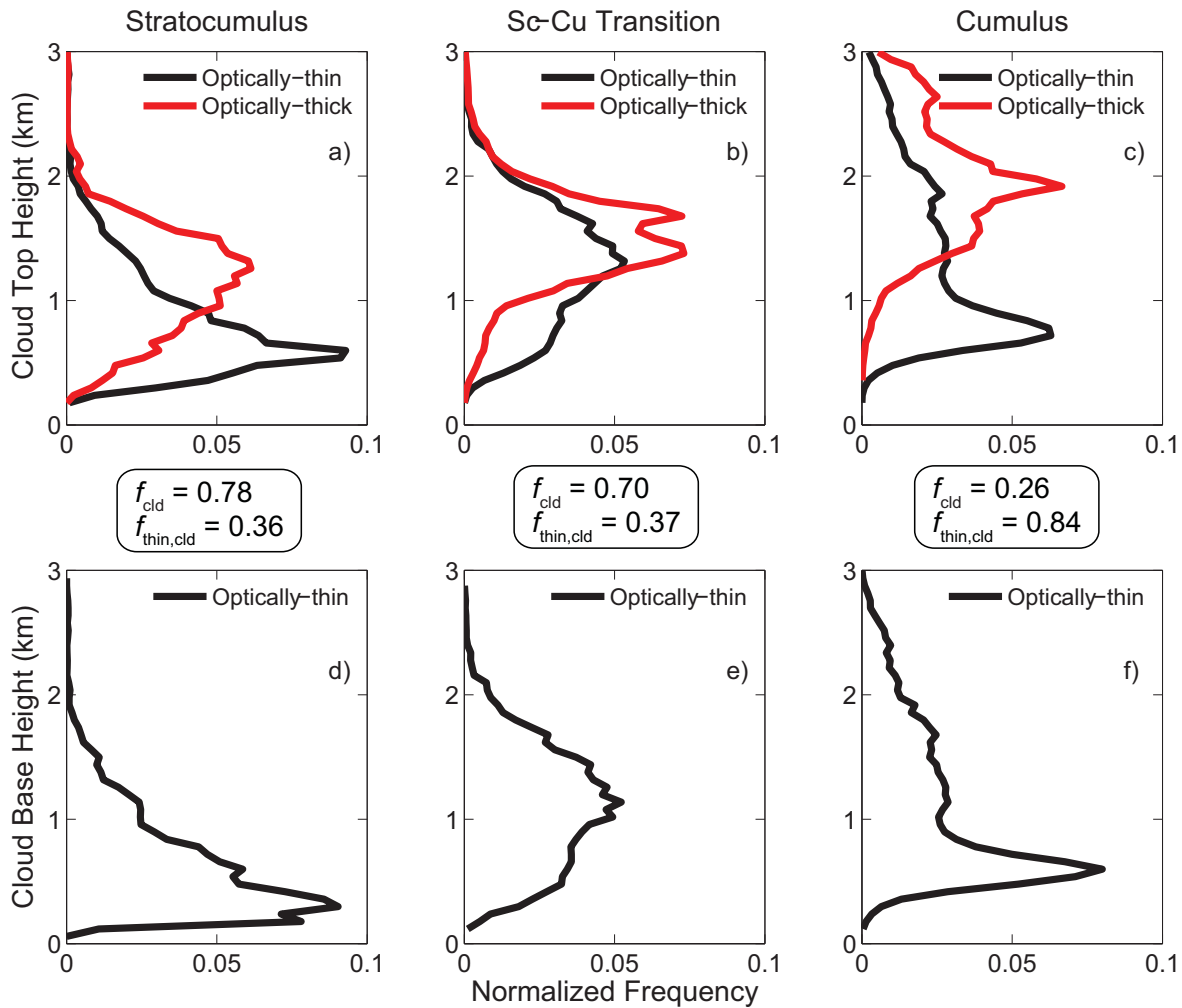


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 1318 **Figure 9.** HSRL (black) and CALIOP (pink) marine low-cloud length distributions with  
 1319 bootstrap derived 95% confidence intervals [Efron and Gong, 1983]. Dashed line is  $\beta = 2$ ,  
 1320 shown for reference. Where cloud size bins overlap, HSRL samples fewer clouds than  
 1321 CALIOP can resolve. HSRL samples more clouds at smaller sizes than CALIOP. Sub-  
 1322 sampling HSRL (blue stars) produces good agreement between HSRL and CALIOP size  
 1323 distributions. CALIOP and HSRL mean  $f_{cld}$  values agree to within 2 standard deviations.  
 1324 However, mean  $f_{thin,cld}$  values do not agree, with HSRL  $f_{thin,cld}$  lower than CALIOP  $f_{thin,cld}$ .  
 1325 Inset schematic describes suggested mechanism for  $f_{thin,cld}$  discrepancy, i.e., it is the result of  
 1326 differences in instrument sampling methods: blue circles represent CALIOP 90 m FOVs,  
 1327 and orange bar shows the relative size of the HSRL FOV compared to the CALIOP FOV.  
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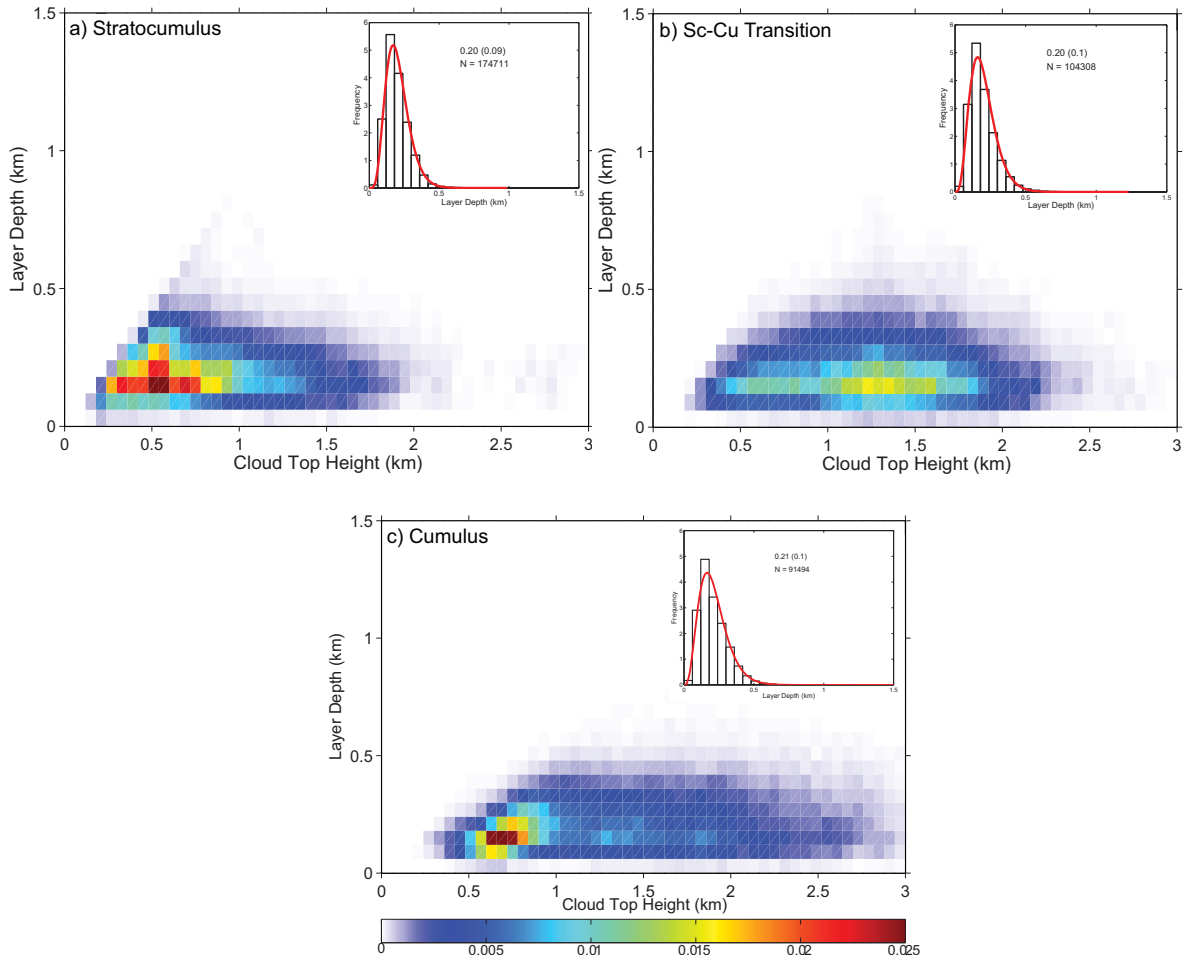
1329  
 1330 **Figure 10.** Optically-thin fraction of marine low-cloud length for non-polar oceans, and  
 1331 three regions with different dominant cloud types [region center latitude, longitude]:  
 1332 California stratocumulus region (Sc) [25°N, 125°W], tropical Pacific stratocumulus to  
 1333 cumulus transition (Sc-Cu Trans) [15°N, 135°W], and tropical Pacific trade wind cumulus  
 1334 (Cu) [15°S, 155°W]. Increasing optically-thin fraction as cloud length decreases is  
 1335 observed across all regions, despite different cloud types. Vertical bars are 95% confidence  
 1336 intervals calculated from the bootstrap method [Efron and Gong, 1983]. Abscissa values  
 1337 are cloud length bin center.

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**Figure 11.** Vertical distribution of two-year mean optically-thick (red) and optically-thin (black) cloud top height (top row) and, optically-thin cloud base height (bottom row) for 3 cloud regimes [region center latitude, longitude]: a) and d) California stratocumulus region (Sc) [25°N, 125°W], b) and e) tropical Pacific stratocumulus to cumulus transition (Sc-Cu Trans) [15°N, 135°W], and c) and f) tropical Pacific trade wind cumulus (Cu) [15°S, 155°W]. Layer boundaries are identified by CALIPSO feature detection algorithm at full horizontal resolution, and 60m vertical resolution.

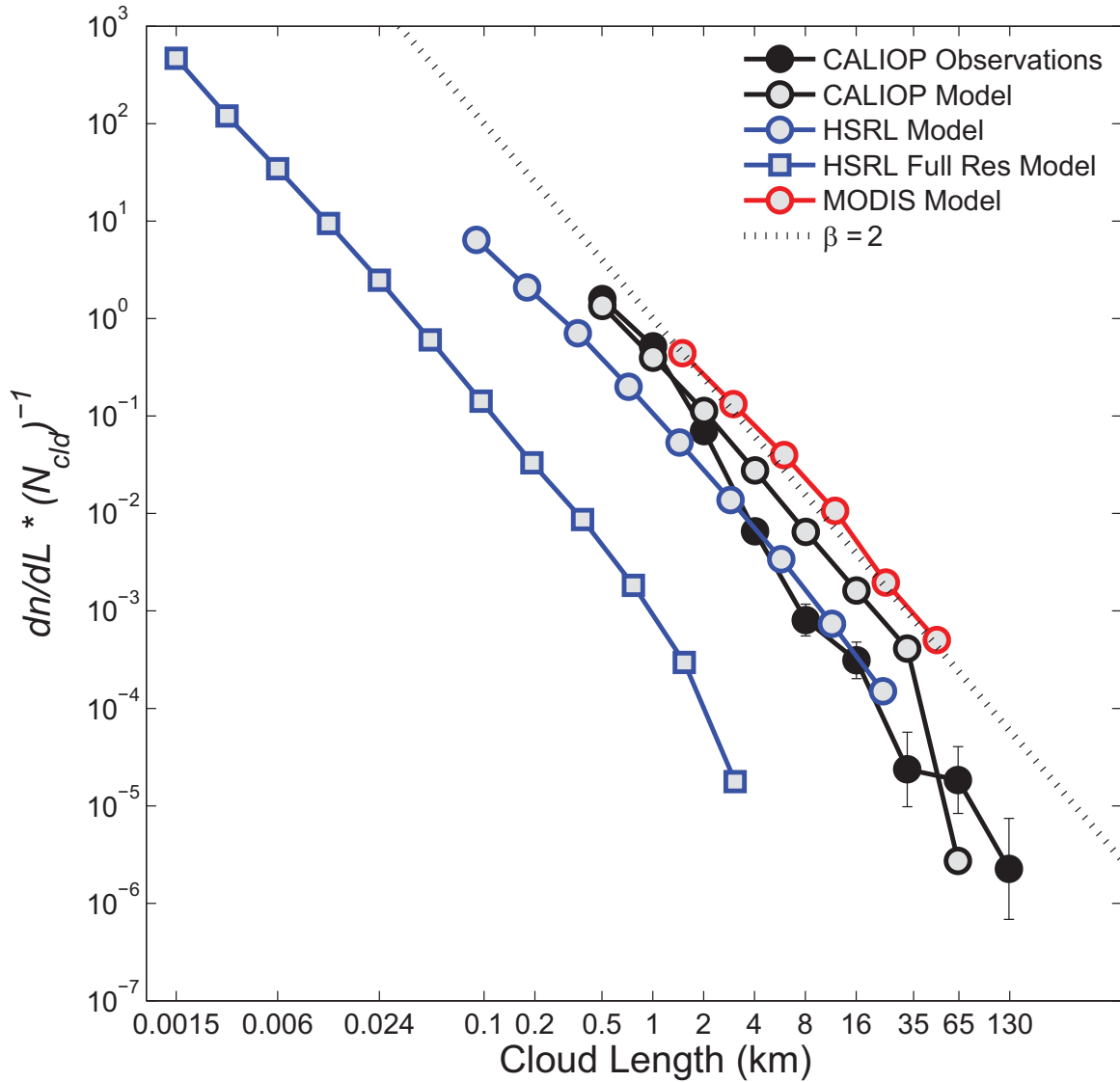


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 1349 **Figure 12.** Normalized density distributions of optically-thin cloud top height and layer  
 1350 depth for 3 cloud types [region center latitude, longitude]: a) California stratocumulus  
 1351 region [25°N, 125°W], b) tropical Pacific stratocumulus to cumulus transition (Sc-Cu  
 1352 Trans) [15°N, 135°W], and c) tropical Pacific trade wind cumulus (Cu) [15°S, 155°W].  
 1353 Inset figures are layer depth distributions for each region, overlaid with a gamma  
 1354 distribution (red). Text is region mean (standard deviation) layer depth (km), and number  
 1355 of full-resolution optically-thin cloud profiles for each region. Layer depth histogram is  
 1356 normalized such that the area under the curve = 1.  
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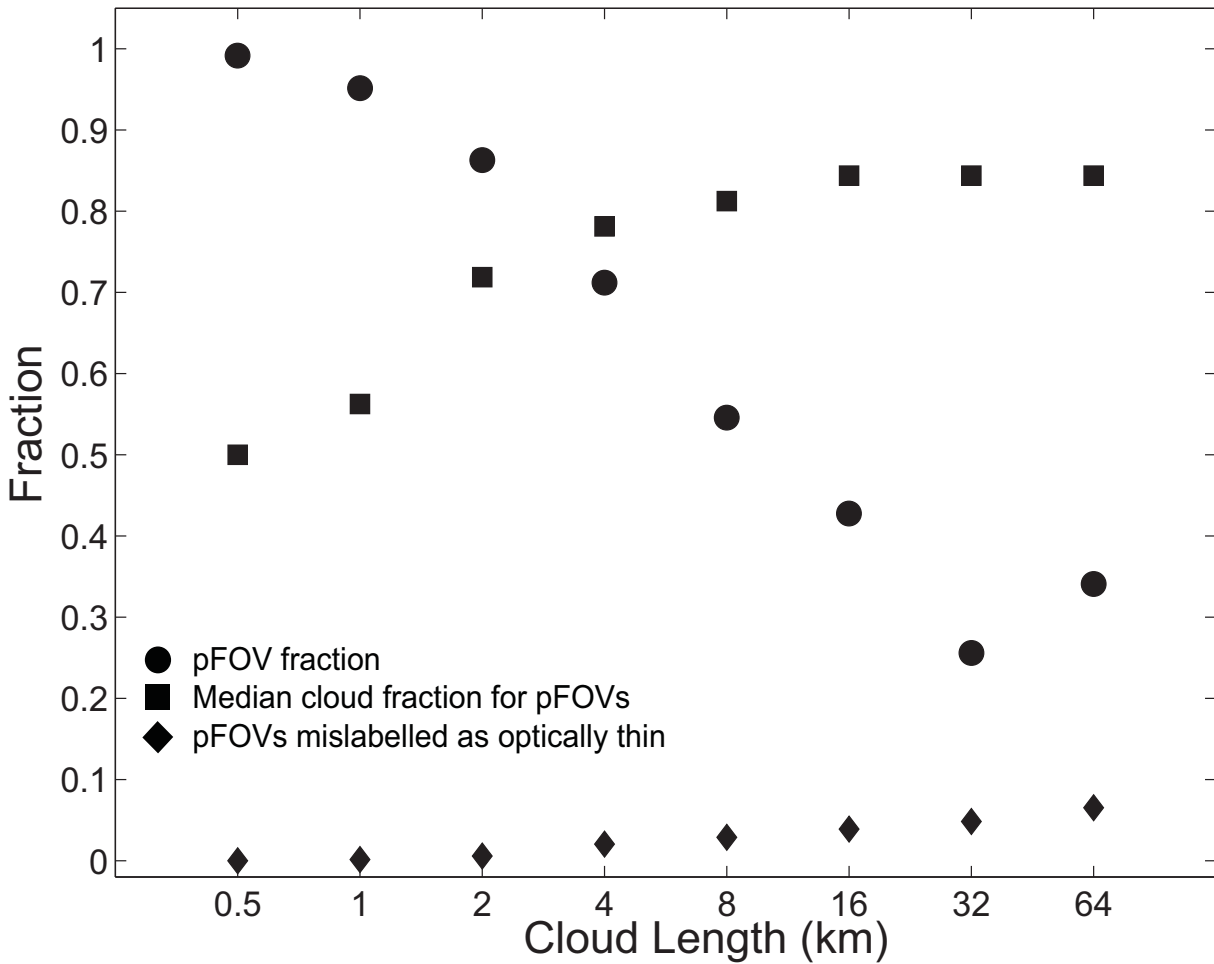


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**Figure 13.** An aerial photograph of a typical trade wind cumulus cloud field observed in the Western Caribbean during the Rain in Cumulus over Ocean (RICO) field program. Apparent is the wide range of cloud depths and horizontal extents. (Image courtesy of B. Stevens).



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 1367 **Figure 14.** Two years of CALIOP tropical Pacific trade wind region observations (black  
 1368 dots)  $\pm$  95% confidence intervals from bootstrap method [Efron and Gong, 1983], and  
 1369 cascade model output simulating this cloud regime, sampled with CALIOP FOV (black  
 1370 circles), MODIS FOV (red circles), HSRL FOV (blue circles), and HSRL Full Res i.e.,  
 1371 highest resolution sampling of model output (blue squares). Size bins are logarithmically  
 1372 spaced and ordinate axis is normalized frequency.  $\beta = 2$  line (dashed line) shown for  
 1373 reference. Cascade model output from 100 runs with  $H = 0.015$  and  $p = 0.05$ . Abscissa  
 1374 values are cloud length bin center.  
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**Figure 15** Cloudy fields of view (FOVs) and partially cloud-filled FOVs (pFOVs) values binned by cloud length, as estimated from model simulated optical depth ( $\tau$ ) values sampled at CALIOP FOV resolution. Displayed for each cloud length bin is the fraction of all cloud FOVs that are less than 100% cloudy, i.e., pFOV (filled circles); the median cloud fraction for pFOVs (filled squares); the fraction of optically-thin pFOVs incorrectly classified as optically-thin, i.e., a FOV with mean  $n/n_*$  and mean  $\tau$  of the cloudy-only segment greater than 3 (filled diamonds). Cascade model output is from 100 runs with  $H = 0.07$  and  $p = 0.0005$ .

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