

1 **Optimal geometric characterization of forced zonal**
2 **mean tropical precipitation changes.**

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7 **Abstract** The zonal and mean tropical precipitation response to paleoclimate
8 and anthropogenic forcing scenarios ranging from the Last Glacial Maximum
9 (LGM), CO₂ quadrupling (4XCO₂), mid-Holocene, North Atlantic freshwater
10 hosing and volcanic forcing is analyzed in an ensemble of global climate models.
11 Zonally averaged tropical precipitation changes are characterized in terms of
12 three geometric manipulations of the climatological precipitation (hereafter,

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13 modes): meridional shifts, intensifications, and meridional contractions. We
14 employ an optimization procedure that quantifies the magnitude and robust-
15 ness (across different models) of changes in each mode in response to each forc-
16 ing type. Additionally, the fraction of precipitation changes that are explained
17 by the modes –in isolation and combined– is quantified. Shifts are generally
18 less than 1° latitude in magnitude and explain a small fraction ($<10\%$) of
19 tropical precipitation changes. Contractions and intensifications are strongly
20 anti-correlated across all simulations with a robust intensification and con-
21 traction of precipitation under global warming and a robust reduction and
22 expansion under global cooling during the Last Glacial Maximum. The near
23 constant scaling between contractions and intensifications across all simula-
24 tions is used to define a joint contraction/intensification (*CI*) mode of tropical
25 precipitation. The *CI* mode explains nearly 50% of the precipitation change
26 under $4XCO_2$ and LGM forcing by optimizing a single parameter. These re-
27 sults suggest the shifting mode that has been extensively used to interpret
28 paleo-rainfall reconstructions is of limited use for characterizing forced zonal
29 mean precipitation changes and advocates for a reinterpretation of past pre-
30 cipitation changes to account for the *CI* mode.

31 **Keywords** Tropical Precipitation · ITCZ · Paleoclimate · Tropical Contrac-
32 tion

33 **1 Declarations**

34 1.1 Funding

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37 1.2 Conflicts of interest/Competing interests

38 None.

39 1.3 Availability of data and material

40 The model precipitation data used in this study were uploaded to Figshare and
41 published with the following doi (10.6084/m9.figshare.12284282). All CMIP5/PMIP3
42 and CMIP3/PMIP2 climate model data were downloaded from the Earth Sys-
43 tem Grid Federation (ESGF) node hosted by Lawrence Livermore National
44 Laboratory data repositories (<https://esgf-node.llnl.gov/>).

45 1.4 Code availability

46 Not applicable.

47 1.5 Authors contributions

48 AD formulated the optimization procedure, analyzed the model simulations
49 and wrote the manuscript. ARA designed the freshwater hosing experiments,

50 compiled model simulations and edited the manuscript. DSB provided feed-
51 back on model design, edited manuscript and provided feedback on which
52 colors (shades of pink) readers find generally offensive in graphics.

53 **2 Introduction**

54 The concept of a meridional shift of the intertropical convergence zone (ITCZ)
55 is ubiquitous in both the paleoclimate and climate dynamics literature for de-
56 scribing spatial patterns of forced tropical precipitation changes. Here, we
57 define an ITCZ shift as a meridional translation of the climatological trop-
58 ical precipitation pattern. In paleoclimate studies, concurrent increases and
59 decreases of precipitation in different latitude bands within the tropics have
60 commonly been interpreted as a meridional translation (shift) of the ITCZ
61 including: (i) southward ITCZ shifts in response to Heinrich events (*Jacobel*
62 *et al.*, 2016; *Wang et al.*, 2001) characterized by North Atlantic iceberg dis-
63 charge; (ii) southward ITCZ shifts during the Little Ice Age (*Pahnke et al.*,
64 2007; *Haug et al.*, 2001; *Sachs et al.*, 2009) when proxy evidence suggests
65 the Northern Hemisphere was cold relative to the Holocene; (iii) northward
66 ITCZ shifts during the early Holocene (*Haug et al.*, 2001; *Bird et al.*, 2011;
67 *Wang et al.*, 2001; *Sachs et al.*, 2018) when precessional phasing intensified
68 boreal summer insolation; and (iv) southward ITCZ shifts during the Last
69 Glacial Maximum (*Arbuszewski et al.*, 2013; *Koutavas and Lynch-Stieglitz*,
70 2004) when large ice sheets occupied the Northern Hemisphere continents. A
71 meridional translation of a strongly peaked climatological precipitation dis-

72 tribution would result in a dipole of precipitation anomalies. Thus, the con-
73 trasting precipitation changes across different latitude bands discerned from
74 compilations of paleoclimate records are qualitatively consistent with an ITCZ
75 shift (*Arbuszewski et al.*, 2013; *Sachs et al.*, 2009).

76 In idealized dynamical studies (i.e. aquaplanet simulations) forced by hemi-
77 spherically asymmetric heating/cooling, the ITCZ very clearly shifts towards
78 the hemisphere in which the atmosphere is heated more strongly. This ITCZ
79 response is commonly understood in terms of the Hadley cell shifting in order
80 to redistribute the excess energy in the warmed hemisphere to the opposite
81 hemisphere (*Kang et al.*, 2008; *Wei and Bordoni*, 2018; *Frierson and Hwang*,
82 2012). The same zonal mean ITCZ shift in response to hemispherically asym-
83 metric forcing is also seen in comprehensive coupled climate models (*Chiang*
84 *and Bitz*, 2005; *Schneider et al.*, 2014; *Kang*, 2020). However, the spatial pat-
85 tern of forced precipitation changes in models with zonal asymmetries such as
86 geography, ocean circulations and a seasonal cycle is more complicated than
87 a simple shift of the climatological precipitation pattern. ITCZ shifts in com-
88 prehensive models are diagnosed in a myriad of ways, including changes in
89 the “precipitation centroid” (P_{CENT} which is akin to a gravitational center of
90 mass *Frierson and Hwang*, 2012), the hemispheric difference in tropical pre-
91 cipitation (*Hwang et al.*, 2013), the precipitation power-weighted latitude or
92 “moment” (*Adam et al.*, 2016a), or by simply noting the spatial map of pre-
93 cipitation change is characterized by a meridional dipole. These various ITCZ
94 shift metrics have different strengths and weaknesses, but in nearly all cases,

95 little effort has been made to quantify how much of the forced precipitation
96 changes can be explained by an ITCZ shift.

97 Consider the following example of how a tropical precipitation change that
98 does not involve a meridional translation of precipitation might be misdiag-
99 nosed as an “ITCZ shift” starting from a climatological tropical precipitation
100 distribution is bi-modal with peaks north and south of the equator – as is
101 common in many climate model simulations (e.g. the black line in Fig. 1A).
102 Assume that, in response to external forcing, the precipitation peak in the
103 NH is intensified with no other changes in the distribution. This precipitation
104 change would be diagnosed as “northward shift” in the common metrics used in
105 the literature (e.g. P_{CENT} , precipitation weighted moments, hemispheric con-
106 trast of precipitation). However, a meridional translation of the climatological
107 precipitation is a poor characterization of the precipitation change since the
108 locations of the tropical precipitation peaks have not moved. *Adam (2021) re-*
109 *cently demonstrated that the tropical precipitation changes in idealized model*
110 *settings are better characterized by pulsing of the peaks rather than meridional*
111 *translations for states that have bi-modal tropical precipitation peaks.*

112 In this manuscript, we quantify how well a meridional shift in the climato-
113 logical zonal mean precipitation pattern characterizes simulated precipitation
114 changes (ΔP) under a myriad of climate forcings in the following way. The
115 climatological zonal mean tropical precipitation distribution is shifted north
116 and south in search of the resultant pattern that best matches the simulated
117 precipitation in response to forcing; the fraction of ΔP explained by the shift

118 is quantified by how well the optimally shifted distribution matches the forced
119 response. This procedure also provides a new metric for quantifying the mag-
120 nitude of ITCZ shifts (optimized value of the translation in degrees latitude)
121 that can be compared against other commonly used metrics to address how far
122 the ITCZ shifts in model simulations. Additionally, we ask if there are other
123 geometric manipulations of the climatological precipitation (hereafter, modes)
124 that explain a larger fraction of simulated tropical precipitation changes than
125 the shifting mode. Specifically we analyze two additional modes: i) intensifi-
126 cation of the precipitation distribution with no change in shape and ii) merid-
127 ional contraction/expansion of the precipitation distribution about the equa-
128 tor. These three modes – shifting, intensifying and contracting – are optimally
129 fit, in conjunction, to best match the forced changes in precipitation.

130 We consider forced precipitation changes in large ensembles of climate mod-
131 els across a myriad of different climate forcings including: i) CO₂ forcing; ii)
132 Last Glacial Maximum orbital parameters and boundary conditions; iii) mid-
133 Holocene orbital parameters and greenhouse gas concentrations; iv) freshwater
134 hosing of the North Atlantic and v) volcanic forcing. We limit our analyses
135 to annual and zonal mean precipitation changes (hereafter ΔP) in the cur-
136 rent manuscript, to focus on characterizing the modes of tropical precipita-
137 tion changes on the broadest spatial-temporal scale while acknowledging that
138 zonally-homogeneous and seasonal changes of tropical precipitation are robust
139 and comparable in magnitude to zonal, annual mean changes (*Atwood et al.*,

140 2020; Donohoe et al., 2019; Haug et al., 2001; Yan et al., 2015; Braconnot
141 et al., 2012; Liu et al., 2018).

142 This manuscript is organized as follows. The model simulations that we an-
143 alyze and the methodology for optimally characterizing the tropical precipita-
144 tion changes are described in Section 3. The resulting optimal characterization
145 of ΔP in terms of shifts, intensifications and, contractions and the fraction of
146 precipitation changes described by each mode across the different prescribed
147 forcing scenarios are described in Section 4. A summary and discussion follows.

148 **3 Model simulations analyzed and methodology**

149 Here we describe the suite of coupled climate model simulations and forcing
150 scenarios for which ΔP is analyzed and the methodology for optimally char-
151 acterizing ΔP in terms of shifting, intensifying and contracting modes.

152 **3.1 Model simulations analyzed**

153 We make use of model output for idealized anthropogenic (abrupt CO₂ qua-
154 drupling) and paleoclimate (LGM, mid-Holocene) forcing scenarios run across
155 an ensemble of different climate models as part of the Climate Model Inter-
156 comparison Project (Taylor et al., 2012) and Paleoclimate Model Inter-comparison
157 Project (Braconnot et al., 2007a, 2012). Additional volcanic and freshwater
158 forcing simulations are analyzed in a subset of coupled climate models. Atwood
159 et al. (2020) analyzed the zonal structure of tropical precipitation changes in
160 the same suite of model simulations and additional details of the simulations

161 are provided there. 95 sets of forced simulations and companion control sim-
162 ulations are analyzed.

163 **CO₂ quadrupling:** 22 simulations from the abrupt CO₂ quadrupling
164 experiment (hereafter 4XCO₂) of CMIP5 (*Taylor et al.*, 2012) are analyzed
165 in which atmospheric CO₂ concentrations are instantly quadrupled relative to
166 pre-industrial (PI) levels and held fixed thereafter. ΔP is evaluated from the
167 difference between the climatology in years 100-150 after quadrupling in the
168 4XCO₂ simulations and that in the PI control simulation in the same model.

169 **Last Glacial Maximum – LGM:** 13 total LGM simulations are ana-
170 lyzed with 7 from PMIP2 (*Braconnot et al.*, 2007a) and 6 from PMIP3 (*Bra-*
171 *connot et al.*, 2012). LGM simulations are forced by solar insolation from
172 21,000 years before present, greenhouse gas (CO₂, CH₄, N₂O) and aerosol
173 concentrations are set based on ice core data (CO₂ is set to 185 ppm) coast-
174 lines are set to be consistent with a 120m decrease in sea level, and the land ice
175 topography is prescribed as specified in *Kageyama et al.* (2018) which includes
176 the expansive Laurentide ice sheet over North America. ΔP is evaluated from
177 the LGM climatology calculated from years 31-200 after a spin up period.

178 **Mid Holocene:** 12 PIMP3 (*Braconnot et al.*, 2012) mid-Holocene simula-
179 tions are analyzed. Simulations are forced by orbital parameters representative
180 of 6,000 years before present from *Berger* (1978). Most significantly, due to
181 orbital precession, the seasonal cycle of insolation in the Northern Hemisphere
182 was amplified 6,000 years ago while the duration of summer was reduced. We

183 note that in the annual average the sum of all mid-Holocene forcing agents are
184 symmetric about the equator.

185 **Freshwater Hosing:** 18 total simulations of North Atlantic freshwater
186 forcing (aka 'hosing') are analyzed – 16 from the Community Climate System
187 Model version 4 (CCSM4) as described in *Atwood et al. (2020)* and 1 each with
188 Max Planck Institute of Meteorology OM1 model and the United Kingdom
189 Hadley Center CM3 model as described in *Timmermann et al. (2007)*. In each
190 simulation either 1 Sv or 0.1 Sv of freshwater forcing is imposed across the
191 surface of the North Atlantic Ocean (from 50°N-70°N) for 100 years. ΔP is
192 calculated from the climatology averaged over the last 80 years of simulation
193 (it takes approximately 20 years for the Atlantic Meridional Overturning cir-
194 culation to fully respond to the freshwater forcing). Four ensemble members
195 were performed with the default configuration of CCSM4 with 1 Sv meltwater
196 forcing. The same freshwater forcing was also applied to two bias-corrected
197 versions of the model. The bias corrections were: (1) a modification to the
198 orography of Central America (which reduced the biases in the eastern Pacific
199 climatology associated with the poor resolution of Central American topog-
200 raphy; *Baldwin et al., 2021*) and (2) surface heat flux modifications with a
201 cyclostationary seasonal cycle throughout the tropical oceans (30°S to 30°N)
202 to further reduce the bias in the climatological seasonal cycle in SST. In one
203 configuration of the model with three ensemble members, only the Central
204 American topography was modified with no changes in the surface heat fluxes.
205 In a second configuration of the model with four ensemble members, both the

206 topography and surface heat flux corrections were prescribed. A final set of four
207 ensemble members was performed with 0.1 Sv meltwater forcing with both to-
208 pography and surface heat flux corrections prescribed (see *Atwood et al.*, 2020,
209 for more details).

210 **Volcanic Forcing:** The climate response to volcanic forcing is analyzed
211 separately for volcanic events localized in the NH extratropics, SH extratropics
212 and tropics. Volcanic events from two different types of simulations are ana-
213 lyzed: i) volcanic forcing only experiments and ii) last millennium simulations
214 with all climate forcings.

215 The volcanic forcing only simulations consist of 5 CESM4 Last Millennium
216 Ensemble (LME) volcanic-only simulations (with separate composite volcanic
217 responses constructed for eruptions located in the NH, SH and tropics – mak-
218 ing 15 CESM4 composites) and an ensemble of simulations with the Norwe-
219 gian Earth System Model Version 1-M (NorESM) mimicking a high-latitude
220 Northern Hemisphere (NH) summer eruption (the Laki eruption in Iceland;
221 *Pausata et al.*, 2015a,b). In the NorESM simulations, 100 Tg of SO₂ and dust
222 was added to the upper atmosphere over a four month period to simulate
223 the Laki eruption. Eruptions were prescribed at a different starting year of a
224 transient historical simulation (1850-2005) for each of the 48 ensemble mem-
225 bers. 3 composites each consisting of 16 ensemble members are formed to be
226 consistent with the number of eruptions in the last millennium composites
227 (discussed below). The NorESM volcanic forcing runs are compared against

control runs in the same model that were branched from the same transient historical simulation but with no prescribed volcanic forcing.

The response to volcanic forcing is also assessed from selected PMIP3 all-forcing last millennium (LM) transient simulations (CCSM4 and GISS Model E Ensemble Members 122, 125, and 128). In the last millennium simulations, ΔP forced by volcanic eruptions is diagnosed from the composite response in the years after a volcanic event minus that in the closest 5 years preceding the event with minimal volcanic activity. Specifically, for each eruption, the monthly precipitation response was calculated from months with AOD > 0.02 up to 2 years before and after the peak of the eruption and compared to the control climatology, which was taken to be the climatology of the closest 5-year period prior to the onset of the event with minimal volcanic activity (defined as global mean AOD summed over the 5-year period < 0.01). Tropical volcanic events are defined as volcanic events with peak globally-averaged AOD > 0.1 , and extratropical volcanic events are further defined by hemispherically-averaged AOD at least 25% greater in one hemisphere. This procedure results in composites that consist of 9-20 events in each region and simulation (See *Atwood et al.*, 2020, for more details).

3.2 Methodology for defining optimal precipitation shifts, intensifications, and contractions

We characterize the forced ΔP in terms of shifts, intensifications and contractions of the unforced (control) climatological tropical precipitation by op-

250 timally geometric manipulating the control precipitation to best match the
 251 forced precipitation. We begin by evaluating the optimal shift acting in iso-
 252 lation. To search for the optimal shift, the unperturbed (pre-industrial) zonal
 253 mean precipitation, $P_{PI}(\theta)$, is shifted in latitude, θ , both northward and south-
 254 ward by S degrees latitude in increments of 0.02° up to 5° . Mathematically:

$$P_{SHIFT}(\theta) = P_{PI}(\theta - S). \quad (1)$$

255 For each value of S in the search algorithm, the fraction of ΔP explained
 256 by a shift of S ($FRAC_{SHIFT}$) is defined as:

$$FRAC_{SHIFT} = 1 - \frac{\sqrt{\langle (P_{FORCED} - P_{SHIFT})^2 \rangle}}{\sqrt{\langle \Delta P^2 \rangle}} \quad (2)$$

257 where angled brackets, $\langle \rangle$, denote spatial averages over the tropics (defined as
 258 equatorward of 20°). The fraction of ΔP explained by the shift is negatively
 259 proportional to the root mean square (RMS) difference between the shifted
 260 precipitation (P_{SHIFT}) and the target precipitation (P_{FORCED}) normalized
 261 by the root mean square of ΔP . If P_{SHIFT} perfectly matches P_{FORCED} , the
 262 RMS difference is zero and the fraction explained is 1. Whereas for a shift of
 263 zero, $P_{FORCED} - P_{SHIFT,0} = \Delta P$ so the numerator and denominator are equal
 264 and the $FRAC_{SHIFT} = 0$. The optimal shift is chosen such that $FRAC_{SHIFT}$
 265 is maximized or, equivalently, the RMS error between the shifted precipitation
 266 and the target precipitation (P_{FORCED}) is minimized.

267 To demonstrate the optimization algorithm, consider the precipitation re-
 268 sponse to freshwater forcing in the CCSM4 model shown by the green line in

269 Fig. 1B which has enhanced tropical precipitation in the SH tropics relative
 270 to that in the unforced simulation (P_{PI} , black line). The goal is to meridion-
 271 ally translate the pre-industrial precipitation (P_{PI} , black line) to optimally fit
 272 P_{FORCED} (green line). The optimal shift is determined to be -1.7° (with neg-
 273 ative indicating southward) and the optimally shifted precipitation is shown
 274 by the blue curve. $FRAC_{SHIFT} = 0.25$ which indicates that 25% of ΔP is ex-
 275 plained by the shift. By construction, the optimally shifted precipitation has
 276 the same shape as P_{PI} but is meridionally translated; the significant mismatch
 277 between P_{SHIFT} and P_{FORCED} is associated with changes in the amplitudes
 278 of the peaks. The optimally shifted precipitation matches P_{FORCED} signifi-
 279 cantly better than P_{PI} but it is far from a perfect match. We note the optimal
 280 shift is evaluated to within 0.02° latitude which is a finer resolution than the
 281 model output. These sub-grid scale shifts impact precipitation changes on the
 282 coarser model grid.

283 Next we optimally shift and intensify P_{PI} in conjunction to best match
 284 P_{FORCED} to ask: what additional fraction of ΔP can be explained by intro-
 285 ducing an optimal intensification of the climatological precipitation (relative to
 286 the optimal shift only)? Mathematically, the combined shift and intensification
 287 of P_{PI} is

$$P_{SHIFT,INT}(\theta) = I \cdot P_{PI}(\theta - S), \quad (3)$$

288 where I is the (dimensionless) intensification scalar. We calculate $P_{SHIFT,INT}$
 289 for values of I between 0.8 and 1.2 in increments of 0.01 and over the S values

290 previously considered in the shift only case. The optimal values of S and I
 291 are defined as those for which the RMS difference between $P_{SHIFT,INT}$ and
 292 P_{FORCED} is minimized – noting that the optimal value of S for the com-
 293 bined shift and intensification need not equal that for the optimal shift only.
 294 The additional fraction of ΔP explained by the intensification, $FRAC_{INT}$, is
 295 defined as the improvement of RMS error in the fit to P_{FORCED} relative to
 296 that of the optimally shifted only precipitation, normalized by RMS ΔP :

$$FRAC_{INT} = \frac{\sqrt{\langle (P_{FORCED} - P_{SHIFT})^2 \rangle}}{\sqrt{\langle \Delta P^2 \rangle}} - \frac{\sqrt{\langle (P_{FORCED} - P_{SHIFT,INT})^2 \rangle}}{\sqrt{\langle \Delta P^2 \rangle}}. \quad (4)$$

297 A clear example of the how the intensification mode contributes to explaining
 298 ΔP is demonstrated by the optimal fit to the simulated precipitation response
 299 to $4XCO_2$ in the FGOALS model (Fig. 1A). In this case, the shift alone (op-
 300 timal value of 0.5° northward) describes a very small fraction of ΔP (2%)
 301 whereas the combined shift and intensification (optimized $I = 1.16$ intensifi-
 302 cation) explain 39% of ΔP . The difference in fraction of ΔP explained between
 303 these two optimized fits (37%) is defined as $FRAC_{INT}$.

304 Next, we introduce a meridional contraction of P_{PI} in conjunction with
 305 the shifting and intensifying modes using the form

$$P_{SHIFT,INT,CON}(\theta) = I \cdot P_{PI} \left(\frac{1}{C} \cdot (\theta - S) \right), \quad (5)$$

306 where C is the contraction scalar with $C > 1$ corresponding to an expansion
 307 and $C < 1$ corresponding to a contraction. We search over values of C between

0.8 and 1.2 in increments of 0.01 and over the S and I values previously considered. The optimal values of S , I and, C are defined as those for which the RMS difference between $P_{SHIFT,INT,CON}$ and P_{FORCED} is minimized – noting that the optimal value of S and I for the optimal $P_{SHIFT,INT,CON}$ need not equal those defined for $P_{SHIFT,INT}$ and/or P_{SHIFT} . The additional fraction of ΔP explained by the contraction, $FRAC_{CON}$, is defined as the improvement of RMS error in the fit to P_{FORCED} relative to that of the optimally shifted and intensified precipitation ($P_{SHIFT,INT}$) normalized by the RMS ΔP :

$$FRAC_{CON} = \frac{\sqrt{\langle (P_{FORCED} - P_{SHIFT,INT})^2 \rangle}}{\sqrt{\langle \Delta P^2 \rangle}} - \frac{\sqrt{\langle (P_{FORCED} - P_{SHIFT,INT,CON})^2 \rangle}}{\sqrt{\langle \Delta P^2 \rangle}} \quad (6)$$

317 .

The improved fraction of ΔP explained due to introducing an optimal contraction can be visualized in Fig. 1A by the mismatch between P_{FORCED} (green) and $P_{SHIFT,INT,CON}$ (orange) as compared to that of $P_{SHIFT,INT}$ (red). In the example of the simulated precipitation response to $4XCO_2$ in the FGOALS model (Fig. 1A), the intensity of the bi-modal peaks (near $7^\circ N$ and $7^\circ S$) in the optimized $P_{SHIFT,INT}$ (before the contraction mode is introduced) is too small whereas the intensity of the precipitation on the poleward flanks of the peaks is too large and, thus, further increasing the I scalar in $P_{SHIFT,INT}$ to match the peaks height increases the RMS mismatch with P_{FORCED} . However, introducing a contraction in conjunction with a shift

328 and intensification of the precipitation distribution allows the optimal solu-
 329 tion to match both the increased precipitation in the peaks of the distribution
 330 and the nearly unchanged precipitation on the flanks. The combination of all
 331 three optimized modes ($P_{SHIFT,INT,CON}$) explains an impressive 75.4% of
 332 ΔP . This improved fit relative to $P_{SHIFT,INT}$, which only explains 39.2% of
 333 ΔP , is attributed to the contraction mode ($FRAC_{CON} = 36.2\% = 75.4\% -$
 334 39.2%).

335 We have chosen to apply the precipitation modes in the order: (1) shift,
 336 (2) intensification, (3) contraction. This methodology answers the following
 337 sequence of questions: i) What fraction of ΔP is described by an optimal shift
 338 alone? ii) How much additional ΔP can be described by adding an optimal
 339 intensification in conjunction with the shift? and iii) How much additional ΔP
 340 can be described by introducing an optimal contraction in addition to the shift
 341 and intensification? Because the three precipitation modes are not generally
 342 spatially orthogonal, the order of applying the modes does impact the resul-
 343 tant $FRAC_{SHIFT}$, $FRAC_{INT}$ and $FRAC_{CON}$. In practice, the shifting mode
 344 is nearly spatially orthogonal to the other modes whereas the contraction and
 345 intensification are not. This characteristic can be understood as follows; since
 346 precipitation is larger in the deep tropics (i.e. near the equator) than at the
 347 edges of the tropical domain, an expansion of P_{PI} ($C > 1$) results in a domain
 348 wide intensification and, thus, projects onto the intensifying mode. Qualita-
 349 tively similar results for $FRAC_{SHIFT}$, $FRAC_{INT}$ and, $FRAC_{CON}$ are found
 350 for the five alternative permutations of the order of modes aside from one ex-

351 ception: the value of C found for optimizing ΔP is significantly impacted by
352 whether contraction is introduced before or after the intensification mode due
353 to the previously mentioned increase in tropical domain average precipitation
354 that results from meridional expansion alone. For this reason, we report the
355 modal parameters (S , I and, C) that result from optimizing all three modes si-
356 multaneously. We report FRAC_{SHIFT} , FRAC_{INT} and, FRAC_{INT} that result
357 from the specific order of modes outlined above.

358 4 Results

359 We begin by characterizing the ensemble mean and spread of ΔP under each
360 forcing scenario in terms of optimal shifts, intensifications and contractions
361 of unperturbed precipitation (P_{PI}). In Section 4.2 we then generalize across
362 experiment type to assess the relative importance of shifts, intensifications
363 and contraction for describing tropical precipitation changes. In Section 4.3
364 we compare the magnitude of optimal shifts (S) to other common metrics of
365 ITCZ shifts diagnosed from the same simulations and discuss the differences.

366 4.1 Characterization of tropical precipitation changes in response to different 367 forcings

368 4.1.1 Last Glacial Maximum and $4XCO_2$

369 The tropical precipitation changes from the cold LGM climate (solid blue
370 line in Fig. 2A) to the pre-industrial (black line) and from the pre-industrial

371 to the warm 4XCO₂ climate (solid red line) are best characterized as an in-
 372 tensification and meridional contraction of the precipitation distribution as
 373 the planet warms. The precipitation contraction with warming is visually ev-
 374 ident from the equatorward migration of the bi-modal precipitation peaks
 375 in warmer climates, especially the SH peak under 4XCO₂ forcing. The opti-
 376 mally shifted, intensified and contracted fit ($P_{SHIFT,INT,CON}$) to P_{FORCED}
 377 is shown by the dashed lines in Fig. 2A, which captures a large fraction of
 378 ΔP ; $P_{SHIFT,INT,CON}$ explains 62% and 61% of the ensemble mean ΔP for
 379 the LGM and 4XCO₂ simulations, respectively. The optimal modal fit of the
 380 LGM precipitation ($P_{SHIFT,INT,CON}$) features a robust decrease in the inten-
 381 sity of tropical precipitation across models ($I = 0.91 \pm 0.02$) and an expansion
 382 of the precipitation distribution ($C = 1.04 \pm 0.02$) where the stated range is
 383 two standard deviations of the ensemble mean, roughly corresponding to the
 384 95% confidence interval (values are shown in the table inset of Fig. 2A). In con-
 385 trast, in response to 4XCO₂ forcing there is robust intensification of tropical
 386 precipitation ($I = 1.13 \pm 0.02$) and a contraction in the precipitation distribu-
 387 tion ($C = 0.89 \pm 0.02$). There is also a modest but robust southward shift in
 388 precipitation under LGM forcing ($S = -0.4^\circ \pm 0.2^\circ$), whereas the shift under
 389 4XCO₂ is not robust – a southward shift is found in 9 models and northward
 390 shift is found in 13 models, in agreement with the analysis of *Donohoe and*
 391 *Voigt (2015)*.

392 The optimal shift under LGM forcing is southward in 11 of the 13 simula-
 393 tions (abscissa in Fig. 3A) with an ensemble mean value of -0.4° and is $< 1^\circ$

394 in the majority of models. The magnitude of these shifts is almost an order
395 of magnitude less than those inferred at regional scales from paleoclimate re-
396 constructions such as the 7° southward shift of the Atlantic ITCZ suggested
397 by *Arbuszewski et al.* (2013). Similarly, the magnitude of the optimal shift
398 in response to $4XCO_2$ is less than 1° in all ensemble members (red dots in
399 Fig. 3A). In addition to the small magnitude of the optimal shift, the shifting
400 mode explains a small fraction of ΔP under both LGM and $4XCO_2$ forcing;
401 the ensemble average $FRAC_{SHIFT}$ is 4.1% and 2.4% for the LGM and $4XCO_2$
402 simulations and exceeds 10% in only two of the LGM simulations and one of
403 the $4XCO_2$ simulations (Fig. 4A).

404 In addition to the robust ensemble mean intensification and contraction
405 of tropical precipitation under global warming, the intensification and con-
406 traction are highly (negatively) correlated across model simulations, with cor-
407 relation coefficients of -0.95 and -0.52 for the LGM and $4XCO_2$ simulations,
408 respectively (Fig. 3B). This result suggests that the degree of contraction and
409 the degree of intensification of the tropical precipitation distribution in re-
410 sponse to forcing go hand and hand. We return to this point in Section 4.2.
411 The intensification and contraction modes each explain a significantly greater
412 fraction of ΔP than the shift mode in response to LGM and $4XCO_2$ forc-
413 ings (Fig. 4A). The intensification mode explains 48% and 30% of ΔP in the
414 LGM and $4XCO_2$ simulations, respectively, with values exceeding 10% in all
415 of the LGM simulations and 19 of 22 of the $4XCO_2$ simulations (Fig. 4A).
416 The contraction mode explains 10% and 29% of ΔP in the LGM and $4XCO_2$

417 simulations, respectively, with values exceeding 10% in 7 of 13 of the LGM
418 simulations and 20 of 22 of the 4XCO₂ simulations (Fig. 4A). These results
419 collectively demonstrate that contractions and intensifications of tropical pre-
420 cipitation respond robustly to CO₂ and LGM forcing and better characterize
421 the changes in tropical precipitation changes than does meridional shift.

422 4.1.2 Freshwater hosing

423 The ensemble mean tropical precipitation under 1 Sv of freshwater hosing of
424 the North Atlantic (solid orange in Fig. 2B) features a reduction in the peak
425 amplitude in the NH and an amplification of the precipitation maximum in
426 the SH relative to precipitation in the PI simulation (black solid line). ΔP
427 in response to freshwater hosing is best characterized by the shifting mode
428 (ensemble mean FRAC_{SHIFT} of 20% — Fig. 4A) with a robust ensemble
429 mean southward shift of $-1.1^\circ \pm 0.3^\circ$ (Fig. 2B). ΔP in response to 0.1 Sv
430 of freshwater forcing (4 CCSM4 simulations— not shown) is much smaller in
431 magnitude than the response to 1 Sv of freshwater forcing with a similar
432 spatial characterization with an amplification of the SH precipitation peak
433 and reduction of the NH peak. The ensemble mean ΔP response to 0.1 Sv
434 is not significantly different than 10% of the ensemble mean ΔP in response
435 to 1 Sv of hosing in the CCSM4 (12 simulations) with an ensemble mean
436 (southward) optimal shift (S) of $-0.12 \pm 0.03^\circ$ in response to 0.1 Sv as compared
437 to $-1.1 \pm 0.3^\circ$ in response to 1 Sv. This result indicates that the number of
438 simulations considered can not reject the null hypothesis that the ΔP response

439 to hosing is linear with respect to the magnitude of the prescribed freshwater
440 flux.

441 In the 1 Sv freshwater hosing simulations, there is also a modest but ro-
442 bust reduction of the ensemble mean precipitation ($I = 0.97 \pm 0.01$) that
443 explains a small fraction of ΔP (1%), whereas the ensemble mean contraction
444 is not significantly different from zero. Despite the ambiguity of the direc-
445 tion of contraction across ensemble members (Fig. 3B), the contracting mode
446 describes a non-negligible fraction of ΔP (4%) across the ensemble members
447 and up to 22% in individual ensemble members (Fig. 4A). We note that, even
448 though the shifting mode explains a larger fraction ΔP than the other two
449 modes, overall the ΔP is poorly characterized (26%) by the geometry of the
450 optimal modes considered here, as is evident in the mismatch between the
451 optimal fit and P_{FORCED} (c.f. the solid and dashed orange lines in Fig. 2B).
452 The reason for this mismatch is as follows: the three modes chosen in this
453 study do not allow for the relative magnitudes of the precipitation peaks in
454 the NH and SH to change but, rather, only allow the peaks to simultaneously
455 increase/decrease (via the intensity mode), translate together (via the shift
456 mode) or move closer/farther apart (via the contraction mode). If we were,
457 instead, to employ a precipitation mode that optimized the intensity of the
458 precipitation differentially in the NH and SH, the optimized fit would better
459 match P_{FORCED} under freshwater forcing. We return to this point in Section
460 4.3.

461 4.1.3 Mid-Holocene

462 The ensemble average ΔP in response to mid-Holocene forcing is small in mag-
463 nitude compared that of other forcings considered here (green line in Fig. 2B).
464 Furthermore, there is no robust direction of the optimal shift, intensification,
465 and contraction across the ensemble of experiments. For individual ensemble
466 members, the magnitude of simulated ΔP in response to mid-Holocene forc-
467 ing is significantly smaller (approximately 20% of) that in the other forcings
468 considered here with an ensemble average tropical domain RMS of ΔP value
469 of 0.08 mm/day as compared to 0.37, 0.40 and 0.34 mm/day in the LGM,
470 4XCO₂ and hosing simulations, respectively. Additionally, Fig. 4A shows that
471 the total fraction of ΔP explained by all three optimal modes is significantly
472 smaller for the mid-Holocene simulations (ensemble average of 25%) than for
473 the LGM and 4xCO₂ simulations (ensemble averages of 61% and 61%, respec-
474 tively). The contraction mode describes the largest fraction of ΔP across the
475 mid-Holocene ensemble (12%) followed by the shifting mode (11%) whereas
476 the intensification mode is relatively inefficient at characterizing ΔP (2%).
477 We note that there are robust seasonally and longitudinally inhomogeneous
478 changes in simulated tropical precipitation under mid-Holocene forcing (*Bra-*
479 *connot et al.*, 2007b; *Liu et al.*, 2018) but the imprint of these changes on the
480 zonal and annual mean ΔP is small and poorly characterized by the modes of
481 tropical precipitation considered here.

4.1.4 Volcanic forcing

The tropical precipitation response to volcanic forcing is separated into ensemble averages of NH volcanoes (red line - Fig. 2C), SH volcanoes (blue line) and, tropical volcanoes (green line). The ordinate scale for each ensemble average (NH, SH, and tropical) are offset for visual purposes. The ensemble mean P_{FORCED} in response to NH volcanic events features a reduction in the amplitude of the NH precipitation peak and an enhancement of the SH precipitation peak relative to P_{PI} . This precipitation response is qualitatively similar to that seen under North Atlantic freshwater hosing but with reduced magnitude. The response to SH volcanoes is nearly a mirror image of the response to NH volcanoes but with an additional tropical domain average precipitation reduction such that a reduction in the amplitude of the SH precipitation peak is the primary feature of ΔP . ΔP under extratropical volcanic forcing is robustly characterized by an ensemble mean shift away from the source of volcanic forcing: the optimal shift in response to NH volcanic forcing is southward in all 15 ensemble members (ensemble mean $-0.4^\circ \pm 0.2^\circ$) and the optimal shift in response to SH volcanic forcing is northward in all nine ensemble members (ensemble mean $0.4^\circ \pm 0.2^\circ$). In addition to the shift under extratropical volcanic forcing, there is also a robust de-intensification of precipitation under SH volcanic forcing and a contraction of precipitation under NH forcing. These two robust ensemble mean responses (de-intensification under SH forcing, contraction under NH forcing) are most pronounced in the three GISS Model E members that have a larger magnitude volcanic forcing

505 than the other simulations (see *Atwood et al.*, 2020, for a discussion) whereas
506 the shifting response to extratropical volcanoes is robust across models (Fig.
507 4A).

508 Although the intensifying and contracting modes explain a small fraction of
509 the ensemble mean ΔP in response to extratropical volcanic forcing, the inten-
510 sification and contracting modes do explain a significant fraction of ΔP within
511 individual models. The intensifying and contracting modes explain 19% (4%)
512 and 10% (13 %) of ΔP across the SH (NH) volcanic forcing (Fig. 4A). The
513 fraction of ΔP explained by the shifting mode across the ensemble members
514 is 18% for the NH volcanic events and 12% for SH volcanic events. This result
515 suggests that, **while the tropical precipitation clearly redistributes precipita-**
516 **tion away from the cooled hemisphere in response to extratropical volcanic**
517 **forcing**, the precipitation response to extratropical volcanic forcing is poorly
518 characterized by a meridional translation of the climatological precipitation.
519 As was discussed for ΔP under freshwater hosing, the shifting mode cannot
520 capture the opposing changes in the magnitude of the NH and SH precipitation
521 peaks and, thus, fails to capture a large fraction of ΔP .

522 The ensemble mean P_{FORCED} in response to tropical volcanoes features
523 a nearly latitudinally homogeneous reduction in precipitation throughout the
524 tropical domain with no discernible change in the shape of the distribution
525 (green line in Fig. 2C). The optimal characterization of the ensemble mean
526 ΔP under tropical volcanic forcing is a robust reduction in intensity (6% re-
527 duction in the intensity scalar, I) with no statistically significant contraction

528 or shift. Across the ensemble members, the intensification mode explains 38%
529 of ΔP whereas the contraction explains 19% despite the differing direction of
530 contraction/expansion across the the ensemble members (contracting in 5 of
531 9 members and expanding in 4 of 9 members). The shifting mode explains
532 very little (2%) of ΔP in response to tropical forcing which suggests that,
533 in the particular case of tropical volcanic forcing, changes in energy input to
534 the tropical atmosphere (*Schneider et al.*, 2014) has minimal impact on the
535 ITCZ location perhaps because the annual mean (mean state) ITCZ is near
536 the equator and, thus, there is minimal area between the energy flux equator
537 and geographic equator. Previous work has demonstrated that the precipita-
538 tion response to volcanic forcing is “El Nino like ” for eruptions in the NH
539 (*Stevenson et al.*, 2016) and it is possible that the methodology for defining
540 the precipitation response to volcanic forcing used here includes a portion of
541 the transient El Nino response. However, the tropical precipitation response
542 to El Nino in observations is primarily characterized by a contraction of pre-
543 cipitation with a very modest shift (not shown) and, thus, it is unlikely that
544 the transient ENSO response to volcanic forcing contaminates the diagnostics
545 (or explanatory power) of the shift mode. Our general finding that the tropical
546 precipitation is redistributed away from extratropical volcanoes and is reduced
547 in response to tropical volcanoes agrees with the conclusions of *Colose et al.*
548 (2016).

549 4.2 The emergence of a joint contracting and intensifying mode of tropical
550 precipitation

551 In response to each of the forcings, the optimal contraction (C) and intensifica-
552 tion (I) scalars are remarkably anti-correlated ($R=-0.91$, Fig. 3B) with a least
553 squares regression slope of $\frac{C}{I} = -0.8$ (the dashed black line in Fig. 3B). The
554 direction of the response in C and I is remarkably consistent across models in
555 response to the same forcing. In particular, tropical precipitation expands and
556 de-intensifies under LGM forcing (blue dots in Fig. 3B are in the $C>1$ and
557 $I<1$ quadrant) and expands and intensifies under $4XCO_2$ (red dots in Fig. 3B
558 are in the $C<1$ and $I>1$ quadrant). These results suggest that tropical precip-
559 itation contractions and intensifications are tightly coupled and the direction
560 of these forced changes are consistent across models, though the magnitude
561 (of the combined contraction/intensification) is model dependent. Moreover,
562 the scaling between the contraction (C) and intensification (I) scalars seems
563 relatively consistent across models and climate forcings .

564 *4.2.1 Physical basis of the contracting and intensifying mode of tropical*
565 *precipitation*

566 We first evaluate the hypothesis that the tight relationship between C and I
567 under climate change is a result of thermodynamic considerations. Following
568 *Held and Soden (2006)*, if atmospheric circulation and relative humidity are
569 climate state invariant, then the change in precipitation minus evaporation,
570 $\Delta(P - E)$, is given by:

$$\Delta(P - E) = (P - E) \alpha \Delta T \quad (7)$$

571 where α is the linearized dependence of saturation vapor pressure on temper-
 572 ature in the Clausius Clapeyron equation and ΔT is the local temperature
 573 change. Hence, provided that the ΔT has a weaker spatial structure in the
 574 tropics than the climatological $(P-E)$, thermodynamic constraints suggest that
 575 the spatial structure of $(P-E)$ is climate state invariant but intensifies with
 576 warming and de-intensifies with cooling. Thus, thermodynamic constraints are
 577 equivalent to $\Delta(P-E)$ being explained entirely by a geometric mode akin to
 578 the intensification mode but acting on the climatological spatial distribution
 579 of $(P-E)$, with no changes in the shape of $(P-E)$.

580 The thermodynamic prediction for $(P-E)$ under $4XCO_2$ with scaling cho-
 581 sen to match the ensemble mean global integral of the absolute value of $(P-E)$
 582 is shown by the orange lines in Fig. 5. Under $4XCO_2$, the change in E is
 583 nearly spatially homogeneous over the tropical domain (c.f. the solid red and
 584 dashed red line in Fig. 5A) and, thus, the changes in $(P-E)$ mimic the spa-
 585 tial structure of ΔP . Importantly, the shape of $\Delta(P-E)$ differs significantly
 586 from the climatological $(P-E)$ and, thus, the change in $(P-E)$ predicted from
 587 thermodynamic constraints (orange line in Fig. 5B) is a poor match to that
 588 realized in model simulations (black lines in Fig. 5B). Specifically, the latitude
 589 of zero $(P-E)$ robustly contracts toward the equator under $4XCO_2$ and the
 590 bi-modal peaks in $(P-E)$ also shift equatorward. Thermodynamic constraints
 591 explain only 25% of the ensemble mean $\Delta(P-E)$ under $4XCO_2$ and 34% in

592 the LGM simulations (not shown); similar percentages are found in individual
593 models. This is roughly half of the fraction of the ΔP that is explained by a
594 combined mode of contraction and intensification (53% for 4XCO₂ and 48%
595 for the LGM) described in Section 4.3.

596 We next evaluate if the relationship between C and I is a consequence
597 of conservation of tropical domain average precipitation. Given that C and I
598 both change the tropical domain average precipitation (which increases with
599 both increasing I and decreasing C), one would expect an anticorrelation
600 between C and I if the tropical domain average precipitation was nearly un-
601 changed. We demonstrate that the linear relationship between C and I seen
602 in Fig. 3B ($\frac{C}{I} = -0.8$) is significantly different than that expected from conser-
603 vation of tropical precipitation through the following exercise: (i) P_{PI} in each
604 model is contracted/expanded (via random values of C in equation 5), (ii) we
605 find the value of I for each C that restores the tropical domain average of
606 the contracted/expanded precipitation to that of P_{PI} . The resultant tropical
607 domain average conserving relationship between C and I is $\frac{C}{I} = -2.0$ which
608 differs from that seen across the climate forcing simulations. Therefore, the
609 tight correlation between C and I has a physical basis and is not merely a
610 consequence of conservation of domain average precipitation.

611 *4.2.2 Explanatory power of the joint contracting and intensifying mode of*
 612 *tropical precipitation*

613 Given C and I are highly (anti-) correlated across the ensemble of climate
 614 forcing simulations, we now define a combined mode of contraction and inten-
 615 sification constrained by the linear best fit relationship of $\frac{C}{I}=-0.8$. In other
 616 words, we require that C and I used to define the optimal fit to P_{FORCED}
 617 in Eq. 5 follow $=-0.8 I$, thus, reducing the number of free parameters used
 618 to optimize the fit to P_{FORCED} . We refer to this combined mode of con-
 619 traction/intensification or expansion/deintensification as the CI mode. We
 620 repeat our optimal characterization ΔP using two modes – the shift and CI
 621 mode – as opposed to the three original independent modes (shift, contraction
 622 and, intensification). Since we are fitting for two free parameters instead of
 623 three free parameters using the same geometric basis, a smaller fraction of
 624 ΔP will be described by the two mode methodology as compared to the three
 625 mode methodology. However, the advantage of this approach is it simplifies the
 626 characterization of tropical precipitation changes and reduces the likelihood
 627 of statistical artifacts.

628 The fraction of ΔP explained by the two mode characterization across all
 629 forcing scenarios is shown in Fig. 4B alongside that of the three mode charac-
 630 terization in Fig. 4A. We note, that since the shift mode is optimized to best fit
 631 P_{FORCED} before optimizing the other modes in the calculation, $FRAC_{SHIFT}$
 632 is unchanged between the three mode and two mode characterizations. In the
 633 LGM and $4XCO_2$ simulations, the CI mode explains a very large fraction

634 (48% and 53%, respectively) of ΔP , respectively. There is a modest reduc-
635 tion (8.7% in the LGM simulations and on 5.4% in the 4XCO₂ simulations)
636 in ΔP explained due to switching from a three parameter characterization to
637 the two parameter characterization (c.f. panel 4A and 4B) suggesting that the
638 *CI* mode efficiently represents ΔP across these simulations (and across the
639 different models) with little loss of explanatory power. Additionally, the pre-
640 cipitation response to tropical volcanoes, and to some extent SH volcanoes, is
641 well characterized by the *CI* mode (explaining 36% and 28% of ΔP respec-
642 tively). In contrast, in response to freshwater hosing and NH volcanoes, the
643 *CI* mode explains very little of ΔP , as is expected given that ΔP for these
644 simulations is dominated by the shift mode. Under mid-Holocene forcing, com-
645 bining the contraction and intensification modes into a single *CI* mode results
646 in reduction of ΔP explained by 8%, from 26% in the three mode character-
647 ization to 18% for the two mode characterization. Similarly, under volcanic
648 forcing the fraction of ΔP explained by the two mode optimization (32%) is
649 reduced by 11% relative to that of the three mode optimization (43%). These
650 significant fractional loss of explanatory power by the two mode optimization
651 results from the optimal contraction and intensification scalars in individual
652 ensemble members diverging from the ensemble mean best fit line (Fig. 3B).

653 Overall, the *CI* mode is an efficient representation of ΔP in the simulations
654 characterized by global warming and cooling (e.g. LGM and 4XCO₂) and only
655 modestly reduces the explanatory power of geometric modes for describing ΔP
656 under mid-Holocene, freshwater hosing and volcanic forcing.

657 4.3 Magnitude and explanatory power of the shifting mode

658 The optimal shift (S) defined in this study provides a metric to quantify an
659 ITCZ shift in terms of its most literal definition, that is, a meridional trans-
660 lation of the climatological precipitation pattern. This metric can be com-
661 pared to other commonly used metrics of “ITCZ shifts” such as changes in
662 the precipitation centroid – P_{CENT} , defined as the latitude that demarcates
663 regions of equal area weighted precipitation between 20°S and 20°N (*Frierson*
664 *and Hwang*, 2012) – or changes in the precipitation weighted latitude (“mo-
665 ments”) of the tropical precipitation distribution (*Adam et al.*, 2016a). We
666 present the relationship between S and the change in P_{CENT} in Fig. 3A due
667 to the prominence of the latter metric in the paleoclimate (*McGee et al.*, 2014)
668 and dynamics (*Donohoe et al.*, 2013) literature. S and the change in P_{CENT}
669 are very well correlated ($R=0.93$) across the ensemble of different forcings but
670 the magnitude of S is approximately half of the P_{CENT} change (the slope of
671 the linear best fit in Fig. 3A is 0.52) indicating that the P_{CENT} metric of
672 “ITCZ shifts” is more sensitive to the patterns of ΔP realized in model sim-
673 ulations than the optimal shift metric S . Changes in S are generally small in
674 magnitude (ordinate of Fig. 3B): S is less than 1° in all 4XCO₂, mid-Holocene
675 and volcanic forcing simulations and only exceeds 1° in magnitude for two of
676 the 13 LGM simulations. The largest magnitude S values are found in response
677 to freshwater hosing with magnitudes exceeding 1° in 15 of 18 simulations. S
678 is less than 2° in magnitude across all simulations.

679 It is tempting to conclude from the results above that the optimal shift
680 (S) is a less sensitive metric for measuring shifts tropical precipitation than
681 P_{CENT} . We demonstrate below that this is the wrong interpretation and, in
682 fact, for the case of a ΔP that is well characterized by a meridional translation
683 of P_{PI} , S is a *more sensitive* metric of ITCZ shifts than P_{CENT} . To see
684 this point, we first analyze the quantitative relationship between P_{CENT} and
685 the optimal shift S for the case of a ΔP that is perfectly described by a
686 meridional translation of the climatological precipitation in the following set
687 of analysis: (i) the pre-industrial precipitation (P_{PI}) is meridionally translated
688 via Eq. 1 north and south in increments of 0.1° up to 5°N/S , (ii) for each
689 prescribed S , we calculate the resulting change in P_{CENT} relative to that
690 calculated from P_{PI} . We find that the change in P_{CENT} is approximately one
691 half of prescribed S for all cases (regression coefficient = 0.51). This suggest
692 that S is a more sensitive metric for meridional translations of the ITCZ
693 than the P_{CENT} change metric which can be understood as follows. For the
694 case of a prescribed 5° northward translation of P_{PI} , the precipitation that
695 was originally between 15°N and 20°N has shifted out of the tropical domain
696 defined by P_{CENT} whereas the precipitation originally between 25°S and 20°S
697 has entered the tropical domain. P_{CENT} shifts less than the prescribed 5°
698 northward translation because, using the gravitational analogy, it now feels
699 the gravitational pull of the precipitation that has entered the tropical domain
700 from the southern edge (and has lost the gravitational pull of the precipitation
701 that has exited the tropical domain to the North). P_{CENT} can only change as

702 much as S if the precipitation is zero at the edge of the tropical domain. More
703 generally, for the case of a pure translation of the climatological precipitation,
704 changes in P_{CENT} will be smaller than S by a factor of 1 minus the ratio of
705 precipitation at the edge of the tropical domain to that at location of maximum
706 precipitation.

707 Because precipitation is non-zero on the edge of the tropical domain, one
708 would expect S values to exceed P_{CENT} changes *if ΔP was well character-*
709 *ized by a meridional translation of P_{PI}* which is the opposite of what is seen
710 in our results (Fig. 3A). Instead, our results suggest that S values are less
711 than changes in P_{CENT} because the ΔP changes are so poorly described by
712 meridional translations of the climatological precipitation. Instead, the P_{CENT}
713 changes primarily reflect opposing changes in the amplitude of the NH and SH
714 tropical precipitation peaks (e.g. ΔP under freshwater hosing in Fig. 2B) *as*
715 *was recently demonstrated by Adam (2021)*. Such opposing changes in the NH
716 and SH precipitation peaks strongly impact P_{CENT} but only weakly project
717 on to S . *We note the magnitude of S is also smaller than the precipitation*
718 *power-weighted latitude (Adam et al., 2016a) – by a factor of 6 when using a*
719 *power weighting of $N=10$, not shown – a metric that responds (more sensitively*
720 *for higher N) to the relative magnitude of the bi-modal tropical precipitation*
721 *maxima.*

722 5 Summary and discussion

723 The shifting mode (S) of (annual and zonal mean) tropical precipitation
724 changes (ΔP) – defined here as the optimal meridional translation of the un-
725 perturbed precipitation – is found to be relatively useless for characterizing
726 simulated zonal mean precipitation changes under a myriad of anthropogenic
727 and paleoclimatic forcings. In response to global scale warming and cooling
728 associated with LGM boundary conditions and CO₂ quadrupling, the fraction
729 of ΔP described by S (FRAC_{SHIFT}) is 4% and 2% respectively. Even in the
730 cases of strong hemispherically-asymmetric forcings such as freshwater hos-
731 ing and NH volcanic forcing where FRAC_{SHIFT} (19% and 17% respectively)
732 exceeds FRAC_{INT} and FRAC_{CON} , ΔP is poorly characterized by the shift
733 mode. In these cases, ΔP is also poorly characterized by the optimal fit of
734 all tropical precipitation modes considered here with only 25% and 31% of
735 the total ΔP explained as compared to 59% for all other simulations. Hemi-
736 spherically asymmetric tropical precipitation responses are best characterized
737 by changes in the relative amplitude of the tropical precipitation peaks in
738 the NH and SH. *Adam (2021) demonstrated that common metrics of “ITCZ*
739 *shifts” respond to tropical precipitation changes that are best characterized*
740 *by changes in the relative amplitude of bi-modal tropical precipitation peaks*
741 *in the absence of changes in the location of the peaks. These precipitation*
742 *changes warrant further study from a dynamics perspective and should be*
743 *viewed as distinct* from an ITCZ shift, in its most literal sense.

744 It is possible that our interpretation of the term “ITCZ shift” as a merid-
745 ional translation of the climatological precipitation pattern is overly literal
746 and is different from that intended by the paleoclimate and climate dynam-
747 ics communities. However, our main point here is that a meridional dipole in
748 the pattern of forced precipitation change is not necessarily indicative of (i.e.
749 statistically well described by) a meridional shift of the climatological precipi-
750 tation. We hope our work here motivates future studies to clarify the definition
751 of an “ITCZ shift” and quantify how consistent tropical precipitation changes
752 are with the definition used. The various metrics used in literature to quantify
753 “ITCZ shifts” should be renamed to more accurately represent the patterns
754 of forced precipitation changes and the underlying links to the changes in the
755 Hadley circulation that the metrics respond to.

756 We find that the magnitude of ITCZ shifts quantified by the optimal shift
757 (S) method are generally small – less than 1° latitude in magnitude with the
758 exception of the highly hemispherically asymmetric North Atlantic freshwa-
759 ter forcing (where $|S|$ ranges from 0.4 to 1.7° in response to 1 Sv freshwater
760 hosing). S calculated from the simulated precipitation response to forcing is
761 smaller (by approximately one half) than that inferred using alternative sta-
762 tistical metrics of “ITCZ shifts” such as P_{CENT} (Frierson and Hwang, 2012;
763 Donohoe et al., 2013) colored and the precipitation power weighted latitude
764 (Adam et al., 2016a) .

765 The magnitude of optimal shifts (S) diagnosed in this study provide addi-
766 tional support for the notion that zonal mean ITCZ shifts of more than a couple

degrees latitude are not realized in forced model simulations. Small magnitude ITCZ shifts have also been deduced from paleoclimate proxy sea surface temperature (SST) reconstructions during the LGM, mid-Holocene and across Heinrich events using statistical relationships between the inter-hemispheric SST gradient and the zonal mean ITCZ location derived from modern day observed variability and forced simulations (*McGee et al., 2014*). Compilations of paleo hydrological proxy records throughout the tropics meanwhile have been interpreted as representing larger magnitude ($> 5^\circ$) regional shifts of the ITCZ during the LGM and Little Ice Age (*Arbuszewski et al., 2013; Pahnke et al., 2007; Haug et al., 2001; Sachs et al., 2009*). Simulated ITCZ shifts in response to paleoclimate forcing are highly zonally inhomogeneous (*Atwood et al., 2020*). Together, these results suggest that sparse compilations of proxy records likely can only inform on regional (as opposed to zonal mean) shifts of the ITCZ. Our work here further suggests that, even if we had reliable proxy records for reconstructing *zonal mean precipitation*, such changes should not be interpreted solely in terms of a meridional translation of the climatological precipitation the given the small fraction of simulated ΔP explained by the shift mode relative to the intensity and contraction modes. Indeed, model simulations suggest that a decrease (increase) in precipitation proportional to the mean state precipitation throughout the tropics in response to global mean cooling (warming) is a good zeroth order hypothesis **and has slightly more explanatory power for explaining the spatial structure of tropical hydrological**

789 changes than the thermodynamic constraints on (P-E) changes advanced by
790 *Held and Soden (2006)* as discussed in Section 4.2 .

791 A zonal mean ITCZ shift of 3° (as measured by changes in P_{CENT}) re-
792 quires a sustained hemispheric asymmetry of atmospheric heating equivalent
793 to a simultaneous doubling of CO_2 in one hemisphere and CO_2 halving in
794 the opposing hemisphere (resulting in 1 PW of atmospheric energy transport
795 across the equator) assuming no radiative feedbacks (which generally damp
796 localized forcing *Donohoe and Voigt, 2015*). We cannot comprehend a forcing
797 that could achieve a hemispheric asymmetry in atmospheric heating of the
798 magnitude necessary to shift the ITCZ by more than a couple of degrees –
799 even more so when the damping impact of ocean circulation anomalies is con-
800 sidered (*Green et al., 2019*). The preponderance of modeling, theoretical and
801 empirical observational evidence suggests that large magnitude ($> 3^\circ$) zonal
802 mean ITCZ shifts under paleoclimate forcing are dynamically unrealistic and
803 are unlikely to explain the spatial structure of tropical precipitation changes.

804 We find that contractions and intensifications of the tropical precipita-
805 tion respond robustly and in unison (with fixed scaling) to external forcing:
806 e.g. tropical precipitation intensifies and meridionally contracts under global
807 warming, and reduces and meridionally expands under global cooling. The
808 tight (empirically derived) scaling between contractions and intensifications
809 allow us to define a joint contracting and intensifying mode (*CI*) mode of
810 tropical precipitation that efficiently describes the majority of the tropical
811 precipitation changes under LGM and $4XCO_2$ forcing.

812 What mechanism accounts for the tight coupling between contraction and
813 intensification seen across the ensemble of forcing experiments (Fig. 3B)? The
814 meridional contraction of tropical precipitation in response to global warming
815 under anthropogenic forcing has been noted in several studies (*Lau and Kim,*
816 *2015; Su et al., 2017*). The proposed mechanisms of precipitation contraction
817 under warming include: (i) changes in the efficiency of transient-eddy energy
818 and moisture fluxes (*Byrne and Schneider, 2016a; Nolan et al., 2010; Peters*
819 *et al., 2008*), (ii) cloud radiative effects (*Su et al., 2017*), and (iii) increased
820 gross moist stability (GMS) of the tropics under warming that impacts the
821 meridional width of convective regions (*Chou et al., 2009; Byrne and Schnei-*
822 *der, 2016b*). *Donohoe et al. (2019)* found that a contraction of annual mean
823 precipitation resulted from a reduced seasonal migration of the ITCZ off the
824 equator under global warming and argued that the latter resulted from en-
825 hanced GMS of the tropical atmosphere. *Adam et al. (2016b)* found the mod-
826 els with more meridionally contracted mean state tropical precipitation had
827 enhanced net energy input into the tropical atmosphere which is consistent
828 with enhanced GMS leading to contraction of tropical precipitation. All of
829 these mechanisms would be expected to lead to an expansion of tropical pre-
830 cipitation under global cooling in addition to the contraction seen under global
831 warming as found in the present study. However, it is unclear whether the de-
832 gree of simulated contraction is controlled by the magnitude of tropical surface
833 warming, the GMS change, the degree of precipitation intensification (which
834 would be consistent with our results in Fig. 3B), and/or by changes in season-

835 ality. In future work, we hope to apply the optimal characterization of ΔP to
836 seasonal precipitation changes to disentangle these proposed mechanisms by
837 asking whether annual mean contractions result from contractions during all
838 seasons or, alternatively, reflect seasonal precipitation shifts during the solsti-
839 tial seasons that rectify to an annual contraction. Given the strong seasonality
840 of zonal mean tropical precipitation, we also plan to examine whether the
841 large fraction of ΔP explained by the *CI* extends to seasonal precipitation
842 changes or, rather, the annual mean changes reflect more subtle changes in
843 the seasonality of precipitation.

844 How can our findings here be used to interpret large-scale patterns of past
845 tropical precipitation changes based on compilations of sparsely located pale-
846 oproxy hydroclimate records? One possible path forward is to shift, contract
847 and intensify the observed climatological precipitation pattern (*Xie and Arkin,*
848 1996) to optimally fit (i.e. minimize the RMS mismatch) the discrete proxy
849 data. If we can trust the scaling between the contraction and intensification
850 scalars found across the ensemble of climate model simulations (Fig. 3B), then
851 we can reduce the number of free parameters in the optimal fit to two (the
852 shift and *CI* mode), thus, improving statistical significance of the fit to proxy
853 data. This procedure would allow the optimal shift and *CI* mode change to
854 be quantified in past climates.

855 Our work here has focused exclusively on zonal and annual mean tropical
856 precipitation changes whereas the simulated tropical precipitation response
857 has been shown to feature strong zonal inhomogeneities (*Braconnot et al.,*

858 2007a, 2012) with robust zonally localized responses to forcing (*Atwood et al.*,
859 2020). We plan to extend our optimal shifting, intensifying and contracting
860 framework to characterize zonal asymmetries in the precipitation response
861 to forcing in order to better understand and relate mechanisms of regional
862 precipitation changes to compilations of proxy records.

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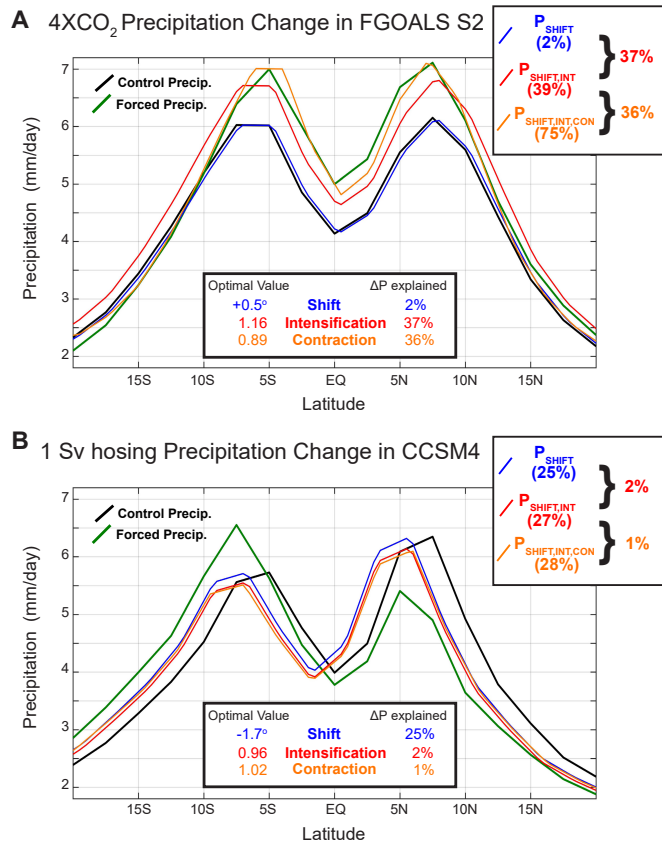


Fig. 1 Illustration of optimal fitting to shifted/contracted/intensified modes of tropical precipitation and the procedure used to define the percent of the precipitation change explained by each mode. The black line shows the climatological (zonal and annual mean) precipitation in the control (pre-industrial) simulations and the green line shows the precipitation after external forcing. First the precipitation is optimally shifted (P_{SHIFT} , blue line). Next the precipitation is optimally intensified and shifted together ($P_{SHIFT,INT}$, red line). Finally, the precipitation is optimally contracted, intensified and shifted together ($P_{SHIFT,INT,CON}$, orange line). The percentage of ΔP explained by each mode is defined as the additional reduction in (normalized) RMS ΔP by successively introducing the mode and all optimal parameters are defined from optimizing the three modes simultaneously. Panel (A) shows an example of a 4XCO₂ precipitation response that is well described by a intensification and contraction. Panel (B) shows the precipitation response to freshwater hosing that is best described by a southward shift.

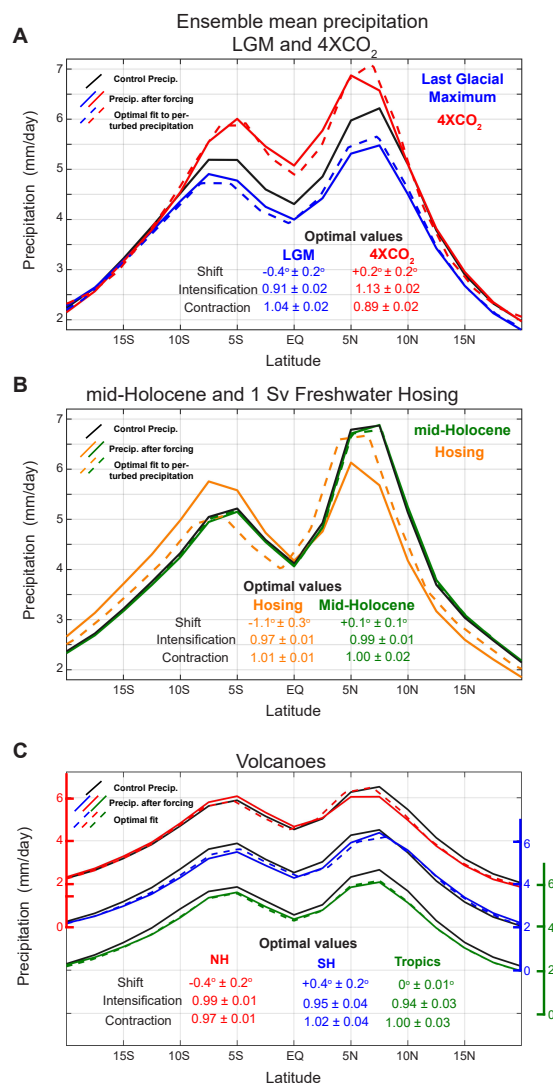


Fig. 2 Ensemble mean, zonal mean and annual mean precipitation for pre-industrial simulations (solid black) and forced simulations (solid colors) with the optimal fit to the forced precipitation shown in dashed colors for the: (A) Last Glacial Maximum (blue) and 4XCO₂ (red) simulations; (B) Mid-Holocene (green) and 1 Sv North Atlantic Freshwater Hosing (orange) and (C) Volcanic forcing simulations subdivided into eruptions originating in the Northern Hemisphere (red), Southern Hemisphere (blue) and tropics (green). Results from each volcanic forcing region have been offset on the ordinate by 2 mm per day (with the PI distribution repeated) for visual purposes as indicate by the colored ordinate axes. Ensemble mean optimal parameters and their inter-model spread (defined as 2 standard deviations of the mean) are shown in the inset tables.

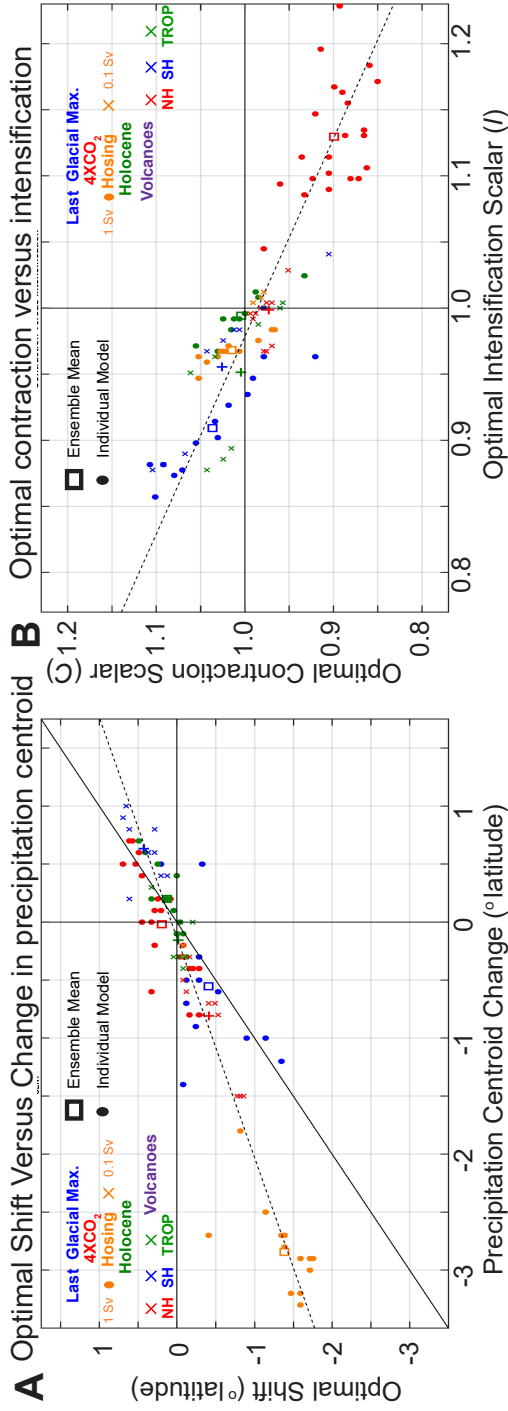


Fig. 3 Optimal shift, contraction and intensification parameters for the zonal and annual mean precipitation changes in response to LGM (blue dots), 4XCO₂ (red dots), Mid-Holocene (green dots), freshwater hosing (orange dots for 1 Sv and crosses for 0.1 Sv) and volcanic forcing which are further subdivided into volcanoes in the Northern Hemisphere (NH – red crosses), Southern Hemisphere (SH – blue crosses) and, Tropics (TROP – green crosses). (A) Optimal shift (ordinate) versus the change in precipitation centroid (abscissa) that is commonly used to quantify ITCZ shifts. The 1:1 line is shown by the solid black line and the linear best fit is shown by the dashed black line. (B) Optimal contraction scalar (ordinate) versus optimal intensification scalar (abscissa). Dots (and smaller crosses for volcanoes) represent individual models and the open squares (larger plus signs for volcanoes) show the ensemble mean of simulations for each forcing experiment. The ensemble mean square for the hosing simulations is the mean of the 1 Sv freshwater forcing simulations only.

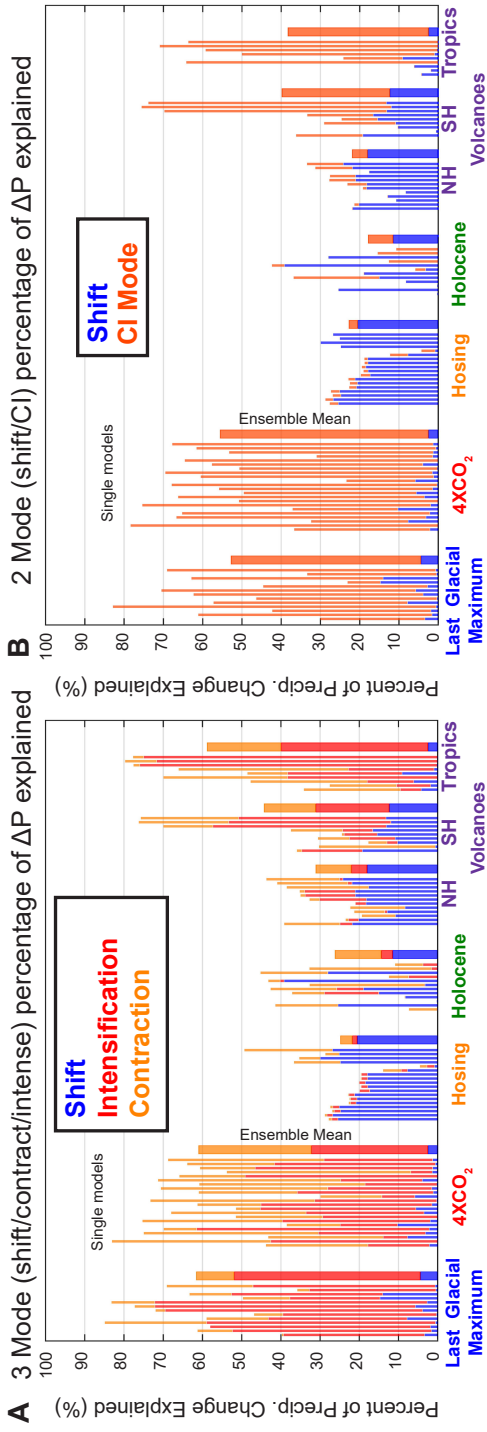


Fig. 4 The percentage of ΔP explained by each mode of tropical precipitation changes in the ensemble of forced simulations. Results are grouped by forcing type with individual models shown by thin bars and the ensemble average shown in the thick bar to the right. (A) Three mode optimization with percentages defined starting from the shifting mode only (blue) then adding the intensification mode (red) and the adding the contraction (orange). (B) Two mode optimization with percentages defined starting from the shifting mode (blue) then adding the combined contraction-intensification mode (CI—reddish orange) with the ratio of the contraction to intensification defined from the linear best fit across all simulations shown in Fig. 1B.

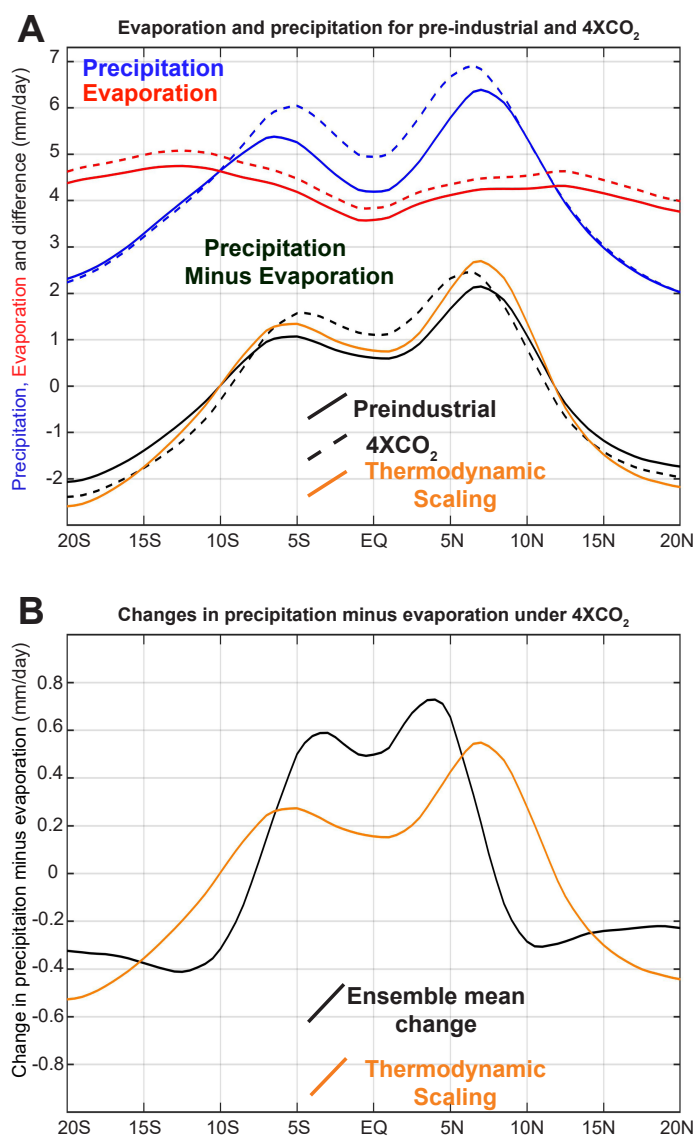


Fig. 5 (A) Ensemble mean, zonal mean and annual mean precipitation (blue), evaporation (red) and precipitation minus evaporation (black) for the pre-industrial simulations (solid lines) and 4XCO₂ simulations (dashed lines). The thermodynamic scaling of *Held and Soden* (2006) which predicts precipitation minus evaporation will scale with water vapor content of the atmosphere is shown in the orange line. (B) The ensemble mean change precipitation minus evaporation (black) compared to the thermodynamic prediction.