Observable Surface Anomalies Preceding Simulated Isolated Convective Initiation

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ABSTRACT

Idealized ensemble simulations of unforced, isolated convective initiation (CI) are analyzed to identify storm-scale features in surface weather fields that precede initiation in a variety of background environments and the observations that would be needed to resolve these features. We identify precipitating storms with an object-based method and generate composites of surface anomalies for the variables of interest surrounding times and locations of initiation. Anomaly magnitudes and correlation length scales throughout the CI process are examined in detail and compare favorably to anomaly estimates from previous observational and modeling studies. We find negative temperature anomalies due to cloud shadowing to be the most prominent storm-scale feature prior to initiation. Significant spatial correlations are shown to extend from the surface throughout the boundary layer and even into the cloud-bearing layer once deep convective clouds become established. The findings are discussed in the context of data assimilation, particularly with respect to current assumptions about surface observation error. We show that, to resolve the storm-scale anomalies in these simulations, the minimum necessary temperature and wind observation densities would likely be limited by spatial correlation length scale while moisture and pressure observations are more limited by observation error.
1. Introduction

The prediction of exactly when and where precipitating convective storms will develop within an unstable air mass (“convective initiation”; CI) remains a challenging problem. This is particularly true for the prediction of “air-mass” type thunderstorms which can develop within a broad region of instability without apparent large-scale forcing (“unforced”). Numerous studies ascribe failure to skilfully predict storm-scale convective development to a lack of spatially and temporally dense observations (e.g. Stensrud and Fritsch 1994; Dabberdt and Schlatter 1996; Mass et al. 2002; Roebber et al. 2002; Fowle and Roebber 2003; Dabberdt et al. 2005; Gallus et al. 2005; Snook et al. 2014). Of current observation platforms, only Doppler radars and some satellite products are able to produce observations on a temporal and spatial scale consistent with these requirements. As such, these platforms have received the majority of research attention. However, these observations provide limited utility in assessing the pre-convective environment, essential for the prediction of CI.

Surface weather observations have received more limited attention for describing storm-scale features in the pre-convective environment, chiefly due to their relatively poor spatial density (e.g. Mass et al. 2002; Sun et al. 2013). Recent studies have demonstrated that frequent, dense surface observations can contribute to improved CI forecasts (e.g. Liu and Xue 2008; Sobash and Stensrud 2015). To date, this forecast improvement has mostly been attributed to improved representation of broader mesoscale forcings (e.g. fronts or drylines) and reducing near-surface model biases (Sobash and Stensrud 2015). Here we wish to expand these findings by examining the potential for dense surface observations to describe structures on the scale of individual thunderstorms.

Surface observations are available at sub-hourly, kilometer-scale density or better in many locations from extant, yet underutilized networks such as private citizen weather stations (Madaus
et al. 2014) and smartphone barometers (Mass and Madaus 2014). Exploiting these observations and enhancing existing surface observation networks may provide a storm-scale surface network capable of satisfying the requirements for CI forecasting. Yet, there has been little work to date describing how the convective initiation process evolves in surface weather fields, particularly for “unforced” convection. As such, there is little description of which surface variables are most relevant to observe, how sensitive these observations would need to be, or the requisite density of these observations in time and space. To assess these questions, the surface evolution of the CI process in modern, kilometer-scale numerical models must be better documented. This work uses idealized simulations of “unforced” convective initiation to investigate if there are observable storm-scale features in surface weather fields that precede isolated convective initiation in light of current observation limitations.

The remainder of this paper is outlined as follows. A review of previous attempts to quantify the relationship between surface characteristics and convective initiation is given in Section 2. Our methods, including the model configuration and analysis techniques are given in Section 3. The results are presented in Section 4 with a discussion of the implications of these results in Section 5. We conclude by summarizing our findings and their limitations in Section 6.

2. Background

a. Role of the surface in convective initiation

The process of convective initiation has typically been viewed in the parcel model context. Surface (or near-surface) observations of atmospheric properties (temperature, moisture, wind and pressure) are connected to the convective initiation process from the parcel model in the following ways:
• **Temperature** — Relatively warm regions have greater free energy, and therefore an increased potential for positive buoyancy and reduced convective inhibition (CIN).

• **Moisture** — Moister rising parcels reach their lifted condensation level (LCL) at a lower altitude. Once the LCL is reached, latent heat release from condensation contributes to positive buoyancy and reduced CIN.

• **Wind** — Areas of convergence near the surface lead to upward vertical motion through mass conservation, sustaining rising parcels against CIN.

• **Pressure** — Areas of lower pressure near the surface are often promote low-level convergence and can indicate rising motion above.

Characterizing the variability in these fields is crucial for successful CI forecasting. Near-surface temperature and moisture may vary for a number of reasons, but surface fluxes strongly influence local variability. Several studies note that surface heat and moisture flux variations can impact when and where convection will develop (e.g. Koch 1984; Yan and Anthes 1988; Childs et al. 2006; Holt et al. 2006). However, the relative importance of heat and moisture fluxes remains uncertain. Crook (1996) and Kang and Bryan (2011) find, through observation and idealized modeling studies, that local surface temperature variations are more important than moisture variability in determining CI occurrence. However, other modeling studies based on specific cases (e.g. Droegemeier and Wilhelmsen 1985; Clark et al. 2004; Martin and Xue 2006; Zhang et al. 2015) suggest that CI is most sensitive to moisture variability. The relative importance of these variables has implications for which surface observations are more likely to discriminate locations of CI, and needs to be investigated further.

There are additional ways for storm-scale variability in surface fields to locally enable or inhibit CI. Variations in cloudiness can create gradients in diabatic heating at the surface (and, conse-
quent, surface temperature and moisture) due to cloud “shadowing” (Koch 1984; Lohou and Patton 2013; Rieck et al. 2015). Observational studies show that diabatic heating differences can change local pressure distributions, generating areas of increased convergence and vertical motion (e.g. Benjamin and Carlson 1986; Bluestein et al. 1988). Once convective storms develop, significant downdrafts due to evaporative cooling can contribute to cold pool formation at the surface, regardless of whether or not precipitation actually reaches the ground. Cold pool boundaries have frequently been examined as focal points for CI (e.g. Droegemeier and Wilhelmson 1985). Here we are more interested in quantifying the initial development of convection in a previously undisturbed environment, as subsequent forecasts of new initiation are likely affected by errors in the initial outbreak.

b. Estimates of Surface Variability in the Pre-Convective Environment

Observational studies suggest that variations in boundary-layer fields preceding CI can be very subtle. Weckwerth et al. (1999) estimates that convergence due to only 1-2 m/s variability in winds may promote sustained upward motion. Arnott et al. (2006) observes convergent regions near the surface preceding convective initiation as small as 2-5 km in extent during the International H₂O Project (IHOP; Weckwerth and Parsons 2006). Using IHOP radar refractivity measurements, Fabry (2006) shows that moisture variability in the convective boundary layer (CBL) can be as small as 0.25 kg/m³ over a 5-km length scale. Furthermore, model-based estimates suggest that local variations in moisture of as little as 1 g/kg are enough to support or discourage CI (Crook 1996; Martin and Xue 2006). Simulations suggest that observations on a spatial density of at least the CBL height are necessary to capture horizontal variability (Jonker et al. 1999). Evaluation of these sensitivities and variability in model simulations of CI is critical for evaluating the utility of a potential dense observing network.
c. Scope of This Study

As noted above, it is well-established that large mesoscale forcing (e.g. fronts, drylines, cold pool boundaries) can act as a focal point for CI (e.g. Owen 1966; Ziegler et al. 1997). Dense surface observations have already been shown to improve forecasts of broader mesoscale features (e.g. Wheatley and Stensrud 2009; Madaus et al. 2014) which has, in turn, led to more accurate CI forecasts (Sobash and Stensrud 2015). Therefore, the practical predictability of convective initiation (and the near-surface properties of the pre-convective environment) can immediately be enhanced by more accurate analyses of these features. Here, we wish to examine the surface storm-scale variability in the pre-convective environment absent these effects.

More generally, studies have repeatedly demonstrated that small, storm-scale forecast errors surrounding convective initiation can rapidly propagate upscale, limiting predictability to a matter of hours (e.g. Brooks et al. 1992; Zhang et al. 2003; Hohenegger and Schaar 2007). Though the downscale propagation of small errors on broad meso- and synoptic scales also constrains and ultimately limits CI predictability (e.g. Durran and Weyn 2015; Torn and Romine 2015), for short-term CI forecasts (less than 6 hours) storm-scale errors have critical importance. Therefore, this study numerically simulates idealized isolated convective initiation in the absence of prescribed dynamic forcing to examine how surface observations might constrain these errors, and what quality and quantity of observations might be required.

3. Methods

a. Convective Environments

To examine convective initiation in idealized simulations without imposing a large-scale forcing mechanism, it is necessary to choose appropriate environments. Twenty-one observed soundings
are subjectively chosen to provide environmental conditions in our simulations (Fig. 1). These soundings span the continental United States (US) east of the Rocky Mountains during June–October, 2014. Specifically, synoptic charts, radar data and storm reports are scrutinized to identify times and locations where:

- Precipitating convection developed within 12 hours of the sounding time
- Large-scale forcing (e.g. fronts, upper-level shortwaves) was distant or absent from the event

The soundings at locations shown in Fig. 1 cover a variety of warm-season environments in the US. Figure 2 describes the surface-based and most-unstable Convective Available Potential Energy (CAPE) and Convective Inhibition (CIN) computed from these soundings. CAPE ranges over 0—2500 J/kg and CIN over 0—150 J/kg (for the most unstable parcel). We expect that the environmental diversity captured by these soundings provides an estimate of the variability and broader applicability of the findings.

b. Forecast Model

For idealized simulations, we use Cloud Model 1, Revision 17 (CM1; Bryan and Fritsch 2002) which has enjoyed widespread use in the severe storms community. From each sounding identified in Fig. 1, homogeneous initial conditions are generated for a 104x104km domain. Land surface properties are uniform and match the the United States Geological Survey (USGS) land-use category that is most prevalent within a 100km radius of the sounding site. Boundary conditions are doubly periodic and the only external forcing is diurnal radiation. Simulations are run for 18 hours from the time of the sounding launch.

In many idealized convective simulations, deep convection is “forced” to initiate by the inclusion of some storm-scale perturbation, often a “warm bubble” or prescribed convergence. Here we
wish to examine a free process of convective initiation without prescribing an initial lifting mechanism. In these simulations, we only impose random, white-noise, gridpoint perturbations to the initial potential temperature field with a maximum magnitude of 0.2K to promote random spatial variability. The environments chosen have sufficient conditional instability that deep convection initiates on its own after diurnal heating and convective mixing have developed a sufficiently deep convective boundary layer. In each environment examined, the first convective clouds appear between four and eight hours after the beginning of the simulation and are uncorrelated with the initial noise perturbations. While this “unforced” convective initiation method limits the possible environments that can be evaluated, these simulations allow assessment of the CI process absent a prescribed, forced initiation.

For each sounding environment, a 25-member ensemble is produced at 1km horizontal resolution. Ensemble diversity is solely due to different initial random potential temperature perturbations. There is debate over the horizontal resolution required to adequately simulate deep convection (e.g. Mass et al. 2002; Bryan et al. 2003; VandenBerg et al. 2014). While higher resolution (O(100m)) may be required to adequately simulate boundary-layer convective rolls common to conditionally unstable environments (Weckwerth et al. 2014; Nowotarski et al. 2014), 1km resolution is chosen as it can provide acceptable forecasts of individual convective storms and more closely matches what current and near-future operational forecasting systems may hope to achieve (e.g. Skamarock 2004; Schwartz et al. 2009; VandenBerg et al. 2014). Future studies may wish to examine simulations at multiple resolutions to compare how the scale and magnitude of observed features may change going from realistic, observation-like (higher-resolution) to operational forecasts (lower-resolution).

Vertical resolution is 80 meters in the lowest 3.2km and linearly increases to 500 meters from 9km to the model top at 18km. Rayleigh damping of vertical motions is applied above 15km.
Parameterizations used include the YSU planetary boundary layer (PBL) scheme (Hong et al. 2006) and NASA/Goddard longwave and shortwave radiation (Chou and Suarez 1999, 2001). The NASA Goddard-LFO microphysics scheme is used, as in the CM1 model the radiation scheme is specifically designed to be used with this microphysics scheme.

c. Analysis Methods

1) COMPOSITING ANALYSIS

To examine the typical characteristics of surface fields surrounding the convective initiation process, we employ a compositing method centered on times and locations where convective initiation occurs. To determine these times and locations, we use an object-based method of identifying and tracking cloud and storm objects loosely based on the Method for Object-based Diagnostic Evaluation (MODE, Davis et al. 2006). Specifically, for each ensemble member forecast, the simulated composite reflectivity field (CREF) is scrutinized at five-minute intervals to find regions where $\text{CREF} > 20 \text{ dBZ}$. Within each identified region, all local maxima of CREF are identified and, if there are more than one, a watershed algorithm is applied to separate the region into unique “objects” surrounding each maxima. All objects in this set with a spatial area less than $9 \text{km}^2$ are discarded. The remaining objects are considered to be the set of deep “clouds” present in the model at that time.

These weighted centroids of the CREF cloud objects are tracked through time using the Python TrackPy package (Allan et al. 2014) at five-minute intervals. Only objects successfully tracked for at least 20 minutes are considered candidate storms. At each time, we examine the accumulated surface precipitation field within the boundaries of each cloud object. The cloud object is labeled a “precipitating storm” if the precipitation has increased within the object boundaries since the previous time. We identify times and locations where convective initiation occurs by noting the
time and centroid of any object when and where precipitation first reaches the ground in that object. We find that our definition of CI, based on observed precipitation at the ground and not on an arbitrary threshold of simulated reflectivity, provides greater consistency and reliable detection of storm objects from simulations of different environments.

For our compositing analysis, we select the first four precipitating storm objects that form in each ensemble member. Preliminary analysis showed that storms that form later in the simulations tended to form along the edges of cold pools produced by earlier storms, and the signals of these cold pools dominated our analysis of pre-convective surface features. While cold pools are an important potential focal point for CI, surface observations have already been seen to improve cold pool representation in models (e.g. Wheatley and Stensrud 2009; Dong et al. 2011). As noted in Section 2, here we want to focus on surface features surrounding the initial wave of CI in the model, absent more organized forcing. With 25-member ensembles, selecting the first four precipitating storms in each member gives a sample size of 100 storms in each composite.

To create composites, for each storm the domain is centered on the time and location when CI occurred (by our definition above). Anomalies are computed in the fields of interest at that time by subtracting the domain-wide horizontal mean of those fields. This process is repeated at five-minute intervals for times up to 180 minutes before and 60 minutes after CI, with the domain remaining centered on the location of CI. No accommodation is made for variations in storm motion vectors; for the initial wave of CI the storm motion from the cells examined is generally parallel. The anomaly fields from all storms at each time relative to CI are then averaged together to form a composite. An analysis of these composites follows in Sections 4a and 4b.
2) CORRELATION LENGTH SCALES

When developing an observation network to observe the pre-convective environment, it is important to consider the average spatial extent of the impact for a given observation, as this has implications for the necessary observation density. To investigate this, average correlation length scales for relevant surface variables are computed for each environment’s ensemble simulation. Specifically, for a given variable, at each point in the domain the ensemble correlation of that point to each other point in the doubly-periodic domain (of the same variable) is computed. The correlation magnitude is then binned by distance from that point to the nearest kilometer, and a bin-wise average is computed, again for each point. By averaging this correlation-by-distance profile across all points at each time in the simulation, the spatial scale of correlation at a given time in the simulation may be examined. A discussion of the results of this analysis is given in Section 4c.

In addition to horizontal correlation length scales, average vertical correlations as a function of time are computed for selected variables. These correlations illustrate the potential impact of a surface observation on atmospheric properties aloft. We discuss one such correlation analysis in Section 4d.

4. Results

In the ensemble simulations for all environments, the boundary layer depth is observed to grow throughout the first several hours of integration in response to increasing heat flux from the surface as solar radiation increases. Cloud formation is observed between 4-8 hours into the simulations with the first surface precipitation occurring between 5-10 hours after the simulation begins (not shown).
a. Compositing Analysis

Figure 3 shows an example of the CI-relative composites of composite reflectivity (CREF), 2m temperature (T2), 2m mixing ratio (Q2), 10m v-wind component (V10) and surface pressure (PSFC) for the 8 August 2014 KBMX environment. The composites from this environment exemplify the salient features seen in composites from the other environments. Below is a description of these features for each variable. Discussion of the magnitudes of these features follows in Section 4b.

1) COMPOSITE REFLECTIVITY

The CREF composites in Fig. 3 (top row) illustrate the reliability of the object-based compositing technique. No signal is present until approximately 30-40 minutes before CI (again, here defined as the onset of precipitation), after which a smoothed CREF maximum grows, passing over the location of CI (black square) at the time of CI (0 min., fifth column) before moving off. This affirms that the composite is reliably co-locating simulated storm objects.

2) SURFACE TEMPERATURE

In composite 2m temperature (T2) (Fig. 3, second row), a warm anomaly appears as early as 120 minutes before CI for this particular environment; in other environments the time of the first discernible warm anomaly varies from 90 minutes to 170 minutes prior to CI. This warm anomaly propagates toward the location of CI, attaining a maximum magnitude between 45-60 minutes before CI. The anomaly area greater than 0.1 K is broad, extending some 30 km in its widest dimension, but the area of >0.5 K is very localized—only 2 km in diameter.

Approximately 30 minutes prior to CI, a cold anomaly develops, collocated with the CREF maximum. Examination of the surface heat fluxes and column cloud water (Fig. 4) relates this anomaly
to cloud shadowing: the reduction of incoming shortwave radiation at the surface due to the cloud overhead. In Fig. 4, the composite column-integrated frozen and liquid water (CLOUD; first row) shows a cloud object appearing approximately 40 minutes prior to CI. This cloud object is collocated with a negative shortwave radiation anomaly at the surface (RADSW), reduced upward heat flux at the surface (HFX) and the negative T2 anomaly described above. These anomalies grow and propagate in tandem as the cloud object evolves toward initiation. A cloud shadowed T2 anomaly is observed in all of the simulated environments. We note that since the NASA-Goddard radiation scheme uses a two-stream approximation (where only the vertical component of radiation is evaluated, regardless of the solar zenith angle), this cloud shadow always occurs directly underneath the developing cloud. Further discussion of the cloud shadowing effect follows in the discussion of Fig. 5 and with respect to vertical correlations in Section 4d.

In the 30 minutes prior to CI, the cloud-shadowed cool temperature anomaly grows in size while the warm anomaly reduces in both magnitude and extent. After precipitation begins, the area of the cold anomaly rapidly expands and the magnitude continues to increase. However, at this point the cloud shadowing is augmented by cool downdraft air descending in the now-precipitating storm and this cold anomaly could be better characterized as a cold pool.

3) Surface Moisture

Anomalies in the surface moisture field (Fig. 3; Q2) do not become noticeable until about 60 minutes prior to CI. A distinct dipole structure is apparent during the hour before precipitation begins, with a negative moisture anomaly preceding a positive anomaly. The positive anomaly tends to remain directly underneath the cloud object (compare to CREF in top row). The moisture anomaly is slightly smaller in extent than the temperature anomaly, but the area of > 0.1 g/kg still becomes 20 km across at its widest dimension in the 30 minutes prior to CI. After CI, the
positive moisture anomaly is quickly replaced by an expanding negative anomaly as the cold pool develops. While the air in the developing cold pool is nearly saturated, the saturation mixing ratio in the cold pool is lower than the ambient environmental mixing ratio, yielding a negative anomaly.

4) 10M WINDS

Near-surface winds can be considered in vector form (both $u$- and $v$-wind components) or by examining each component individually; here we select the $v$-wind component ($V_{10}$) as it showed the strongest signal given the background wind profile of the 8 August 2014 KBMX environment and for ease of comparison with the other variables shown. Note that the mean storm motion has not been subtracted from these winds. The expected signature of low-level convergence first becomes apparent by 60 minutes prior to CI. The convergence strengthens through the following 30 minutes, with the $v$-extrema separated by 7 km, on average. Immediately prior to CI, the convergent anomalies weaken and are rapidly replaced by diverging wind anomalies. These diverging winds again represent cold pool generation with the now-precipitating downdraft spreading out as it reaches the surface.

5) SURFACE PRESSURE

The surface pressure (PSFC) anomalies compliment the 10m wind and 2m temperature anomalies in the hour leading up to CI. A negative pressure anomaly is apparent beginning at 60 minutes prior to CI and is visible through 20 minutes prior to CI (not shown), consistent with the concurrent positive temperature anomaly and convergent winds. By 10 minutes prior to CI, the negative anomaly has been replaced with a growing positive anomaly, likely due to the hydrostatic response to a developing downdraft aloft (as pressure responds to the integrated mass above a point) and
the persistent cooling due to cloud shadowing. This positive anomaly grows rapidly following initiation—another symptom of cold pool development.

b. Composite Anomaly Magnitudes

The above descriptions (particularly the magnitudes and scales) apply specifically to the 8 August 2014 KBMX environment, but the analyses are performed on all environments and show similar evolution. Figure 5 shows time series of the maximum and minimum anomaly magnitude in the entire composited domain for the variables described above (excepting CREF) and for each environment simulated.

For T2 (Fig. 5, upper left), the positive anomalies reach a mean magnitude of 0.3 K, which, on an environment-by-environment basis, generally remains larger than the mean negative magnitude until 40-60 minutes prior to CI. During this time, the cloud shadowing effect becomes apparent in all environments, generating increasingly negative temperature anomalies to about 2 K in magnitude. A slight kink in the negative T2 temperature anomalies at the time of CI highlights the transition from cloud-shadowing to cold pool generation as the primary source of the anomalies.

Surface humidity (Q2) has much more variability between environments. A general increase in both the positive and negative anomaly magnitude over time follows the strengthening dipole as shown in Fig. 3. The mean anomaly magnitudes increase from 0.1 to 0.5 g/kg during the two hours prior to CI. In some environments, there are rapid increases in the positive or negative moisture anomaly beginning 30-40 minutes before CI. The maximum positive moisture anomaly during the 30 minutes prior to CI is well-correlated (r=0.66, p<0.01) to the lifted condensation level (LCL) in the original sounding for each environment (not shown). This correlation suggests that in environments with higher LCLs, a larger moisture anomaly is required for successful CI, agreeing with the parcel model. Negative moisture anomalies become more dominant after CI, but
with a great deal of variance as the moisture properties of the cold pool are highly dependent on the environment throughout the depth of the precipitating cloud.

For surface winds (V10), another slow increase in both positive and negative anomaly magnitude is apparent throughout the two hours prior to CI. These positive and negative anomalies as shown here do not necessarily imply convergence by themselves; however, an analysis of each environment’s anomalies (as above) indicates that this is indeed the case. The average positive and negative anomalies are generally symmetric for each environment, reaching an average magnitude of +/- 1 m/s from −40 minutes through the time of CI. A rapid increase in anomaly magnitude after CI (to +/- 4 m/s) shows the strong divergence (again, verified as above) with cold pool development.

Surface pressure anomalies remain very weak until 30 minutes prior to CI, with a mean amplitude of +/- 0.02 hPa. As the deep convective cloud begins developing (−30 to −40 minutes), positive pressure anomalies slowly increase to 0.1 hPa 10 minutes prior to CI. The average positive anomaly sharply increases around the time of CI to 0.5 hPa as the convective downdrafts become established.

c. Correlation Length Scales

To suitably observe the pre-convective environment, it is important to consider the area over which a given observation may be able to provide useful information about the state. With a state that progresses from near homogeneity through cloud formation and eventual precipitation and cold pool generation, it is expected that this area will vary over time. Furthermore, this area may vary among different variables, indicating that particular observation types may be more efficient at providing state information over a larger area with fewer observations.
For each environment simulated, we use the 25-member ensemble for that environment to compute the ensemble correlation magnitudes as they evolve throughout the simulations. Specifically, at every time in the simulations, for T2, Q2, V10 and PSFC we compute the ensemble correlation between each variable at every gridpoint and the same variable at all surrounding gridpoints as a function of distance. This is averaged across all points in the domain to derive an average decorrelation length scale at a given time. We choose a correlation magnitude of 0.2 as a threshold for decorrelation, as estimates of statistical significance using a standard T-test indicated that, on average, this was the largest correlation magnitude that maintained statistical significance with 95% confidence.

Fig. 6 shows the evolution of these de-correlation length scales as a function of time relative to the ensemble-average time of the first CI event in each environment. For T2, Q2 and V10, the mean correlation length scale is 3-4 km throughout the pre-convective period, ranging from 2-7 km depending on the environment. There is a slow decrease in the length scale as the time of CI approaches, reaching a minimum as clouds begin to populate the domains during the 100 minutes before CI. Correlation length scales increase for all three variables after precipitation begins, as the domains begin to be dominated by large, well-defined cold pools. Surface temperature begins to see an increase in correlation length scale about 90 minutes before CI, which appears related to the cloud shadowing effect, separating the domain into one of two regimes (underneath a cloud or not), similar to cold pools.

Surface pressure correlation length scales (Fig. 6, lower right) are generally broader in extent throughout the simulations, with mean length scales between 4-7 km but a range of 2-12 km in the hours prior to CI. Like T2, PSFC also sees a marked increase in correlation length scale in the 100

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1We note that using correlations at all N points in the domain has fewer degrees of freedom (DOF) than the total number of points in the domain, and compute our significance assuming DOF = 0.2N, which is more appropriate given the relatively short correlation length scales we observe
minutes prior to CI which is likely related to the development and deepening of clouds. We note that this coincides with the growing positive pressure anomaly seen in Figs. 3 and 5.

d. Vertical Correlations

Here we briefly consider the ability of surface observations to constrain not just the boundary layer, but features of the free atmosphere. As noted above, the largest surface anomalies we detect—the cold anomalies in T2—appear closely related to the presence of clouds above. This suggests that, absent other larger-scale variability, in the time prior to CI there could exist a strong correlation between surface cold anomalies and water vapor or cloud water above. Fig. 7 shows the domain-averaged ensemble correlation between T2 and the sum of water vapor and cloud water (Qv+Qc) in the column above each point as a function of time for the same 8 August 2012 KBMX ensemble as in Fig. 3. Note that in this ensemble, the average time when precipitation first occurred is at 520 minutes (about 8.6 hours) into the simulation, or about 15:40 local time.

The growth of the convective boundary layer (CBL) is clearly seen as the depth of non-zero correlations increases during the first 360 minutes of the simulations. During this time, horizontal bands of higher (lower) correlation just above the boundary layer correspond reasonably well with layers of higher (lower) relative humidity in the initial sounding profile (Fig. 7, left panel). During the later half of this period, the horizontal bands appear to transition to vertical bands of correlation within the boundary layer as time progresses. As the CBL deepens in the ensemble, the estimated depth of the boundary layer is strongly positively correlated with the temperature at the surface (not shown). Thus, areas of warmer surface temperatures tend to have a locally deeper CBL, which entrains air from a higher height. In the presence of a varying vertical profile of moisture, the humidity of the entrained air will vary depending on the depth of the CBL. In our idealized setting, this illustrates how observations of surface temperature may provide information about
moisture within and just above the CBL, assuming the broader-scale vertical moisture profile is well-constrained.

After 360 minutes, clouds begin forming in the ensemble and the character of the correlation patterns in Fig. 7 changes. At 360-400 minutes, rapidly deepening cumulus clouds appear in the simulations. The area of negative correlation between surface temperature and moisture aloft deepens with cloud development from a base at 4 km up to over 14 km during 400-600 minutes, affirming the connection between a cold, “shadowed” temperature at the surface and higher moisture in a cloud aloft over a relatively deep layer. Below 4 km, cold pools become the dominant feature after the onset of precipitation (here, again, at 520 minutes), showing a positive correlation between colder air at the surface and drier (with a lower saturation mixing ratio than the ambient environment mixing ratio) downdraft air populating the boundary layer. These correlation patterns illustrate the potential for surface observations to inform the moisture field throughout the CI process over a large depth of the troposphere. Further discussion of this finding, including the limitations of these idealized experiments, follows in Section 6.

5. Implications for Surface Observations of CI

a. Comparing observation effectiveness

As noted in the Introduction, there remains an open question about which surface variable is more important for constraining convective initiation, particularly with respect to temperature and moisture. In the results shown in Fig. 5, the largest anomalies relative to the background variance occur in surface temperature. T2 anomalies are also the earliest to appear in composites and the largest in spatial extent more than one hour before CI. Moisture anomaly magnitudes are highly dependent on the environment, suggesting in some cases local moisture variability may be
more important (e.g., when LCL levels are high) for discriminating locations of CI. The average maximum moisture magnitudes seen in the pre-convective environment (0.5 g/kg) are consistent with observational studies (Crook 1996; Fabry 2006; Martin and Xue 2006) that suggest moisture variability of 1 g/kg or less.

A convergent wind signal is apparent in all the composited pre-CI environments, with the maximum average velocity extrema (+/- 1 m/s) consistent with previous observational studies of pre-convective environments (Weckwerth et al. 1999; Arnott et al. 2006). The surface pressure anomalies are extremely weak (< 0.1 hPa), which would require very well-calibrated and sensitive instruments to detect.

b. Observation accuracy and density requirements

To further consider the practical effectiveness of surface observations to capture the anomalies described above, we consider these anomalies in the context of the data assimilation problem. Specifically, we compute the number of independent observations that would be required to produce an analysis with an error variance sufficiently small to resolve an anomaly of a given magnitude with a specified degree of confidence. We employ the Kalman filter equation to update the analysis error variance upon assimilating a single observation. Since the number of observations that will be required is not known a priori, we use the ensemble square–root filter variant, which is formulated for serial observation processing (Whitaker and Hamill 2002). For each single, scalar variable that is directly observed, the analysis error variance is given by:

\[
\sigma_a^2 = [1 - \beta \sigma_b^2 (\sigma_b^2 + r^2)^{-1}] \sigma_b^2
\]  

(1)

Where \(\sigma_a^2\) is the analysis error variance, \(\sigma_b^2\) is the background (prior) error variance, \(r^2\) is the observation error variance and \(\beta\) is the “square-root factor” that accounts for observation sampling.
error and is a function of \( r^2 \) and \( \sigma_b^2 \) (Whitaker and Hamill 2002). For each observation assimilated, the analysis error variance is reduced. We iterate this equation for multiple independent observations with uncorrelated errors until a target variance is reached.

Figure 8 illustrates the objective of this exercise. The variance of a background distribution (“Prior”) is reduced by serial assimilation of observations until the analysis error variance (“Posterior”), which is also the ensemble variance, is small enough to constrain the error in the ensemble mean to within +/- \( n \) with a specified degree of confidence. The number of independent observations required to achieve this is recorded. Here, confidence levels of 70% and 90% will be considered.

There are three free parameters to consider in this analysis: the observation error variance \( r^2 \), the background (prior) error variance \( \sigma_b^2 \) and the magnitude of the anomaly to resolve, \( n \). Figure 9 illustrates the observation counts required to resolve a range of different anomaly magnitudes.

In the left panel, several observation error variances are examined while assuming an initial background error variance (\( \sigma_b^2 \)) of 1 unit\(^2\) in all cases. In the right panel, a range of different initial background error variances are considered while keeping a fixed observation error variance of 1 unit\(^2\).

We note that this analysis assumes that the model background estimate and the observation values are unbiased, which is not true for most model forecasts and observations. Nevertheless, this idealized perspective is still useful for evaluating observation effectiveness even within a “perfect” scenario.

In most operational data assimilation systems, surface observation error variances are prescribed to be about 1 unit\(^2\) for T2 [K], PSFC [hPa], 10m wind [m/s] and Q2 [g/kg] (e.g. Burton 2013; Hu et al. 2013). We consider an observation error variance of 1 unit\(^2\) in the top row of the left panel of Fig. 9 and in the entire right panel. Recall the mean positive pre-convective environ-
ment T2 anomalies of 0.5 K magnitude shown in Fig. 5. Assuming an initial background error variance of 1 K$^2$ and an observation error variance of 1 K$^2$, for variations of 0.5 K to be resolved with 90% confidence, approximately 25 independent observations are required. However, to resolve with only 70% confidence, only one observation is required. By reducing the background error variance to $<0.2$ K$^2$ (right panel), only one observation would be required to resolve this anomaly with 90% confidence. This indicates that it is likely to be correlation length scale and not observation error that would most constrain the number of temperature observations required. For 10m wind anomalies of 1 m/s, a similar conclusion may be drawn.

However, for Q2 and PSFC, the pre-convective anomaly magnitudes are smaller (0.25 g/kg Q2; $<0.01$ hPa PSFC), and with the observation error variance suggested above (1 unit$^2$) a much larger number of independent observations would be required to resolve these magnitudes. In particular, for PSFC, 300 (100) observations would be required to resolve these magnitudes with 90% (70%) confidence (Fig. 9, top rows). Though PSFC has a slightly longer correlation length scale (Fig. 6), an observation error variance of 1 hPa$^2$ is too large for even an extremely dense PSFC observing network to resolve the anomalies estimated from these simulations.

As the initial background error variance is reduced (Fig. 9, right panel), the required number of independent observations becomes more tractable. One way to maintain a small background error variance would be to cycle an ensemble forecast more frequently, not allowing time for ensemble member forecasts to diverge far from observations or from each other. This is one way to possibly make PSFC and Q2 observations more relevant to storm-scale assimilation in the pre-CI environment.

It is also possible that these observation error variances may be overestimates. Hacker and Rostkier-Edelstein (2007) estimate observation error variance from instruments at the Oklahoma Atmospheric Radiation Measurement (ARM) site and suggest temperature and moisture observa-
tion error variances may be as small as 0.08 K^2 and 0.077 g^2/kg^2, respectively. Even barometric pressure measurements from modern smartphone sensors have a specified relative error variance of 0.04 hPa^2 (Mass and Madaus 2014). While sensor error is only one component of the total observation error, it suggests that a network of well-calibrated barometers may be capable of resolving smaller anomalies than presently assumed.

6. Summary and Conclusions

In this study, we examined isolated convective initiation for a variety of different environments absent large-scale forcing with idealized ensemble simulations. A CI-relative compositing method revealed coherent features in all surface fields examined throughout the CI process. Warm 2m temperature anomalies were present from 90 - 170 minutes prior to the onset of precipitation, with moisture, wind and pressure anomalies also detectable about 1 hour before CI. Wind and moisture anomalies were found to be consistent with observational studies of pre-convective environments. The strongest anomalies prior to CI were related to cloud shadowing, creating a substantial cold anomaly of < −0.7 K as the convective cloud above deepens. Surface pressure anomalies remained relatively weak (<0.1 hPa) throughout the CI process.

In general, all surface fields (except surface pressure) showed horizontal correlation length scales of 3-4 km prior to CI, with surface pressure having a slightly longer length scale of 4-7 km. These suggest an upper bound of about one observation every 5 km for a storm-scale observing network, neglecting temporal correlations. Strong vertical correlations were also observed, particularly between surface temperature anomalies and moisture anomalies aloft. This further suggests that surface observations may help constrain storm-scale features throughout the troposphere, given a well-constrained background environment.
The feasibility of resolving pre-CI surface anomalies was examined in the context of variance reduction through data assimilation. We show that the necessary temperature and wind anomalies would most easily be resolved using current assumptions about observation error variance. Capturing moisture and pressure anomalies would require either an extremely dense network, frequent assimilation cycling or observations with much smaller error variances than currently assumed.

One limitation to these results is the use of a single PBL scheme. The surface fields examined here are chiefly outputs of this scheme, and using a different scheme may yield different results. However, Hacker and Rostkier-Edelstein (2007) suggest that the ability of surface observations to constrain the PBL is not dependent on the PBL parameterization used. Future studies may wish to compare these results with other PBL schemes; here we were limited by the single YSU PBL option available in the CM1 model.

In all environments simulated, the same pattern of anomalies is present during the hours preceding convective initiation, suggesting that there are indeed robust surface features that an appropriate observing network could sample. While these results are encouraging, they must be considered within the context of this idealized modeling setup. For instance, though there are strong patterns of vertical correlation between surface temperature and the moisture in the troposphere aloft (Section 4d), much of this correlation strength is likely due to the nearly identical vertical moisture profiles in all ensemble members without synoptic and broad mesoscale variability. It is unlikely that a real-world ensemble would have such a tightly constrained background environment, and that could lessen the correlation magnitude. Indeed, the role of entrainment in promoting or inhibiting CI is largely ignored in this experiment, as clouds forming in all ensemble members encounter similar conditions aloft.

The correlation length scale evolution and the small variance in the anomaly growth time scale is also likely due to all ensemble members producing clouds and eventually precipitation at about the
same time; in a real-world ensemble with greater variance there is likely to be more diversity in the
time and spatial locations where CI is favorable. However, these results suggest a more rigorous
condition for successful storm-scale CI prediction: even with a tightly constrained background
state, the exact location and timing of CI still is subject to the small scale variability seen here.
Finally, we note that this study does not explicitly connect improved resolution of surface fea-
tures to actual improved forecasts of CI. Important questions remain about the translation from
better surface representation to CI forecast improvement, particularly when faced with a more var-
ied background environment and complicating factors such as entrainment. Future work involving
Observing System Simulation Experiments (OSSEs) with these and other simulations is planned to
test hypothesized observation densities as suggested by this research, examining whether resolving
these anomalies actually leads to improved forecasts of CI.

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